

Photo: courtesy of Skiathos Palace



Lightcone 2013⁺

Venturing off the lightcone • local vs. global features

Topics

Wilson lines off the lightcone • Lightcone field theory • AdS/CFT theory and applications
Hadron structure in OCD • Effective theories • Lattice simulations

May 20-24, 2013, Skiathos, Greece

Wigner Distributions

in

Light-Cone Quark Models



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Pavia U. & INFN, Pavia (Italy)



Outline

Wigner Distributions
Parton distributions in the Phase Space

$$\text{FT } \vec{b}_\perp \leftrightarrow \vec{\Delta}_\perp$$



Generalized Transverse Momentum Dependent Parton Distributions (GTMDs)

GPDs

TMDs

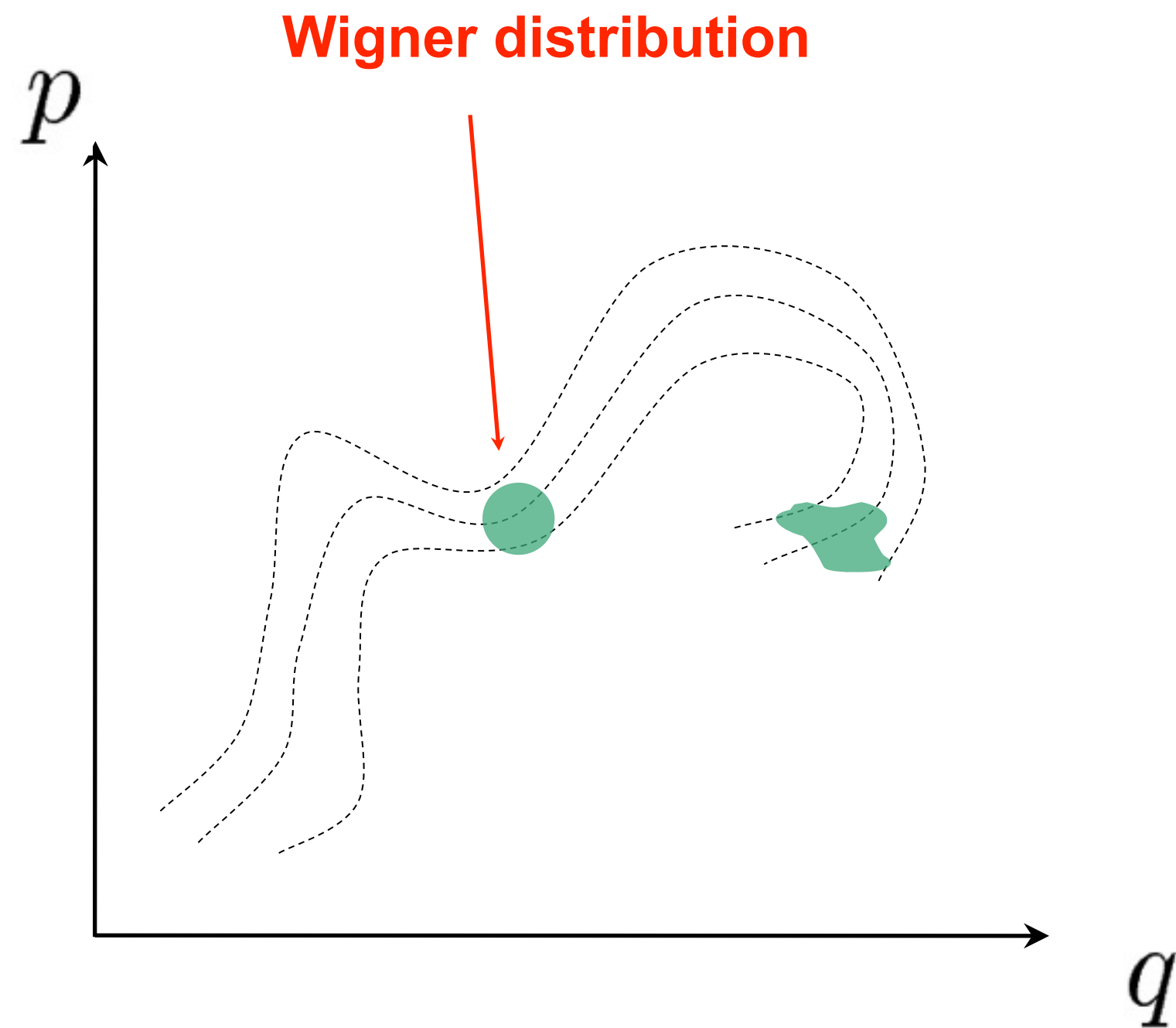
spin and orbital angular momentum structure of the nucleon

insights from quark model calculations

Phase-space distribution

Quantum Mechanics

[Wigner (1932)]
[Moyal (1949)]



Position-space density

$$|\psi(q)|^2 = \int dp P_W(p, q)$$

Momentum-space density

$$|\varphi(p)|^2 = 2\pi \int dq P_W(p, q)$$

Quantum average

$$\langle \hat{O} \rangle = \int dp dq O(p, q) P_W(p, q)$$

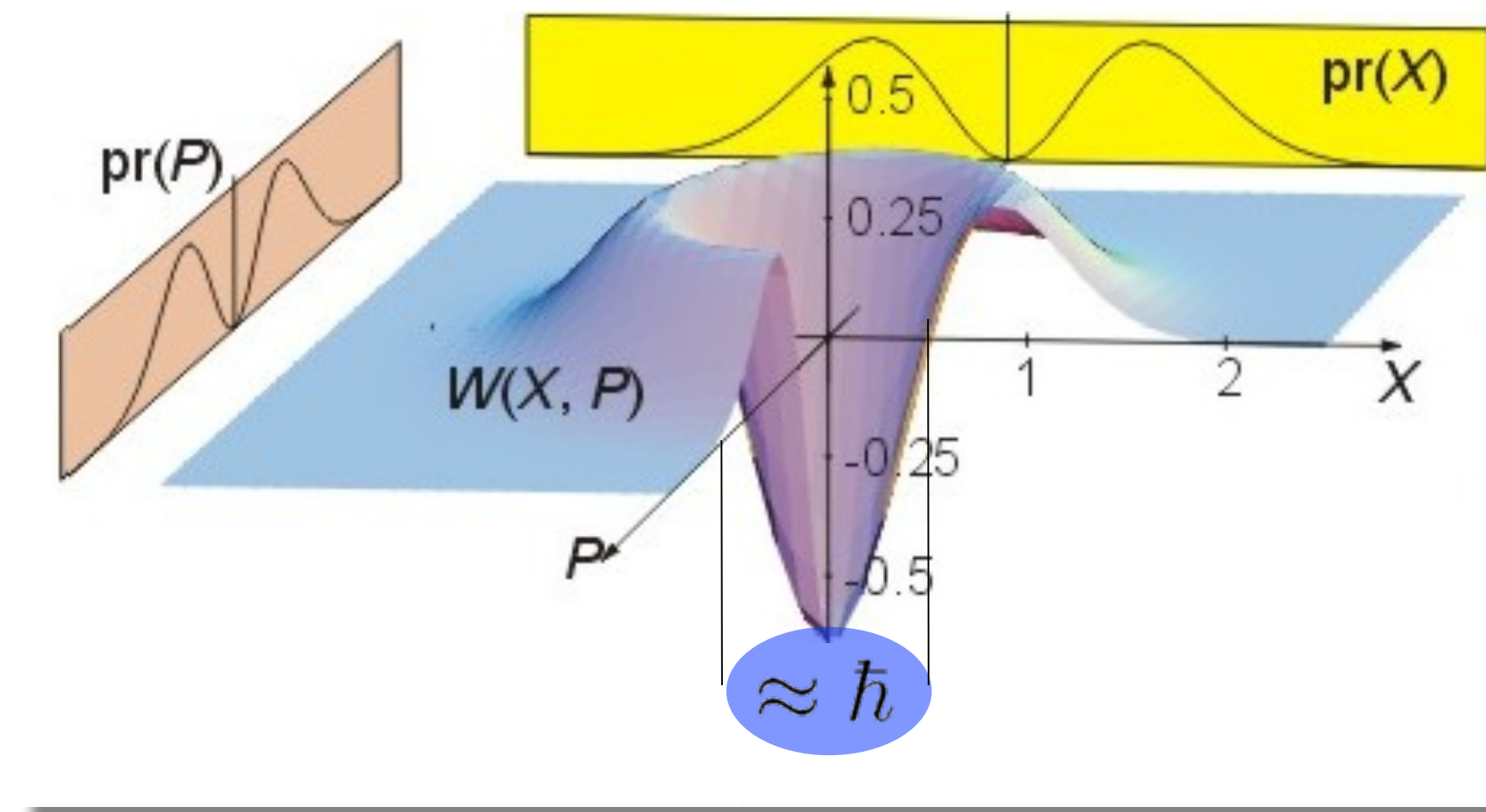
$$\begin{aligned} P_W(p, q) &= \int \frac{dz}{2\pi} e^{-ipz} \psi^*(q - \frac{z}{2}) \psi(q + \frac{z}{2}) \\ &= \int \frac{d\Delta}{(2\pi)^2} e^{-iq\Delta} \varphi^*(p + \frac{\Delta}{2}) \varphi(p - \frac{\Delta}{2}) \end{aligned}$$

Phase-space distribution

Wigner distribution

Numerous applications in

- Nuclear physics
- Quantum chemistry
- Quantum molecular dynamics
- Quantum information
- Quantum optics
- Classical optics
- Signal analysis
- Image processing
- Heavy ion collisions
- ...



