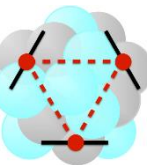


# Nd elastic scattering and the subleading short-range 3NF

R.Skibiński, J.Golak, H.Witała



JAGIELLONIAN  
UNIVERSITY



## LENPIC Collaboration

Jagiellonian University, Kraków  
Ruhr-Universität, Bochum  
Forschungszentrum, Jülich  
Bonn Universität,

Ohio State University  
Iowa State University  
Technische Universität, Darmstadt  
Kyutech, Fukuoka  
IPN, Orsay  
TRIUMF, Vancouver

[www.lenpic.org](http://www.lenpic.org)

24-26.08.2022,  
LENPIC Meeting, Bochum  
1

# Outline

Emulator for the Nd scattering:

Based on:

H.Wiśniewski et al., *Few-Body Syst.* 62 (2021) 23 „*Perturbative Treatment of Three Nucleon Force Contact Terms in Three-Nucleon Faddeev Equations.*”

H.Wiśniewski et al., *Eur. Phys. J. A* 57 (2021) 241 „*Efficient emulator for solving 3N continuum Faddeev equations with chiral 3NF comprising any number of contact terms.*”

H.Wiśniewski et al., *Phys. Rev. C* 105 (2022) 054004 „*Significance of chiral 3NF contact terms for understanding of elastic nucleon-deuteron scattering*”

1. Formalism – new set of equations
2. Tests and the first results on fixing short-range 3NF parameters

# Emulator for Nd scattering

Our standard method to calculate transition amplitude for 3N scattering is to solve:

- the Schrodinger equation -> deuteron

$$(H_0 + V)\Psi_d = E_d \Psi_d$$

- the Lippmann-Schwinger equation -> t-matrix

$$t(E) = V + VG_0(E)V + VG_0VG_0(E)V + \dots = V + VG_0t(E)$$

- the Faddeev equation -> auxiliary operator T

$$T\phi = tP\phi + (1+tG_0)V_{123}^{(1)}(1+P)\phi + tPG_0T\phi + (1+tG_0)V_{123}^{(1)}(1+P)G_0T\phi$$

where  $\phi$  is the deuteron wf times free nucleon state

- Compute amplitudes for :

$$\text{elastic scattering: } U = PG_0^{-1} + V_{123}^{(1)}(1+P)\phi + PT + V_{123}^{(1)}(1+P)G_0T$$

$$\text{deuteron breakup: } U_0 = (1+P)T$$

# Emulator for Nd scattering

$$T\varphi = tP\varphi + (1 + tG_0)V_{123}^{(1)}(1 + P)\varphi + tPG_0T\varphi + (1 + tG_0)V_{123}^{(1)}(1 + P)G_0T\varphi$$

We work in the PWD scheme

$$\begin{aligned} \text{3N state: } \quad |pq\alpha\rangle &\equiv |pq(ls)j(\lambda \frac{1}{2})I(jI)JM_J(t \frac{1}{2})TM_T\rangle & \vec{p} &= \frac{1}{2}(\vec{p}_2 - \vec{p}_3) \\ & & \vec{q} &= \frac{1}{3}(2\vec{p}_1 - \vec{p}_2 - \vec{p}_3) \end{aligned}$$

what means, that PWD of appearing in this equation operators has to be performed.

$$\langle p'q'(l's')j'(\lambda' \frac{1}{2})I'(j'I')J'M_{J'}(t' \frac{1}{2})T'M_{T'} | \hat{O} | pq(ls)j(\lambda \frac{1}{2})I(jI)JM_J(t \frac{1}{2})TM_T \rangle$$

After decomposing 3NF:

CPU time required for one run (i.e. one reaction energy, one NN+3NF potential) amounts from approx. 2-10 hrs., depending on number of partial waves, computer parameters, disk space available). Some hardware (GPU, fast memory) or software (parallelization) improvements are still possible but the cake's not worth the candle.

## Fixing parameters of 3NF

- Up to now, i.e. when working at N2LO there are only two free parameters  $c_D$  and  $c_E$ .
- Typically  $^3\text{H}$  and the  $^2a_{\text{nd}}$  or the differential Nd elastic scattering cross section at one or few energies are used.  
The latter requires solving the triton many times and the Faddeev equation 10-20 times.
- However, now we expect:
- No new 3NF free parameters at N3LO, but three new offshell LECs in the chiral NN force.
- 13 contact terms at N4LO (more precisely, due to some identities between operators, one expects in total 13 free parameters of 3NF at N4LO).
- Thus finding an efficient emulator for solving the 3N Faddeev equation seems to be essential and of high priority.

# Emulator for Nd scattering – new results

- In [1] H.Wiłała et al., Few-Body Syst. 62 (2021) 23  
we proposed such an emulator which enabled us to reduce significantly the required time of computations. We tested its efficiency as well as ability to accurately reproduce exact solution of the 3N Faddeev equations.
- In [2] H.Wiłała et al., Eur. Phys. J. A 57 (2021) 241  
we introduced a computational scheme based on the perturbative approach of Ref. [1], which even by far more reduced the computer time necessary to obtain the observables in the elastic nucleon-deuteron scattering and deuteron breakup reactions, and which is well suited for calculations with varying strengths of the contact terms in a chiral 3NF
- In [3] H.Wiłała et al., Phys. Rev. C105 (2022) 054004  
we used the SMS N4LO+ NN potential in combination with the N2LO chiral 3NF supplemented by all the N4LO contact terms. Our aim was to verify if it would be possible to fix strengths of all the contact terms by performing a least squares fit of theory to Nd elastic-scattering data.

# Emulator for Nd scattering - algorithm

- The contact terms are restricted to small 3N total angular momenta and to only few partial-wave states for a given total 3N angular momentum J and parity  $\pi$

- Let us split 3NF  $\theta = \{c_1, c_2, \dots, c_n\}$   
 $V_{123}^{(1)} = V(\theta_0) + \Delta V(\theta) \equiv V(\theta_0) + \sum_{i=1}^n c_i \Delta V_i$   $\theta_0 = \{0, 0, \dots, 0\}$

- We divide the 3N partial-wave states into two sets:
  1. The  $\beta$  set is defined by non-vanishing matrix elements of parameters dependent short-range 3NF:  $\Delta V(\theta)$ .
  2. The  $\alpha$  set comprises remaining states.