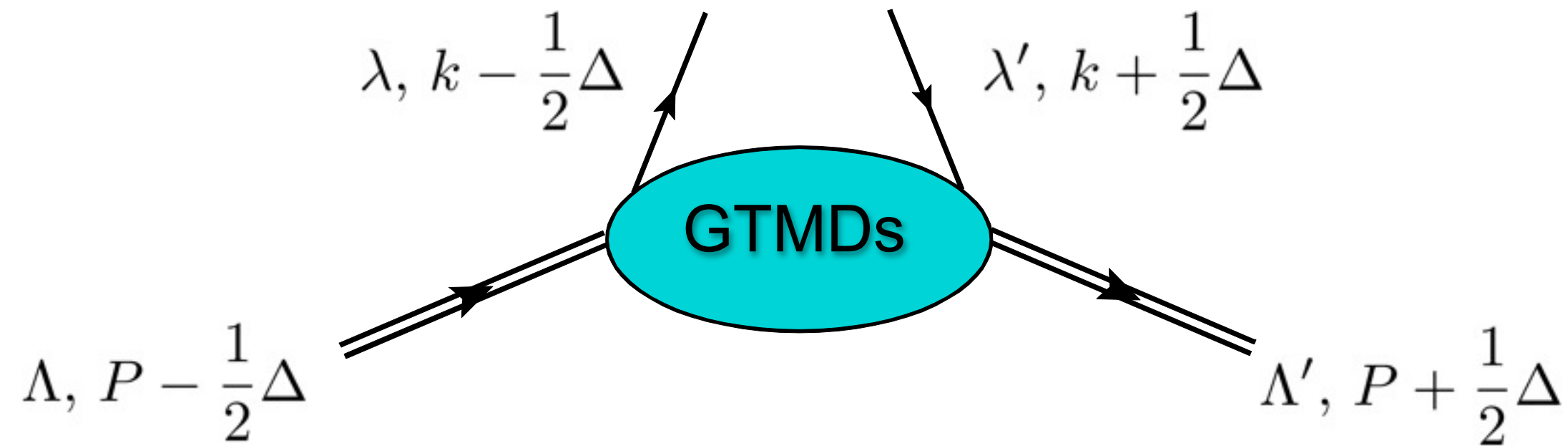
[Antonov et al. (1980-1989)]

Heisenberg's uncertainty relation



Generalized TMDs and Wigner Distributions

[Meißner, Metz, Schlegel (2009)]



Quark polarization



$$W_{\Lambda'\Lambda}^{[\Gamma]} = \frac{1}{2} \int \frac{dz^- d^2 z_\perp}{(2\pi)^3} \langle P', \Lambda' | \bar{\psi}_{\lambda'}^q(-\frac{z}{2}) \Gamma \mathcal{W}(-\frac{z}{2}, \frac{z}{2} | n) \psi_\lambda^q(\frac{z}{2}) | P, \Lambda \rangle e^{i(xP^+ z^- - \vec{k}_\perp \cdot \vec{z}_\perp)}$$



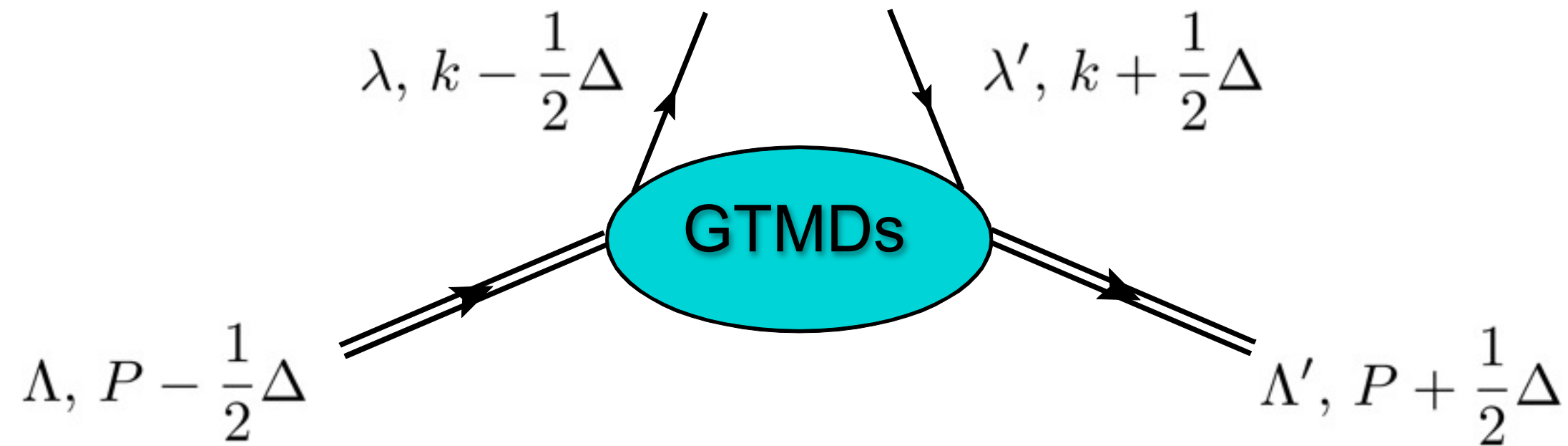
Nucleon polarization

4 X 4 = 16 polarizations \longleftrightarrow 16 complex GTMDs (at twist-2)

$$W_{\Lambda',\Lambda}^\Gamma(x, \xi, \vec{k}_\perp, \vec{\Delta}_\perp)$$

Generalized TMDs and Wigner Distributions

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x: average fraction of quark longitudinal momentum

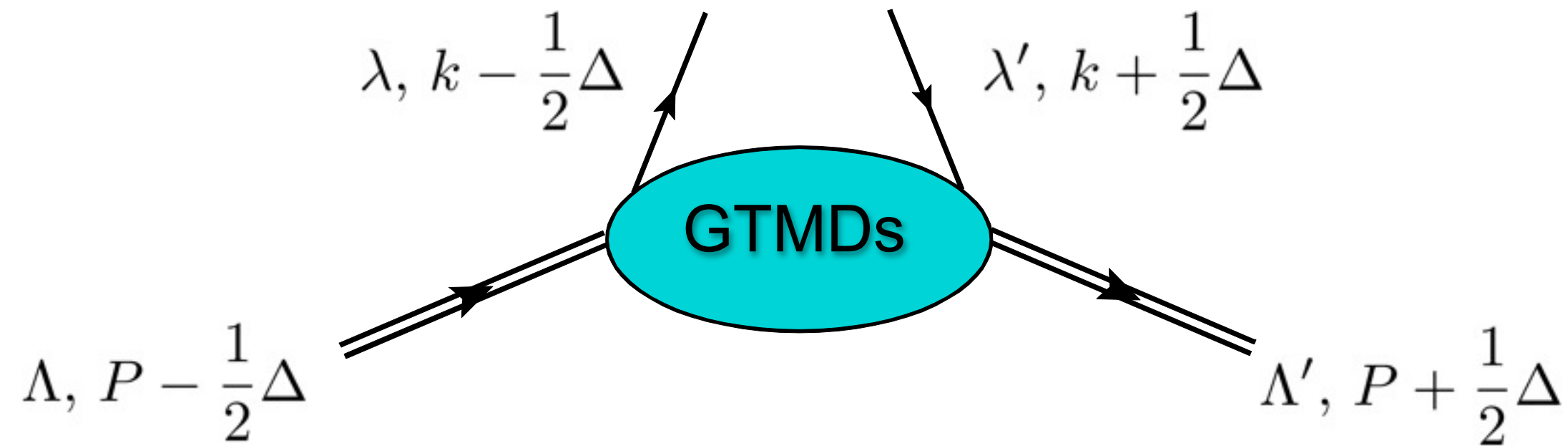
ξ: fraction of longitudinal momentum transfer

\vec{k}_\perp : average quark transverse momentum

$\vec{\Delta}_\perp$: nucleon transverse-momentum transfer

Generalized TMDs and Wigner Distributions

[Meißner, Metz, Schlegel (2009)]



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Nucleon polarization

4 X 4 = 16 polarizations \longleftrightarrow 16 complex GTMDs (at twist-2)

$$W_{\Lambda',\Lambda}^\Gamma(x, \xi, \vec{k}_\perp, \vec{\Delta}_\perp)$$

Fourier transform $\vec{\Delta}_\perp \leftrightarrow \vec{b}_\perp$

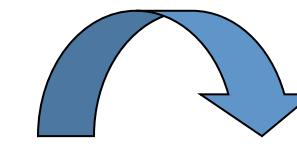
$$\tilde{W}_{\Lambda',\Lambda}^\Gamma(x, \xi, \vec{k}_\perp, \vec{b}_\perp) \quad 16 \text{ real Wigner distributions}$$

GTMDs

$x, \xi, \vec{k}_\perp, \vec{\Delta}_\perp$



2D Fourier
transform

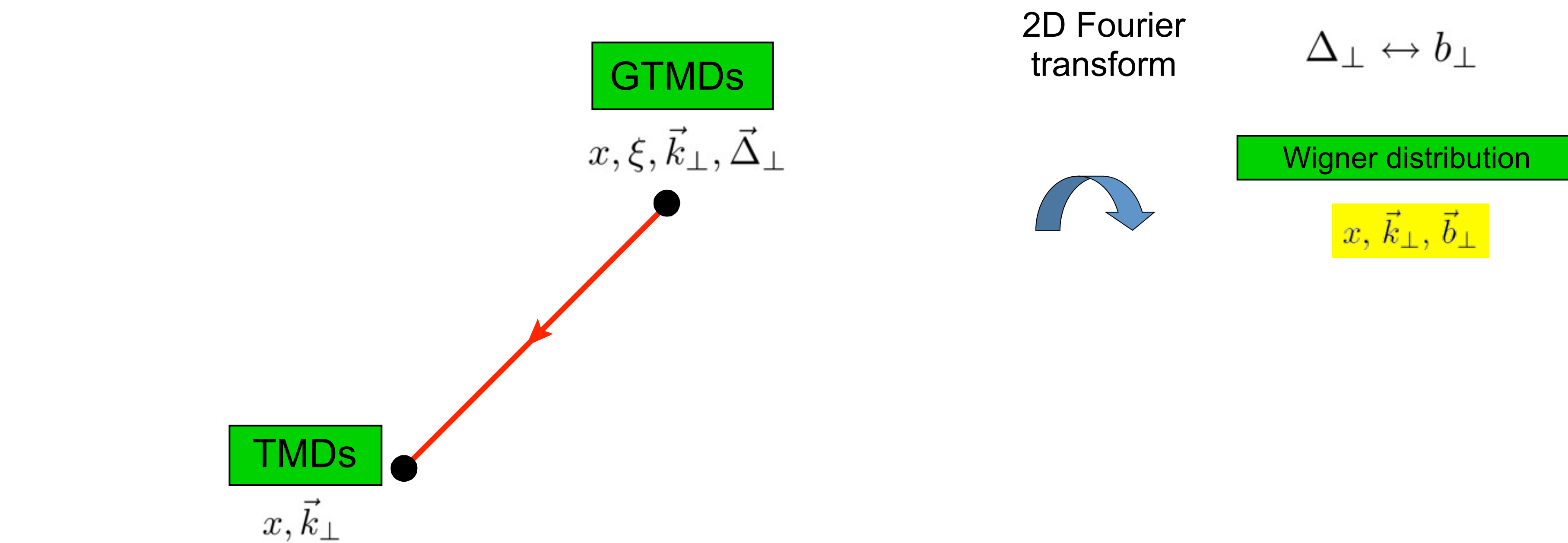


$\Delta_\perp \leftrightarrow b_\perp$

Wigner distribution

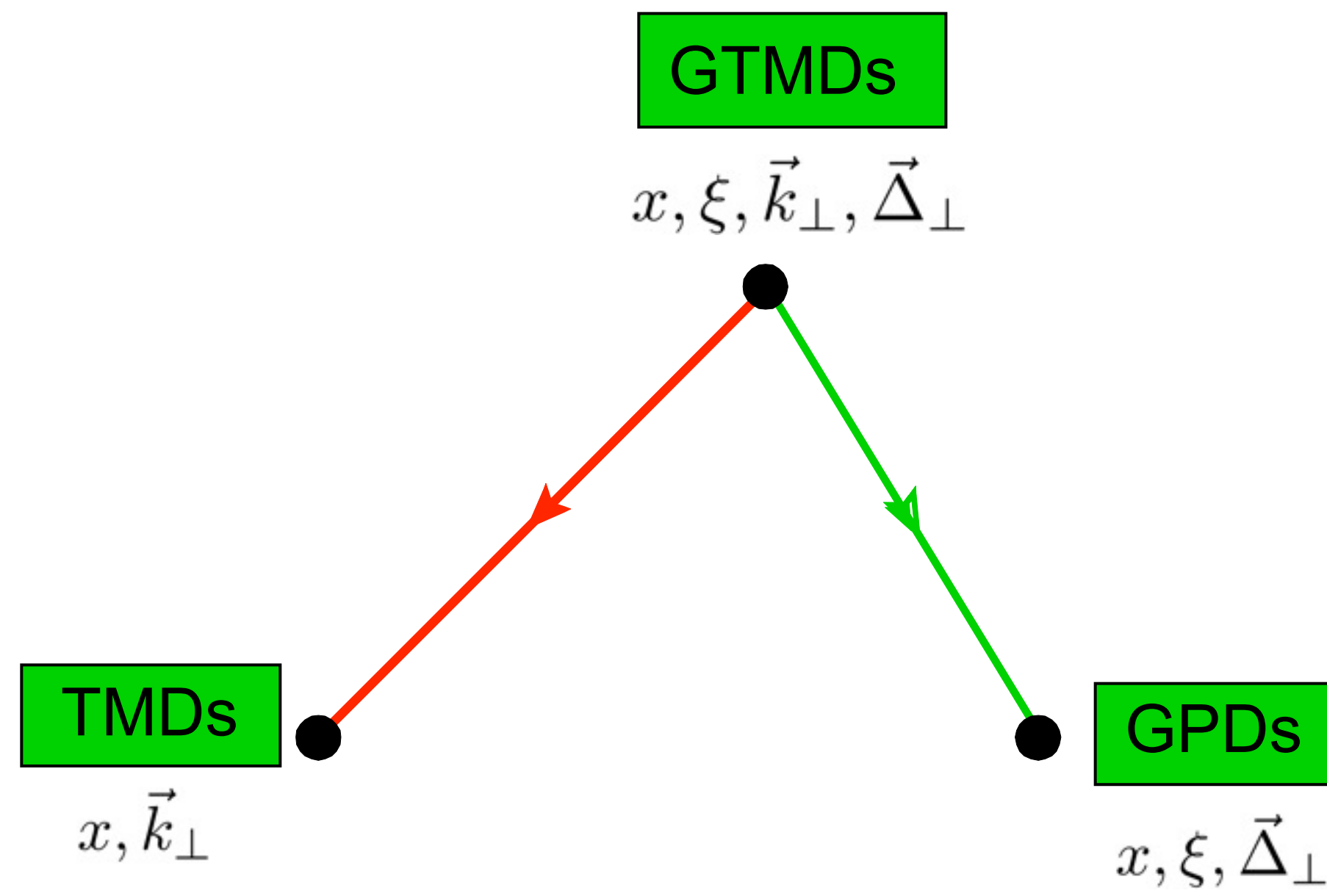
$x, \vec{k}_\perp, \vec{b}_\perp$

[Lorce, BP, Vanderhaeghen, JHEP05 (2011)]

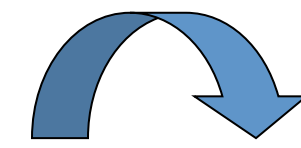


$\rightarrow \vec{\Delta} = 0$

[Lorce, BP, Vanderhaeghen, JHEP05 (2011)]