- Similarly to 3NF  
 $T = T(\theta_0) + \Delta T(\theta)$

# Emulator for Nd scattering - algorithm

- Inserting this to the Faddeev equation leads to sets of equations:
- For  $\langle \alpha |$

$$\begin{aligned}\langle \alpha | T(\theta_0) | \phi \rangle &= \langle \alpha | t P | \phi \rangle + \langle \alpha | (1 + t G_0) V(\theta_0) (1 + P) | \phi \rangle + \langle \alpha | t P G_0 T(\theta_0) | \phi \rangle \\ &\quad + \langle \alpha | (1 + t G_0) V(\theta_0) (1 + P) G_0 T(\theta_0) | \phi \rangle \\ \langle \alpha | \Delta T(\theta) | \phi \rangle &= \langle \alpha | t P G_0 \Delta T(\theta) | \phi \rangle + \langle \alpha | (1 + t G_0) V(\theta_0) (1 + P) G_0 \Delta T(\theta) | \phi \rangle .\end{aligned}\quad (9)$$

- For  $\langle \beta |$

$$\begin{aligned}\langle \beta | T(\theta_0) | \phi \rangle &= \langle \beta | t P | \phi \rangle + \langle \beta | (1 + t G_0) V(\theta_0) (1 + P) | \phi \rangle + \langle \beta | t P G_0 T(\theta_0) | \phi \rangle \\ &\quad + \langle \beta | (1 + t G_0) V(\theta_0) (1 + P) G_0 T(\theta_0) | \phi \rangle \\ \langle \beta | \Delta T(\theta) | \phi \rangle &= \langle \beta | (1 + t G_0) \Delta V(\theta) (1 + P) | \phi \rangle + \langle \beta | (1 + t G_0) \Delta V(\theta) (1 + P) G_0 T(\theta_0) | \phi \rangle \\ &\quad + \langle \beta | (1 + t G_0) [V(\theta_0) + \Delta V(\theta)] (1 + P) G_0 \Delta T(\theta) | \phi \rangle \\ &\quad + \langle \beta | t P G_0 \Delta T(\theta) | \phi \rangle .\end{aligned}\quad (10)$$

- the first equations of both sets of equations are identical, only for different elements, and together give the standard Faddeev equation (with  $V(\theta_0)$  )
- We neglect term  $\sim \Delta V \Delta T$



# Emulator for Nd scattering - algorithm

- We neglect term  $\sim \Delta V \Delta T$ 

$$\begin{aligned} \langle \beta | \Delta T(\theta) | \phi \rangle = & \langle \beta | (1 + tG_0) \Delta V(\theta) (1 + P) | \phi \rangle \\ & + \langle \beta | (1 + tG_0) \Delta V(\theta) (1 + P) G_0 T(\theta_0) | \phi \rangle \\ & + \langle \beta | (1 + tG_0) V(\theta_0) (1 + P) G_0 \Delta T(\theta) | \phi \rangle \\ & + \langle \beta | t P G_0 \Delta T(\theta) | \phi \rangle. \end{aligned} \quad (6)$$

- For single parameter dependent component of  $V$ :  $V_i = c_i V$  we may solve that equation separately at  $c_i = 1$  obtaining corresponding  $\Delta T_i$ 

$$\begin{aligned} \langle \beta | \Delta T_i | \phi \rangle \equiv & \langle \beta | (1 + tG_0) \Delta V_i (1 + P) | \phi \rangle \\ & + \langle \beta | (1 + tG_0) \Delta V_i (1 + P) G_0 T(\theta_0) | \phi \rangle \\ & + \langle \beta | (1 + tG_0) V(\theta_0) (1 + P) G_0 \Delta T_i | \phi \rangle \\ & + \langle \beta | t P G_0 \Delta T_i | \phi \rangle, \end{aligned} \quad (7)$$

- Finally:
$$\langle \beta | \Delta T(\theta) | \phi \rangle = \sum_{i=1}^N c_i \langle \beta | \Delta T_i | \phi \rangle. \quad (8)$$

- Summarizing: one needs to solve  $N+1$  Faddeev equations (one for  $T(\theta_0)$  and  $N$  for  $\langle \beta | \Delta T_i | \phi \rangle$ ), next  $N$  times find  $\langle \alpha | \Delta T_i | \phi \rangle$  by integration.

# Emulator for Nd scattering - algorithm

- In this way we have matrix elements of T

$$\langle \alpha | T(\theta) | \phi \rangle = \langle \alpha | T(\theta_0) | \phi \rangle + \sum_i c_i \langle \alpha | \Delta T_i | \phi \rangle,$$

$$\langle \beta | T(\theta) | \phi \rangle = \langle \beta | T(\theta_0) | \phi \rangle + \sum_i c_i \langle \beta | \Delta T_i | \phi \rangle. \quad (10)$$

- Let us now come back to the scattering amplitudes

$$U = P G_0^{-1} + V_{123}^{(1)} (1 + P) \phi + P T + V_{123}^{(1)} (1 + P) G_0 T$$

$$U_0 = (1 + P) T$$

- They are linear in T: the dependence on the  $c_i$  constants carries over to them, except as a complication for elastic scattering, but they can be written as

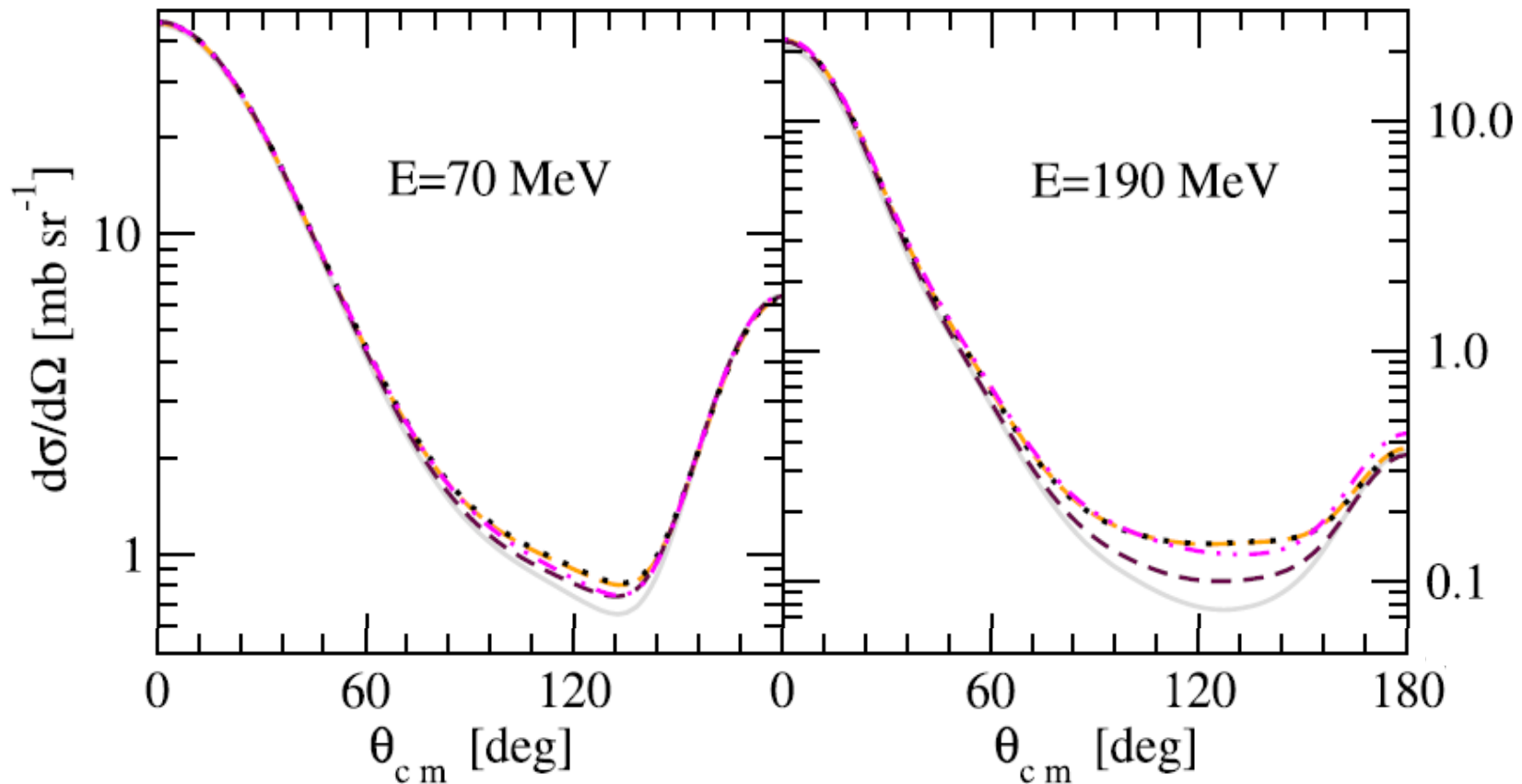
$$U = U(\theta_0) + \sum_i c_i U_i + \sum_{i,k} c_i c_k U_{ik}$$

$$U_0 = U_0(\theta_0) + \sum_i c_i U_{0i}$$

# Emulator for Nd scattering – algorithm - application

- We used SMS N4LO+ NN potential at  $\Lambda=450$  MeV, combined with the N2LO chiral 3NF and supplemented by all subleading N4LO 3NF contact terms from:
  1. L. Girlanda, A. Kievsky, and M. Viviani, Phys. Rev. C 84, 014001 (2011).,
  2. L. Girlanda, A. Kievsky, and M. Viviani, Phys. Rev. C 102, 019903(E) (2020).
- All terms are regulated with the non-local regulator.
- Such a Hamiltonian comprises altogether 15 short-range contributions to 3NF, two from N2LO with the strengths  $c_D$  and  $c_E$ , and thirteen from N4LO with the strengths  $E_i$ ,  $i = 1, \dots, 13$ . However, for two pairs of the  $E_i$  terms matrix elements are identical, thus finally there are 13 unknown parameters.

# Emulator for Nd scattering – test



Exact:

NN N4LO+

NN N4LO+ + 3NF N2LO+E7

( $c_D = -8.2053$ ,  $c_E = -1.0019$ ,  $c_{E7} = 2.0$ )

NN N4LO+ + 3NF N2LO ( $c_D = c_E = c_{E7} = 0.0$ )

Emulator:

NN N4LO+ + 3NF N2LO+E7 ( $\beta = {}^1S_0, {}^3S_1, {}^3D_1$ )

NN N4LO+ + 3NF N2LO+E7 ( $\beta = j \leq 2$ )

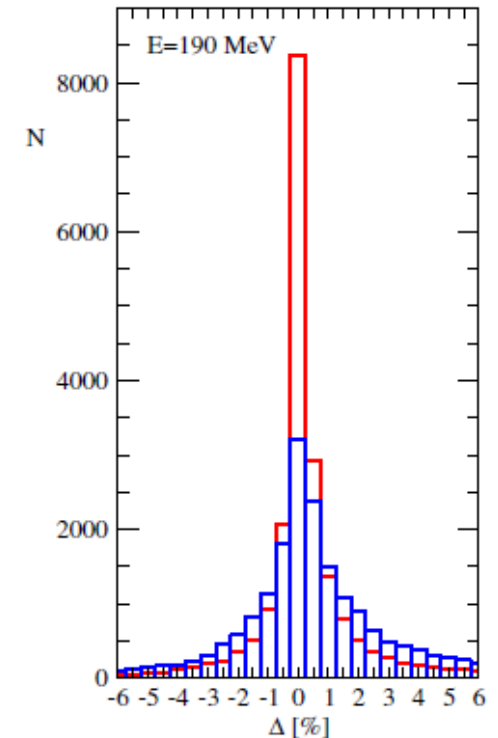
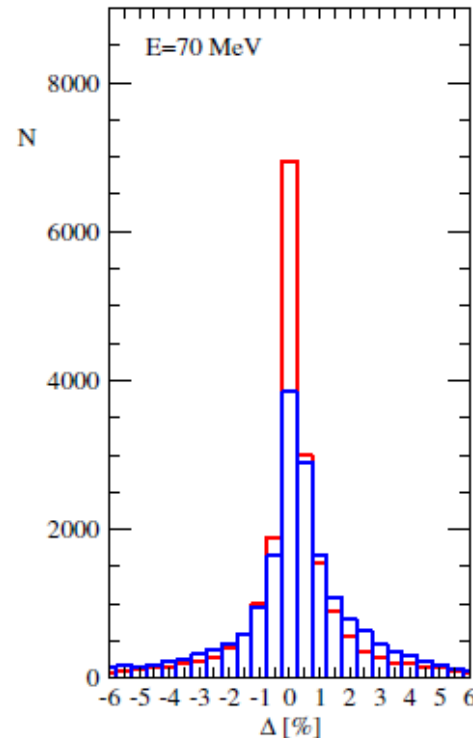
# Emulator for Nd scattering – test