2D Fourier transform



$$\Delta_\perp \leftrightarrow b_\perp$$

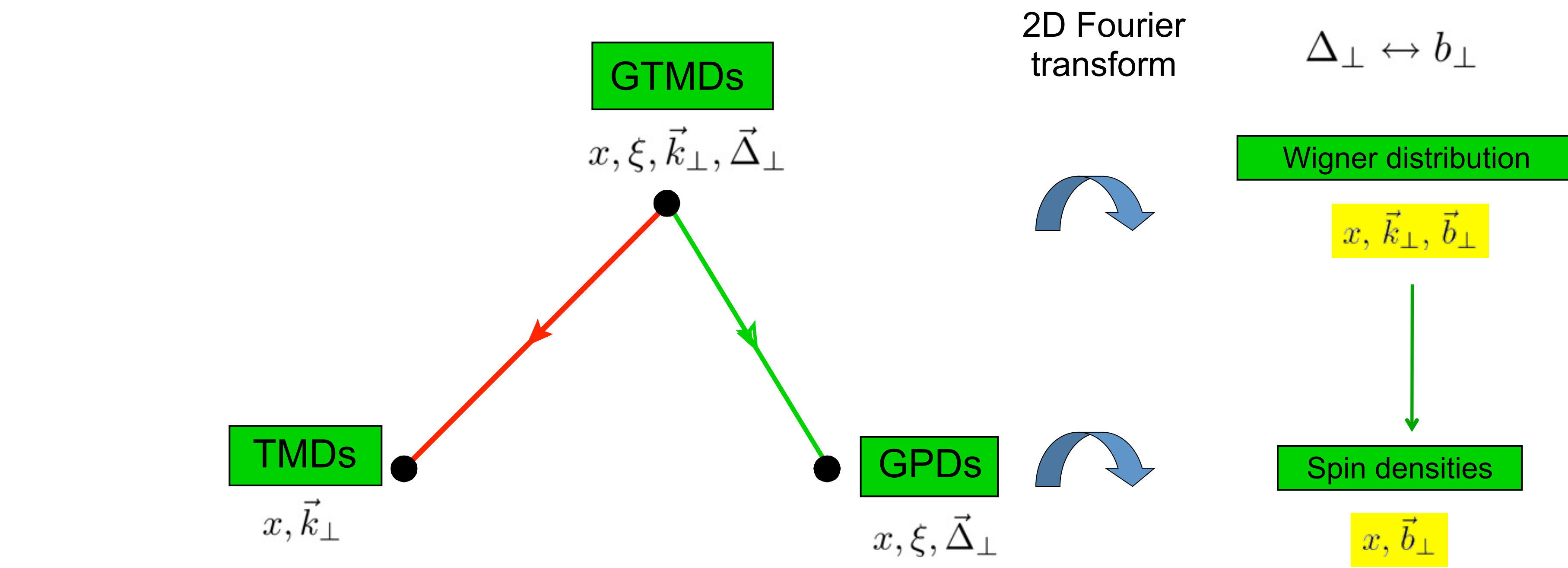
Wigner distribution



$$x, \vec{k}_\perp, \vec{b}_\perp$$

→ $\vec{\Delta} = 0$

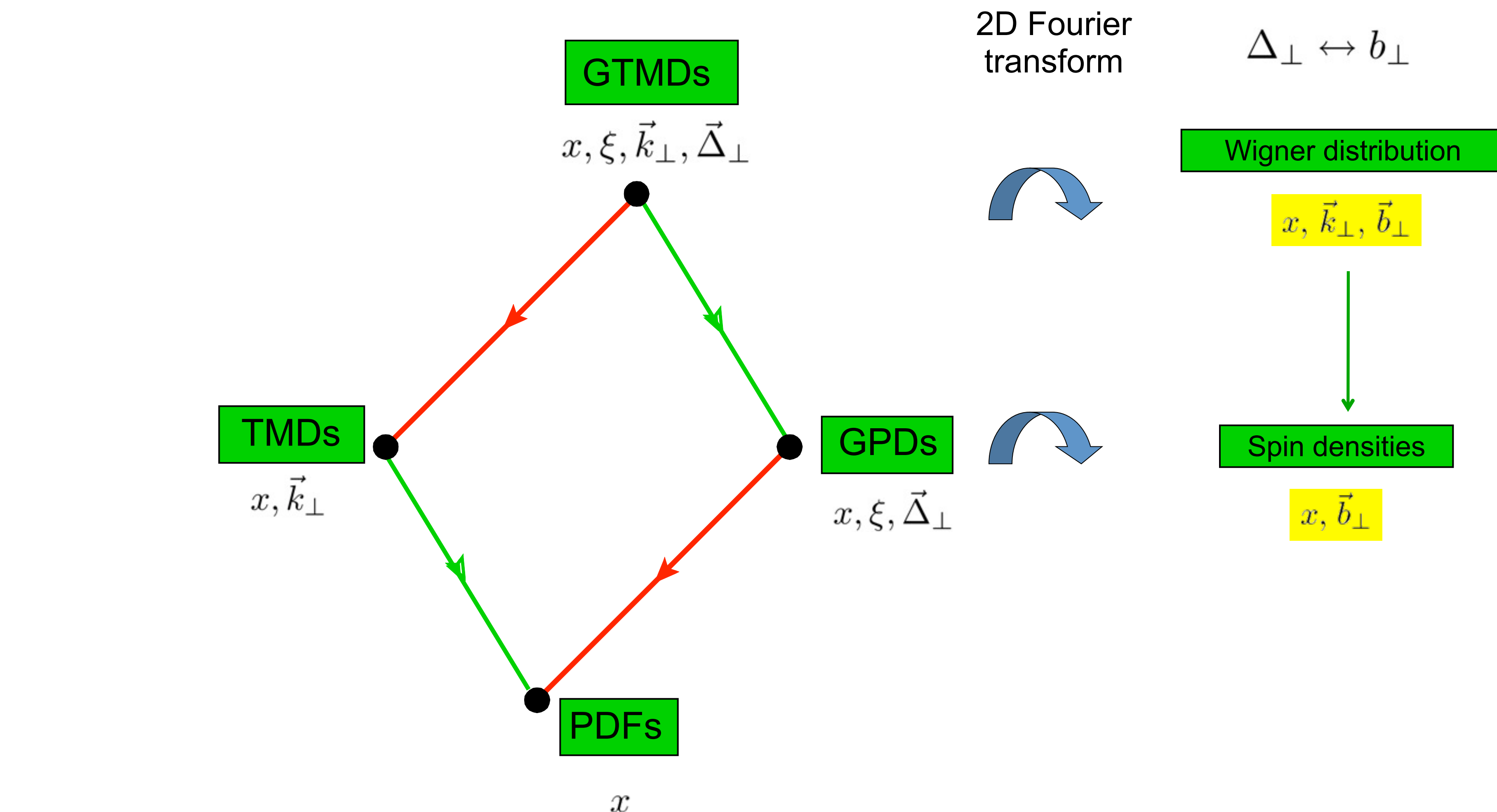
→ $\int dk_\perp$



[Lorce, BP, Vanderhaeghen, JHEP05 (2011)]