- Histogram of

$$\Delta \equiv \frac{Obs(emulator) - Obs(exact)}{Obs(exact)}$$

over all Nd scattering observables and grid of scattering angles and for 6 sets of (cD,cE)

- NN N4LO+ + 3NF N2LO
- **Blue** – the final emulator but with only D and E terms
- **Red** – without neglecting the two terms  $\sim \Delta V \Delta T$  (slightly better but slower)



# Emulator for Nd scattering – algorithm - application

Sensitivity of 3N scattering observables to  $E_i$  terms

Green circles  $V(\theta_0)$

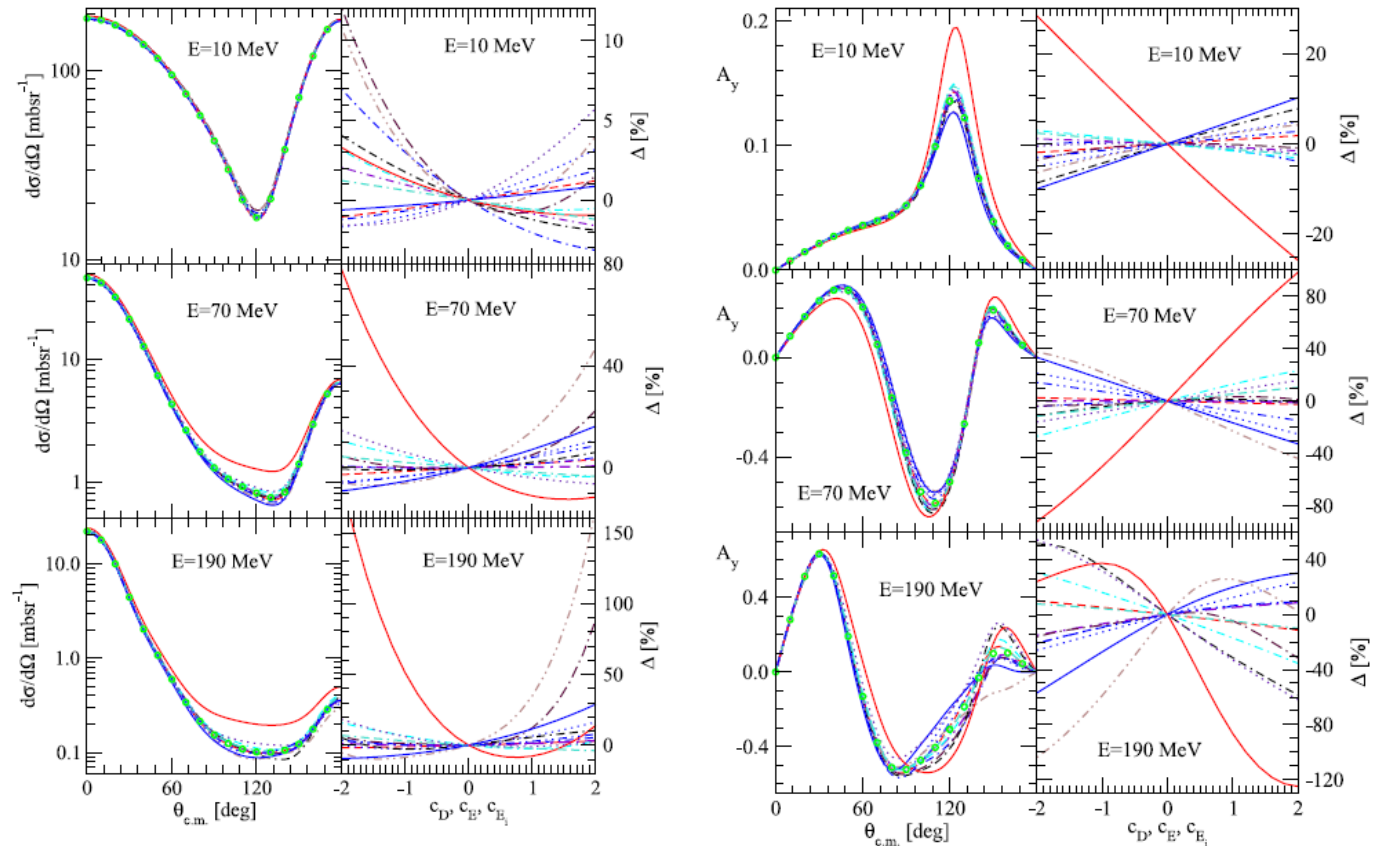
$C_i = -1$  (left)

red solid  $E_8$ ,

blue solid  $E_7$ ,

brown dashed-

double dotted  $E_5$



- N2LO D and E terms do not dominate

- Some observables are more sensitive to specific terms, e.g.  $T_{22}$  to  $E_{10}$

$$\Delta \equiv \Delta(c_i) = \frac{1}{N_\theta} \sum_{\theta_k} \frac{Obs(c_i, \theta_k) - Obs(\theta_0, \theta_k)}{Obs(\theta_0, \theta_k)}$$

# Emulator for Nd scattering – application

## $V_i$ expectation values in ${}^3\text{H}$ at $c_i=1.0$

TABLE I. Contributions of the  $\text{N}^2\text{LO}$  and  $\text{N}^4\text{LO}$  contact terms to the potential energy of the three nucleons in the triton. These expectation values were obtained for the  ${}^3\text{H}$  wave function calculated with the SMS chiral  $\text{N}^4\text{LO}^+$   $NN$  potential ( $\Lambda = 450$  MeV) and assuming strengths of contact terms  $c_i = 1.0$ .

| $V_i$     | $\langle \psi_{{}^3\text{H}}   V_i   \psi_{{}^3\text{H}} \rangle$<br>[MeV] |
|-----------|--|
| $V_D$     | 0.1661   |
| $V_E$     | -1.4294  |
| $V_{E1}$  | 0.3463   |
| $V_{E2}$  | -0.4173  |
| $V_{E3}$  | -0.2754  |
| $V_{E4}$  | -1.0390  |
| $V_{E5}$  | -0.9559  |
| $V_{E6}$  | -1.0699  |
| $V_{E7}$  | $0.1798 \times 10^{-4}$  |
| $V_{E8}$  | $0.8817 \times 10^{-2}$  |
| $V_{E9}$  | -0.2407  |
| $V_{E10}$ | 1.0571   |
| $V_{E11}$ | -0.2407  |
| $V_{E12}$ | 1.0571   |
| $V_{E13}$ | 0.3060   |

Relative strengths

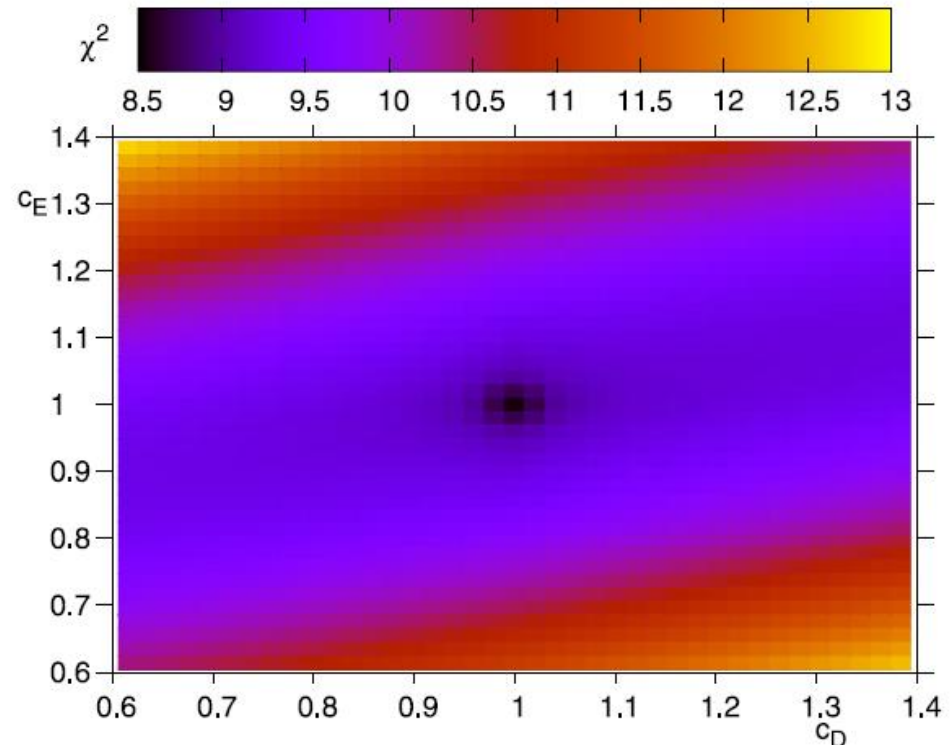
Nearly all terms are important (for  $c_i=1$ ) with exception of  $E_7$  and  $E_8$  terms

# Emulator for Nd scattering – fit to the data - test

- Pseudodata: all  $c_i=1$ , 5 energies: 10,70,135,190,250 MeV, cross section,  $A_Y(N), A_Y(d), T_{20}, T_{21}, T_{22}, \theta_{cm}$  in (40-170) deg, step 5 deg, „experimental” uncertainty 5%.

- The presented plot uses pseudodata at  $E=70$  MeV.

- Reproducing gives  $c_D=1.00\pm 0.08$ ,  
 $c_E=1.00\pm 0.03$   
and  $\chi^2/\text{data}\approx 0.0$   
with a single clear minimum



- Exercise: Pseudodata (3NF N2LO with  $c_D=c_E=1.0$ ) but fit incomplete theory gives  $\chi^2/\text{data}\approx 80$  and  $c_D=3.98\pm 0.08$ ,  $c_E=2.99\pm 0.03$  (see PRC paper)



# Emulator for Nd scattering – fit to the true data at 10,70, and 135 MeV (786 data points)

TABLE III. The values of strengths  $c_i$  found in the least squares fit to the data from Table II at the three energies  $E = 10, 70$ , and 135 MeV.

|              |                  |
|--------------|------------------|
| $c_D$        | $-1.49 \pm 0.06$ |
| $c_E$        | $-1.27 \pm 0.06$ |
| $c_{E_1}$    | $6.40 \pm 0.33$  |
| $c_{E_2}$    | $7.80 \pm 0.36$  |
| $c_{E_3}$    | $6.97 \pm 0.34$  |
| $c_{E_4}$    | $-2.06 \pm 0.13$ |
| $c_{E_5}$    | $-0.36 \pm 0.05$ |
| $c_{E_6}$    | $0.52 \pm 0.03$  |
| $c_{E_7}$    | $-7.40 \pm 0.14$ |
| $c_{E_8}$    | $-2.61 \pm 0.05$ |
| $c_{E_9}$    | $-4.59 \pm 0.22$ |
| $c_{E_{10}}$ | $-0.98 \pm 0.05$ |
| $c_{E_{13}}$ | $-1.14 \pm 0.05$ |