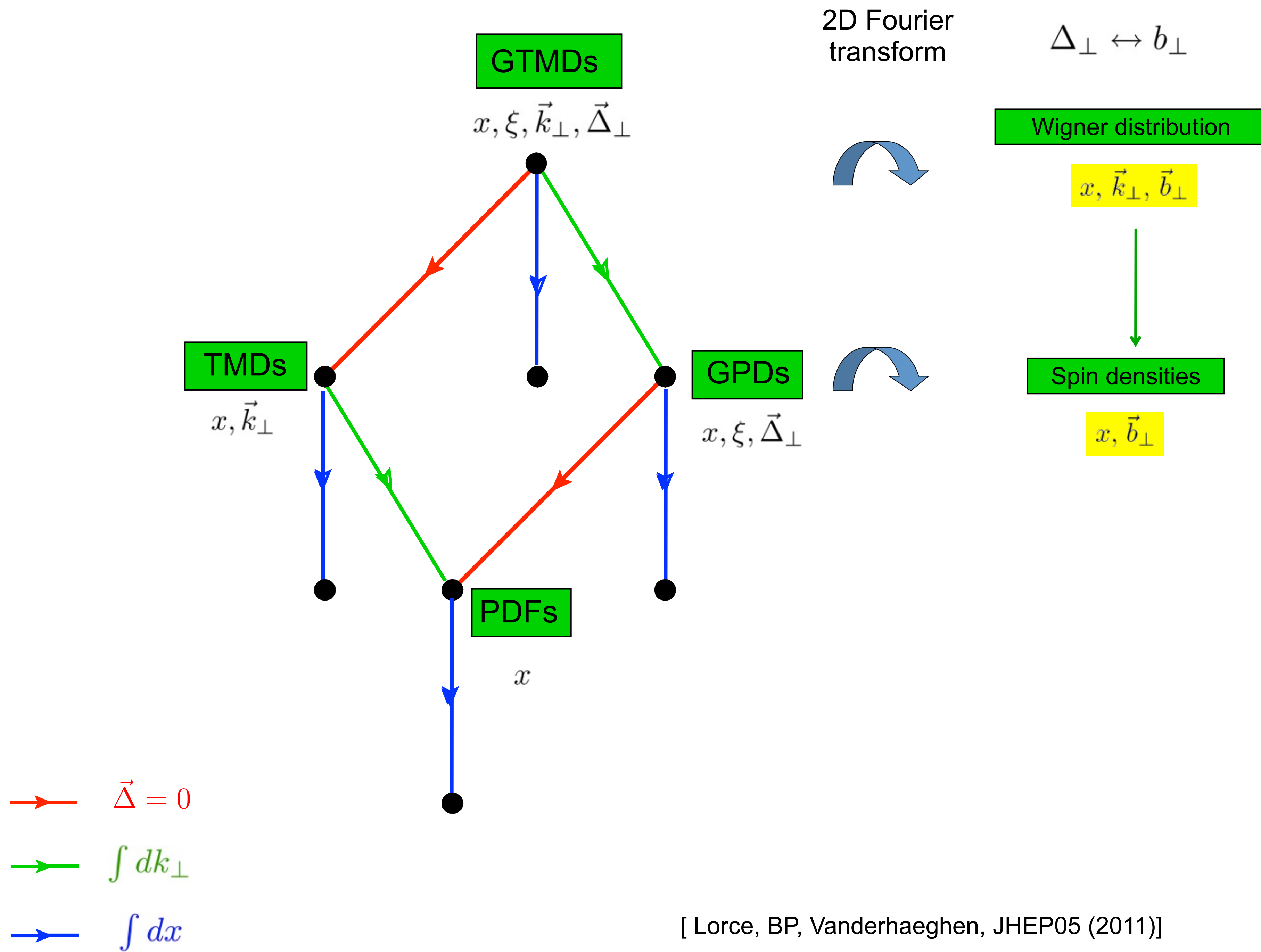
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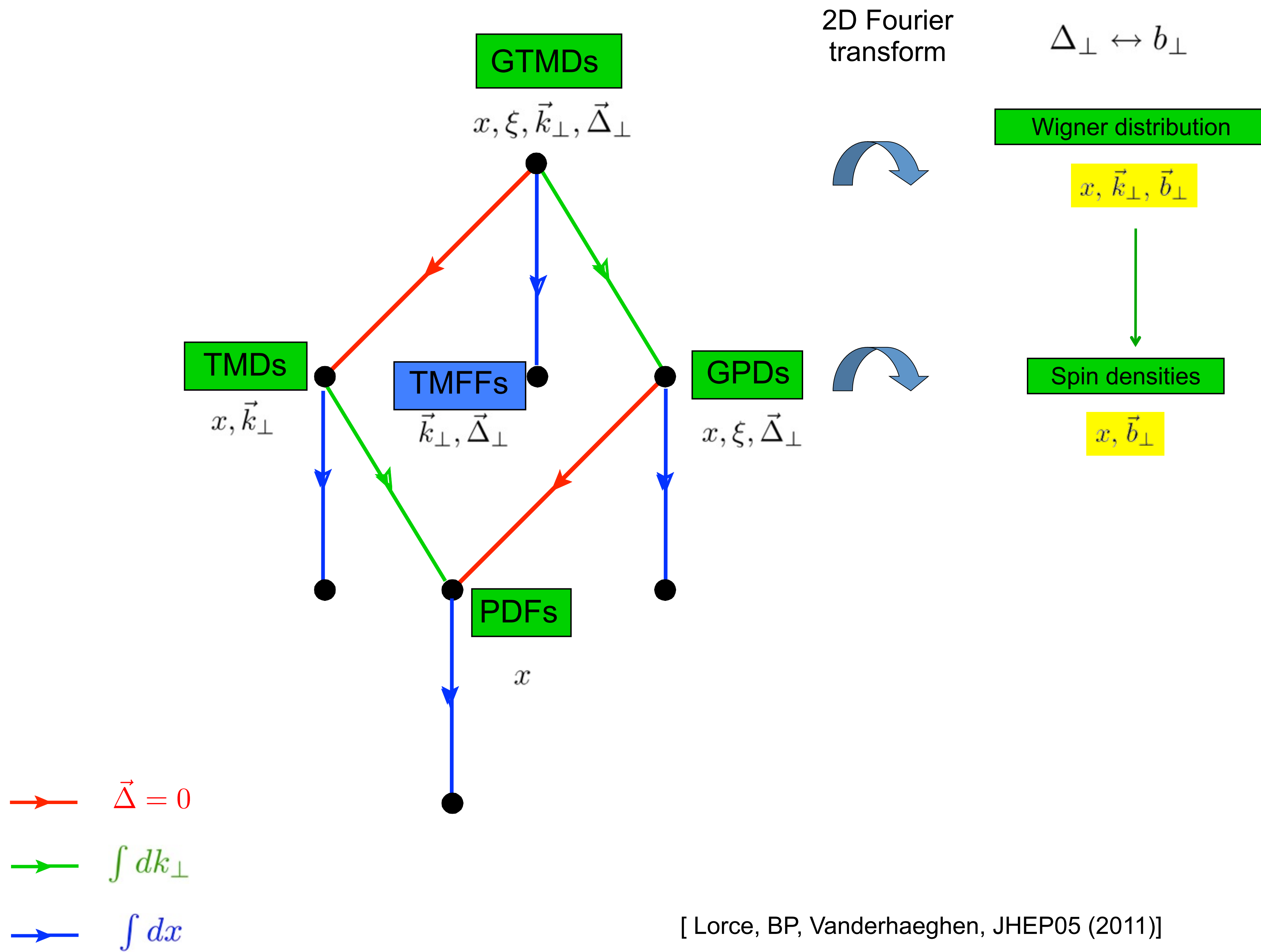


 $\vec{\Delta} = 0$
 $\int dk_\perp$

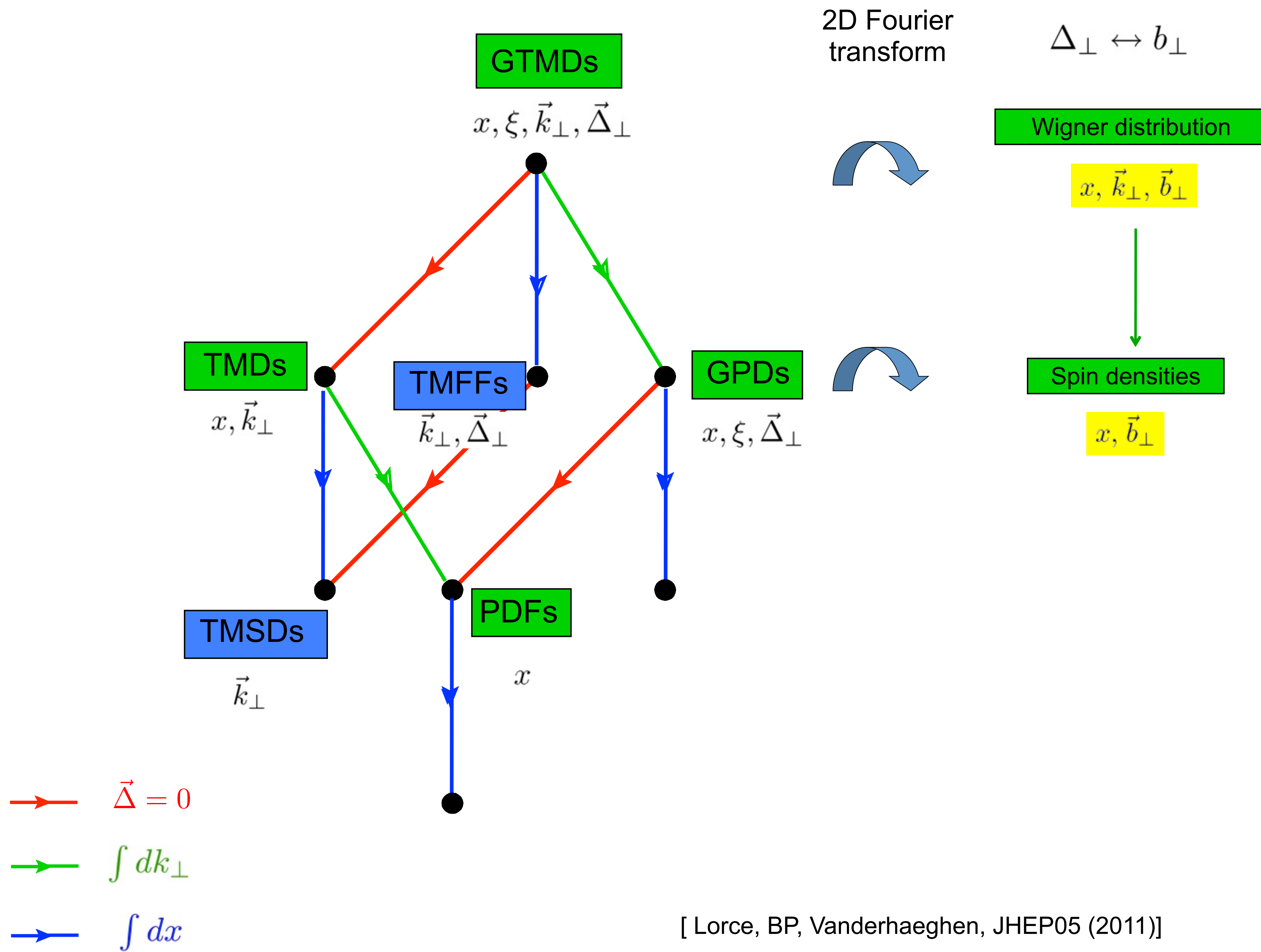
[Lorce, BP, Vanderhaeghen, JHEP05 (2011)]