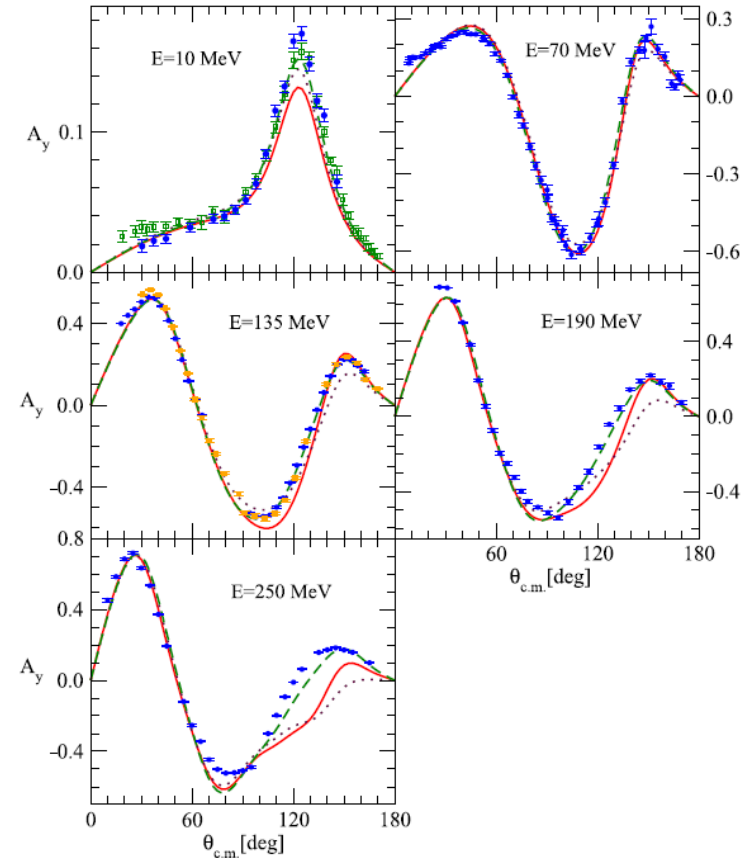
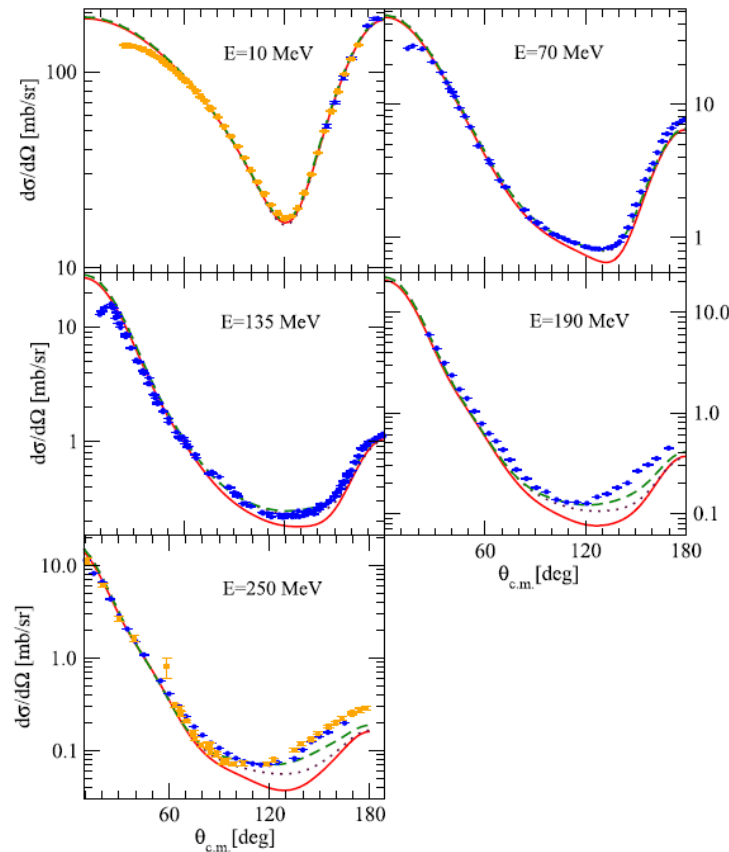
TABLE IV. The covariance matrix for the strengths  $c_i$  determined by the least squares fit of data from Table II at the three energies  $E = 10, 70$ , and 135 MeV [the values shown are  $\text{Cov}(c_i, c_j) \times 1000$ ].

|              | $c_D$ | $c_E$  | $c_{E_1}$ | $c_{E_2}$ | $c_{E_3}$ | $c_{E_4}$ | $c_{E_5}$ | $c_{E_6}$ | $c_{E_7}$ | $c_{E_8}$ | $c_{E_9}$ | $c_{E_{10}}$ | $c_{E_{13}}$ |
|--------------|-------|--------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|-----------|--------------|--------------|
| $c_D$        | 3.914 | -0.456 | 1.412     | 4.573     | 0.843     | 0.844     | -0.729    | -0.892    | 1.109     | 0.267     | -0.726    | 0.123        | -0.207       |
| $c_E$        |       | 3.560  | 0.947     | -3.571    | 1.345     | -0.633    | -0.172    | -0.217    | -2.416    | -0.809    | -1.702    | 0.393        | 0.571        |
| $c_{E_1}$    |       |        | 108.9     | 112.8     | 108.9     | -35.13    | 1.409     | -2.418    | 25.92     | 7.513     | 12.99     | 3.861        | 0.443        |
| $c_{E_2}$    |       |        |           | 130.7     | 113.4     | -35.15    | -1.995    | -3.241    | 32.43     | 9.561     | -0.534    | 0.763        | -3.332       |
| $c_{E_3}$    |       |        |           |           | 112.9     | -38.92    | 1.617     | -1.814    | 27.52     | 8.068     | 8.366     | 1.598        | -0.193       |
| $c_{E_4}$    |       |        |           |           |           | 15.97     | -1.966    | -0.362    | -10.50    | -3.198    | -4.866    | 0.345        | -0.222       |
| $c_{E_5}$    |       |        |           |           |           |           | 2.415     | 0.669     | 0.791     | 0.281     | 9.892     | 1.311        | 1.766        |
| $c_{E_6}$    |       |        |           |           |           |           |           | 0.635     | -0.874    | -0.226    | 1.426     | -0.226       | 0.210        |
| $c_{E_7}$    |       |        |           |           |           |           |           |           | 20.33     | 6.455     | 3.464     | -0.324       | -1.463       |
| $c_{E_8}$    |       |        |           |           |           |           |           |           |           | 2.071     | 1.041     | -0.158       | -0.462       |
| $c_{E_9}$    |       |        |           |           |           |           |           |           |           |           | 50.23     | 9.133        | 8.813        |
| $c_{E_{10}}$ |       |        |           |           |           |           |           |           |           |           |           | 2.625        | 1.910        |
| $c_{E_{13}}$ |       |        |           |           |           |           |           |           |           |           |           |              | 2.499        |

- Big values of  $c_{E_1}, c_{E_2}, c_{E_3}, c_{E_7}, c_{E_9}$
- Correlation coefficients close to  $\pm 1$ :  $\rho(E_1, E_2), \rho(E_2, E_3), \rho(E_1, E_3), \rho(E_3, E_4), \rho(E_7, E_8)$
- Correlation coefficients close to 0:  $(c_D, c_E), (c_D, c_{E_i}), (c_E, c_{E_i})$
- $\chi^2/\text{data} \approx 35$