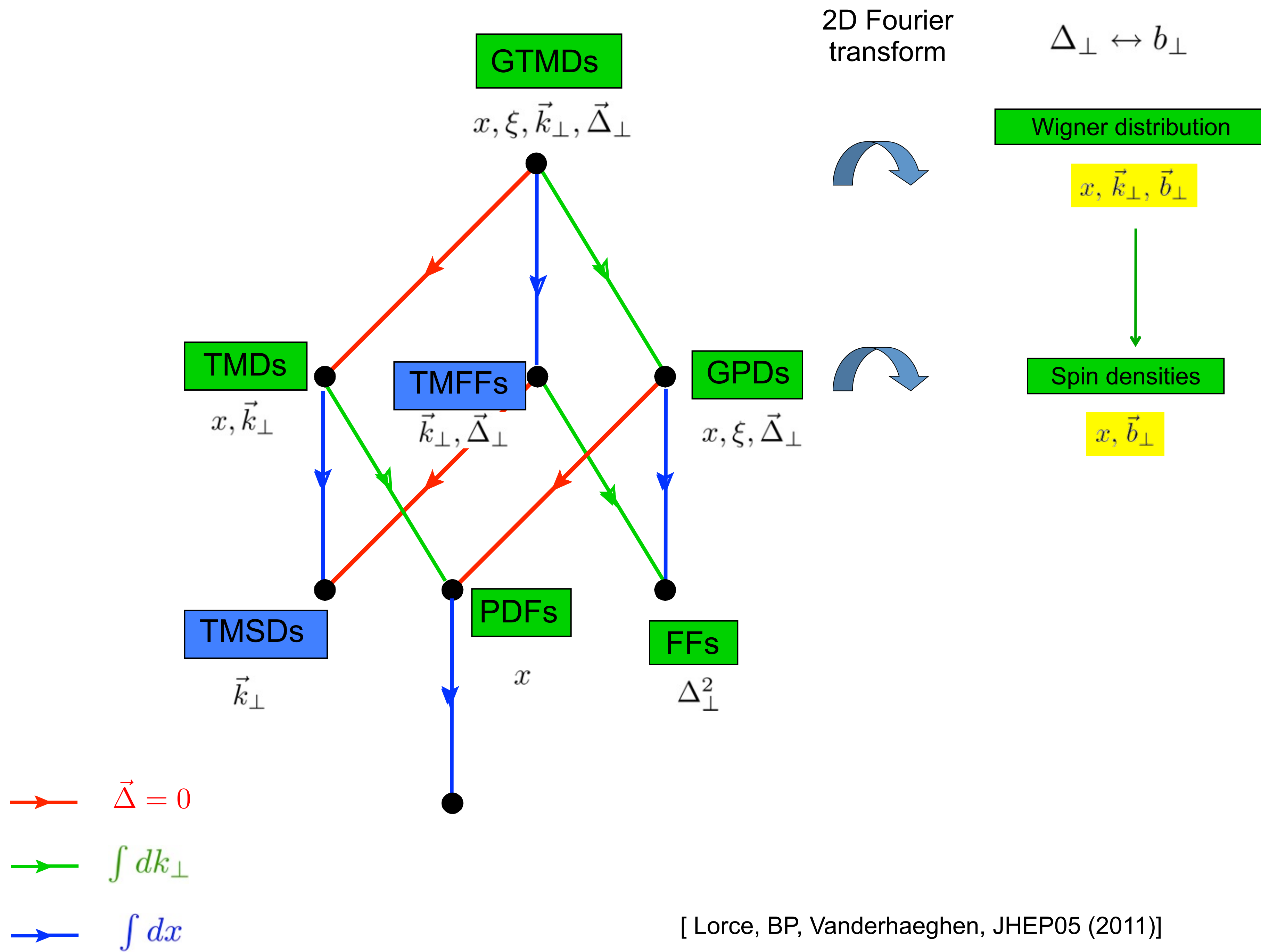
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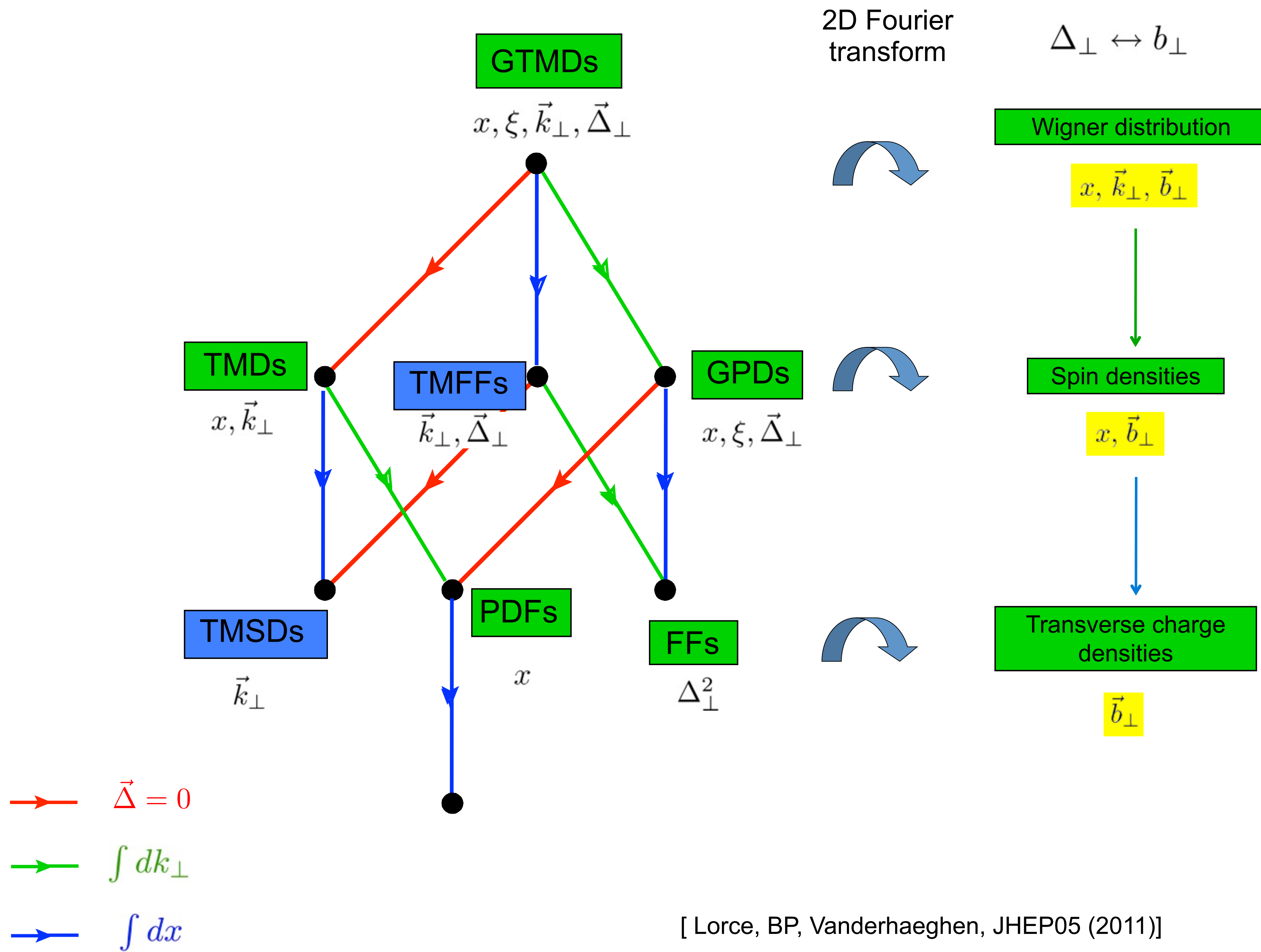
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