# Emulator for Nd scattering – fit to the data: cross section and $A_Y(N)$

- Data at 10, 70 and 135 MeV
- Results at 190 and 250 MeV are predictions



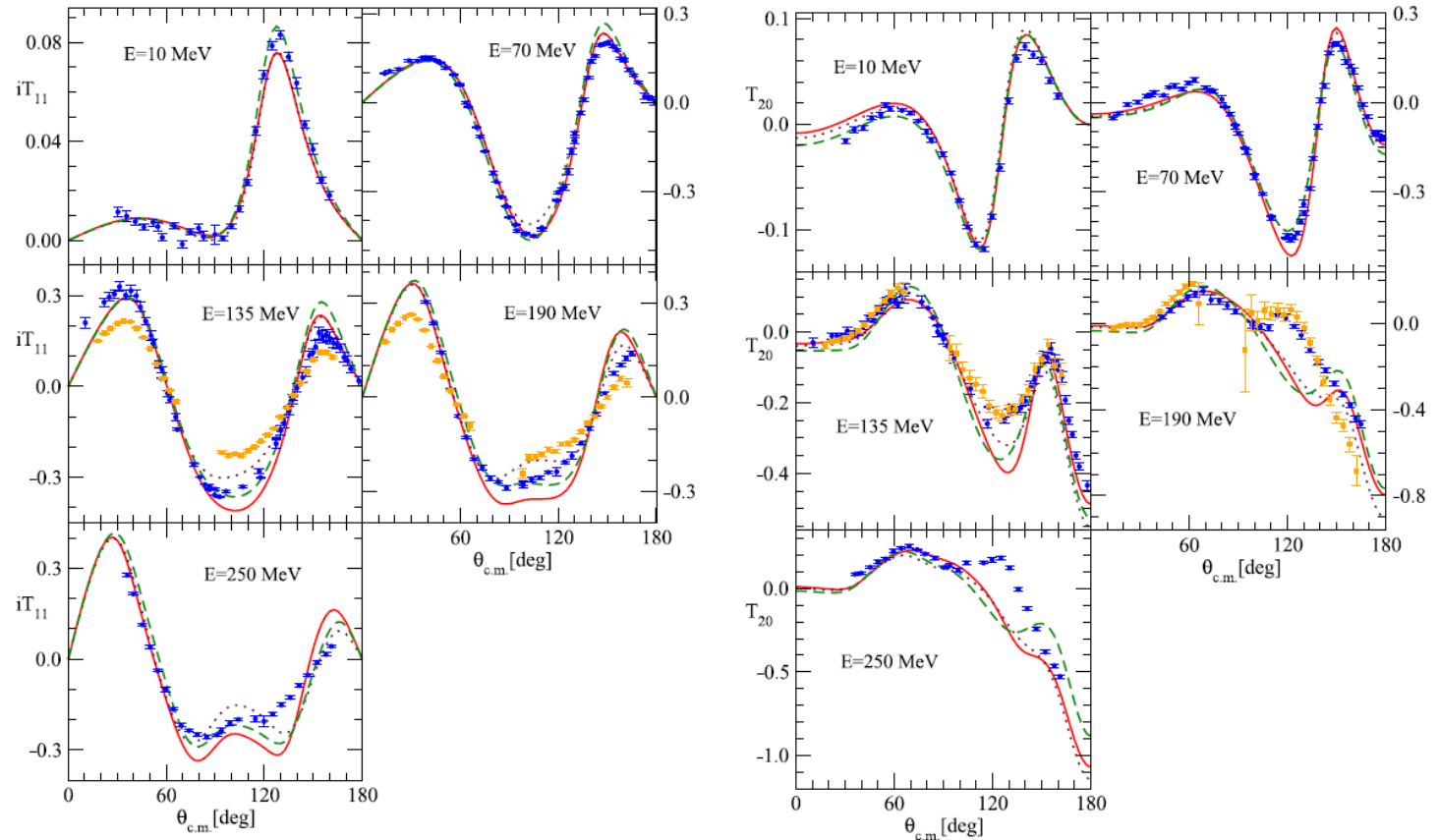
NN N4LO+

..... NN N4LO+ + 3NF N2LO

----- NN N4LO+ + 3NF N2LO +  $E_i$

# Emulator for Nd scattering – fit to the data: $iT_{11}$ and $T_{20}$

- Data at 10, 70 and 135 MeV
- Results at 190 and 250 MeV are predictions



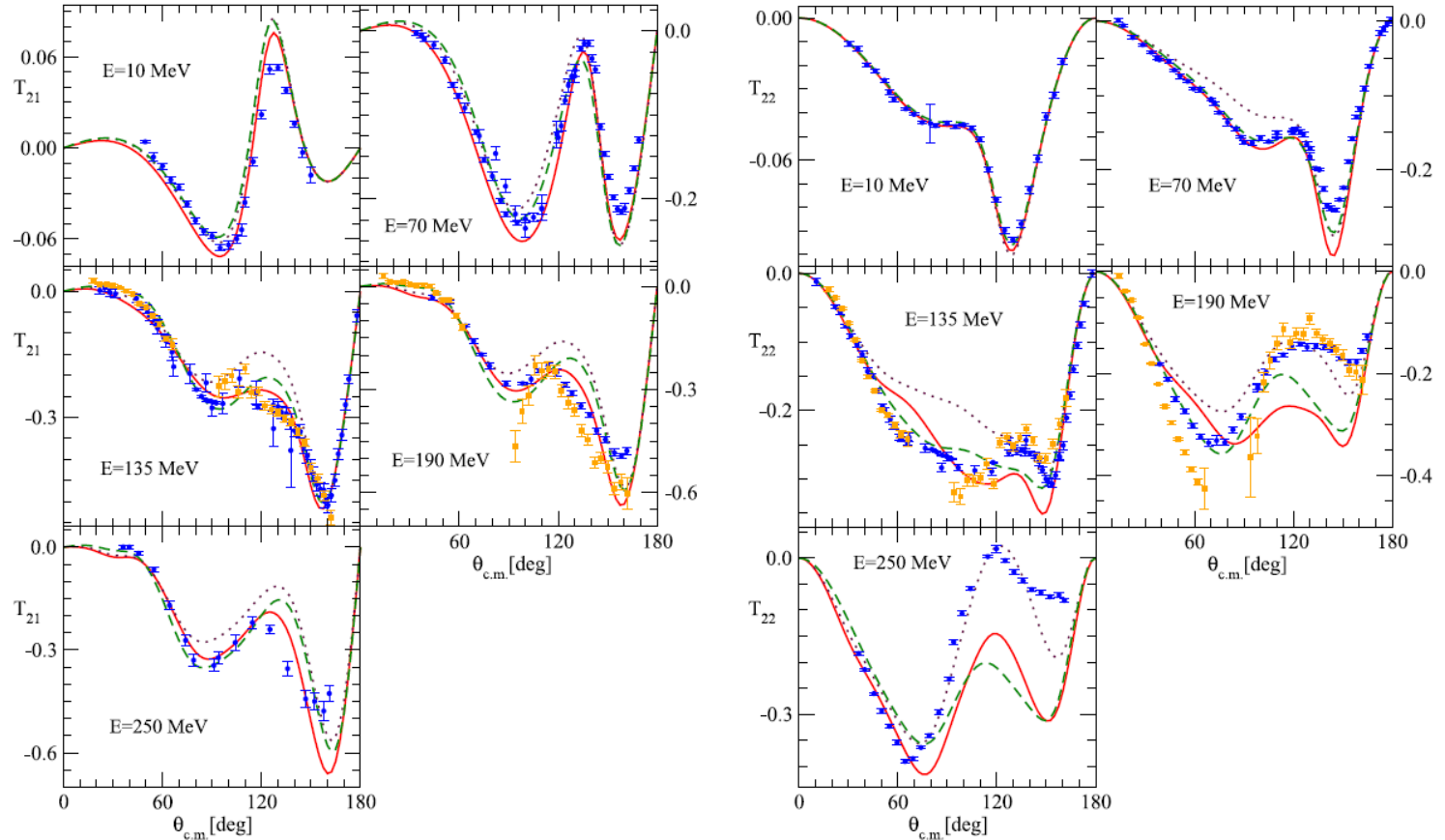
NN N4LO+

..... NN N4LO+ + 3NF N2LO

----- NN N4LO+ + 3NF N2LO +  $E_i$

# Emulator for Nd scattering – fit to the data: $T_{21}$ and $T_{22}$

- Data at 10, 70 and 135 MeV
- Results at 190 and 250 MeV are predictions



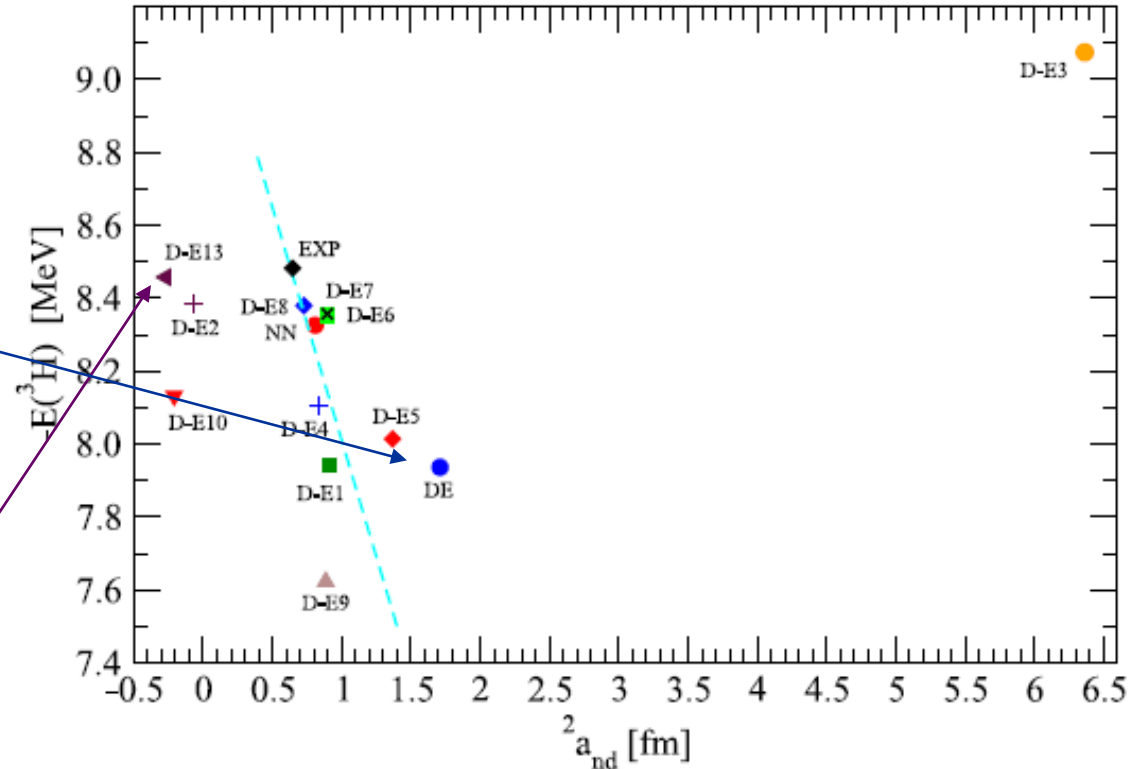
NN N4LO+

..... NN N4LO+ + 3NF N2LO

..... NN N4LO+ + 3NF N2LO + E<sub>i</sub>

# Emulator for Nd scattering – fit to the data: the Phillips line

- EXP (0.645, 8.4820)
- Phenomenological models
- N4LO+ + 3NF N2LO fit  $\equiv$  DE: (1.71, 7.937)
- N4LO+ +3NF N2LO + all  $E_i \equiv$  D-E13 (-0.27046, 8.4581)



There is a problem with simultaneous description of the  $^3H$  and  $^2a_{nd}$

# Summary

- We constructed and tested an efficient and accurate emulator for solving 3N Faddeev equation.
- We applied it to the Nd scattering up to  $E=250$  MeV, using the chiral SMS NN potential at N4LO+ supplemented by 3NF at N2LO and 13 N4LO contact terms.
- Our emulator allows us to fix free parameters of all short-range terms in the 3NF. We found that even at low energies some observables are sensitive to N4LO 3NF contact terms.
- In general, sensitivity of predictions to N4LO 3NF contact terms depends on observable, energy and scattering angle.
- Usually we observe improvements in data description, but very likely above  $\approx 200$  MeV 3NF is not sufficient to explain discrepancies with the data.
- The deuteron breakup data can be used in fitting as well.
- Coulomb correction (if needed) and 3NF at N3LO has to be included for final conclusions.