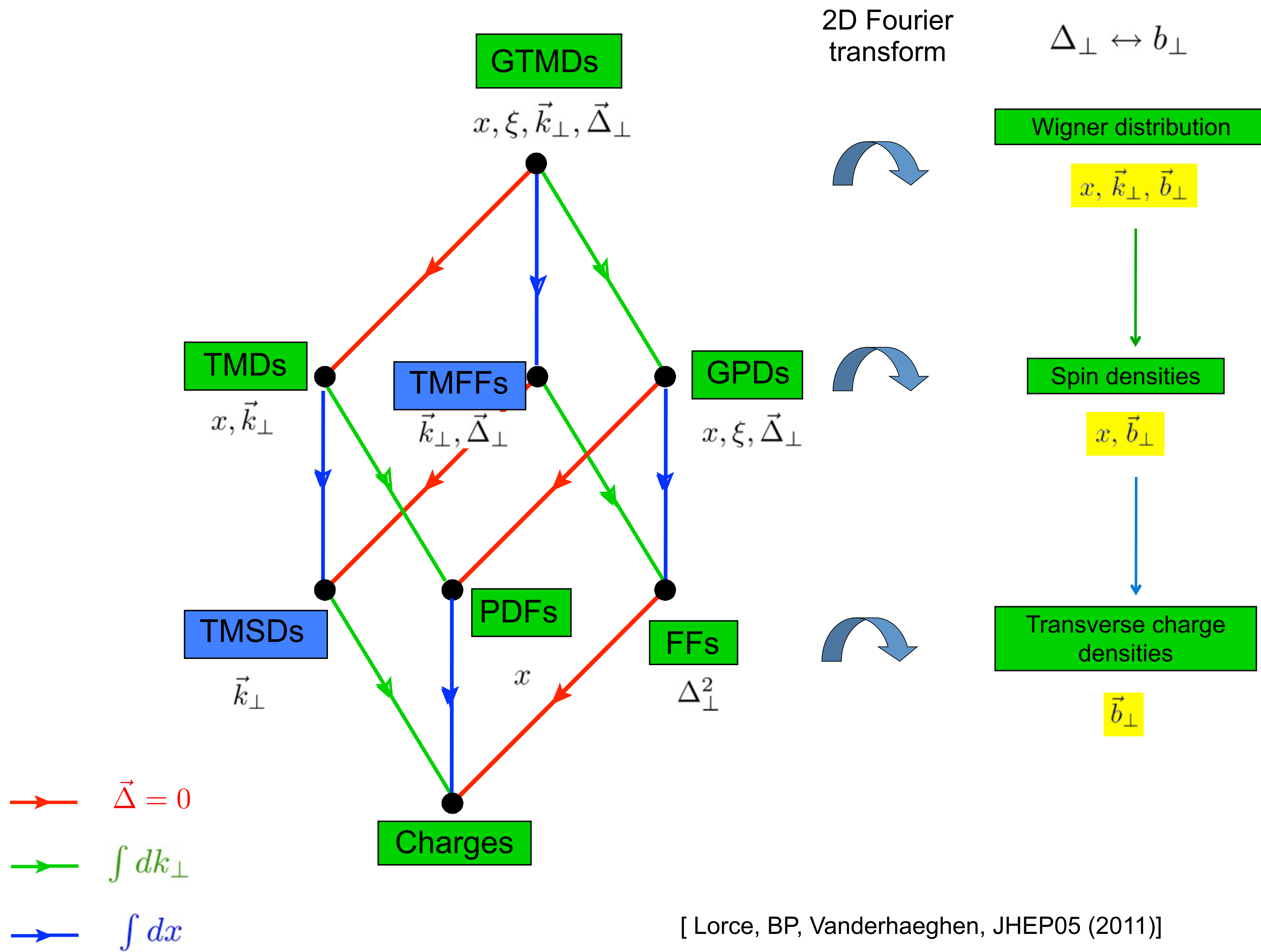
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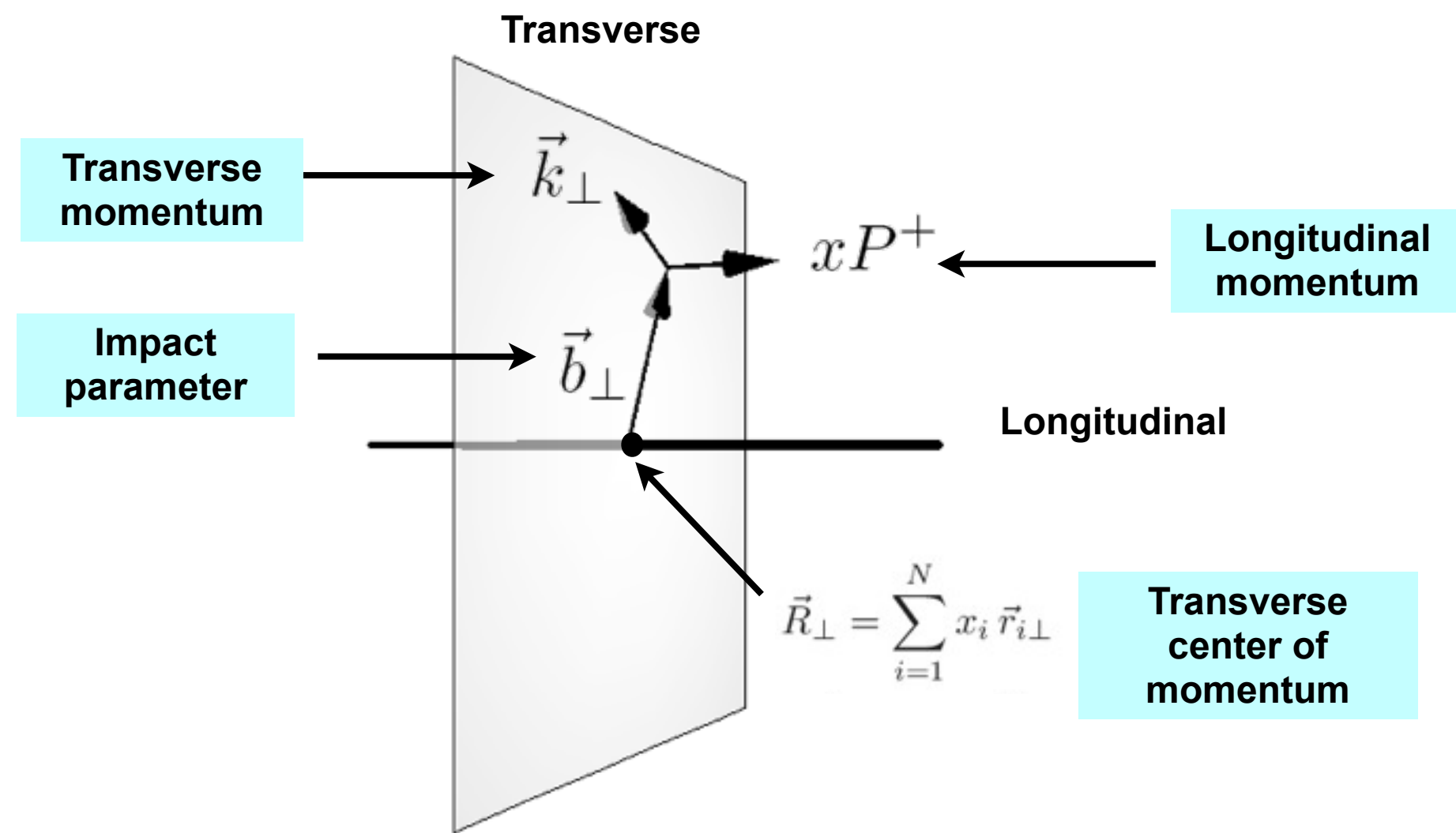


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Wigner Distributions

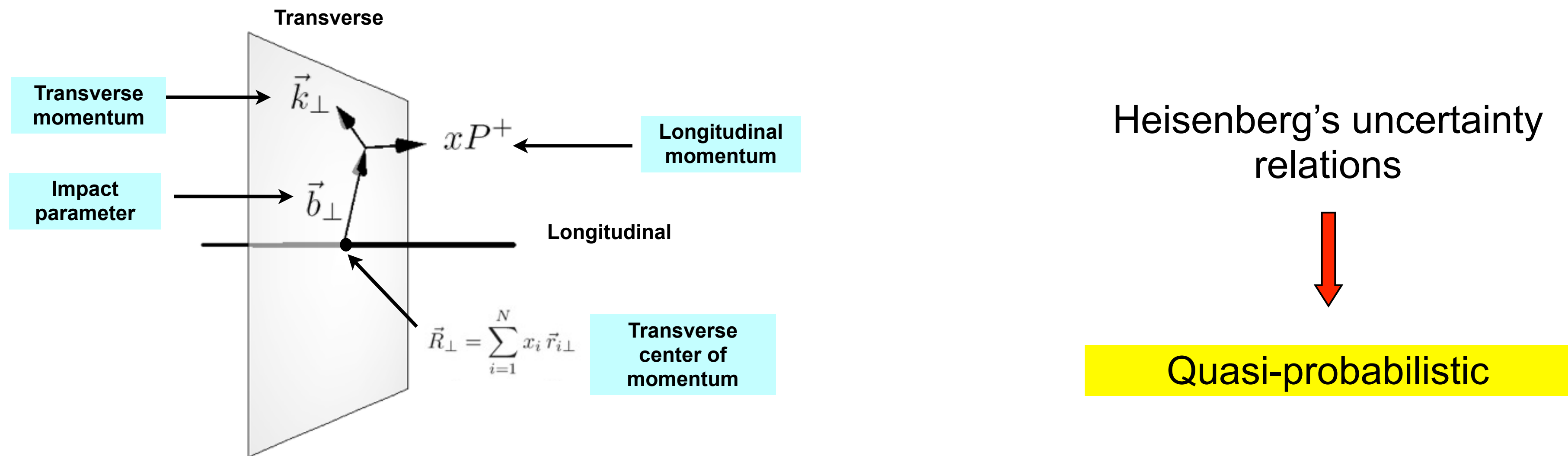


Heisenberg's uncertainty relations



Quasi-probabilistic

Wigner Distributions



❖ real functions, but in general not-positive definite

➡ correlations of quark momentum and position in the transverse plane as function of quark and nucleon polarizations

❖ quantum-mechanical analogous of classical density in the phase space

❖ not directly measurable in experiments

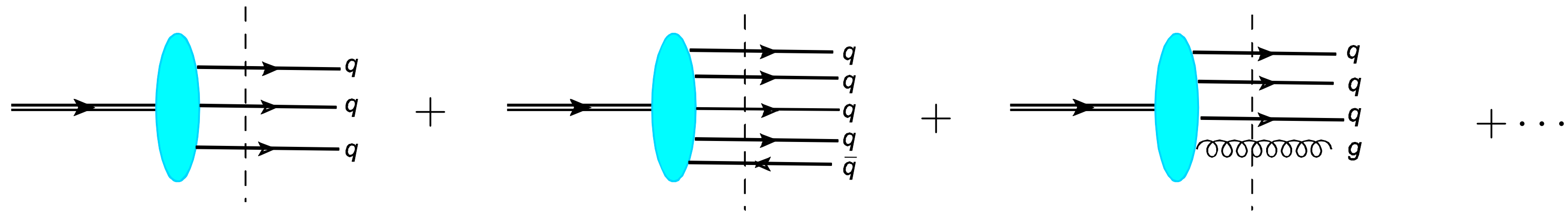
➡ needs phenomenological models with input from experiments on GPDs and TMDs

Light-Front Wave Function

◆ Fock expansion of Nucleon state:

$$|N\rangle = \Psi_{3q}|qqq\rangle + \Psi_{3q q\bar{q}}|3q q\bar{q}\rangle + \Psi_{3q g}|qqqg\rangle + \dots$$

fixed light-cone time ($x^+=0$)

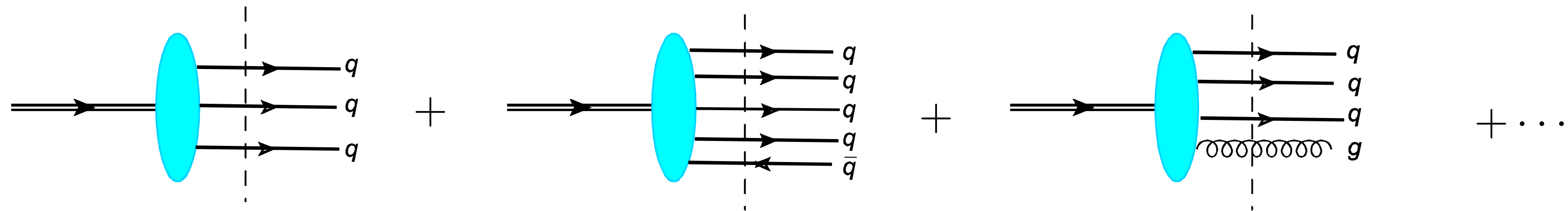


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◆ Eigenstates of momentum

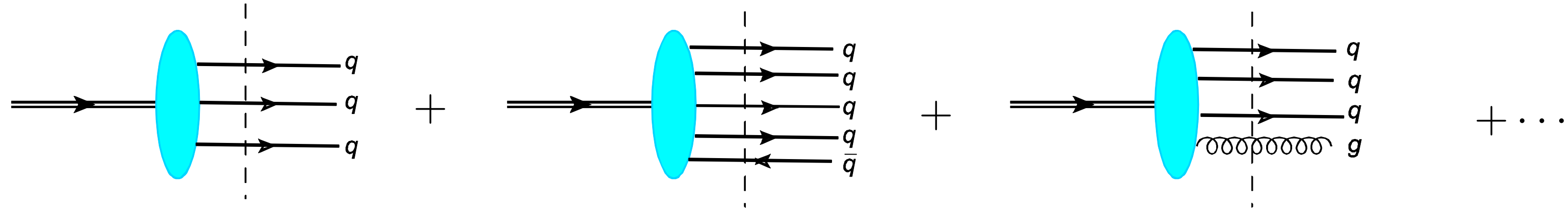
$$P^+ = \sum_{i=1}^N k_i^+ \quad \vec{P}_\perp = \sum_{i=1}^N \vec{k}_{i\perp} = \vec{0}_\perp$$

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◆ Eigenstates of parton light-front helicity

$$\hat{S}_{iz} \Psi_{\lambda_1 \dots \lambda_N}^\Lambda = \lambda_i \Psi_{\lambda_1 \lambda_2 \dots \lambda_N}^\Lambda$$

$$\Lambda = \sum_{i=1}^N \lambda_i + l_z$$

◆ Eigenstates of total orbital angular momentum

$$\hat{L}_z \Psi_{\lambda_1 \dots \lambda_N}^\Lambda = l_z \Psi_{\lambda_1 \lambda_2 \dots \lambda_N}^\Lambda$$

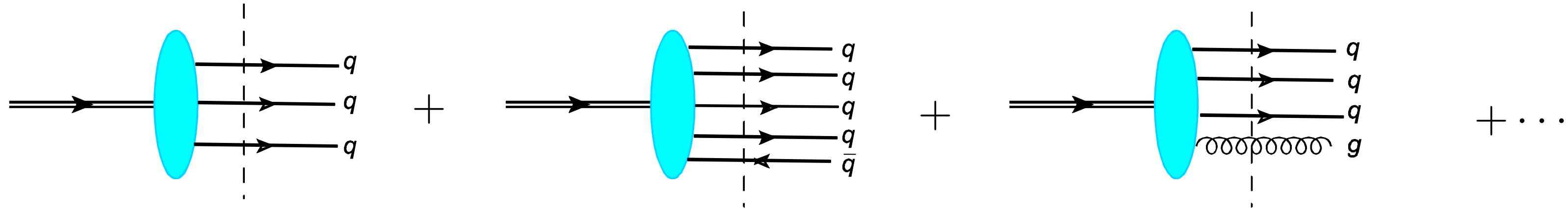
⚠ $A^+ = 0$ gauge

Light-Front Wave Function

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⚠ $A^+ = 0$ gauge

◆ Probability to find N partons in the nucleon

$$\rho_{N,\beta}^\Lambda = \int [dx]_N [d^2 k_\perp]_N |\Psi_{\lambda_1 \dots \lambda_N}^\Lambda|^2$$

normalization $\sum_{N,\beta} \rho_{N,\beta}^\Lambda = 1$

total helicity

$$s_z = \langle \hat{S}_z \rangle = \sum_{N,\beta} \sum_{i=1}^N \lambda_i \rho_{N,\beta}^\Lambda$$

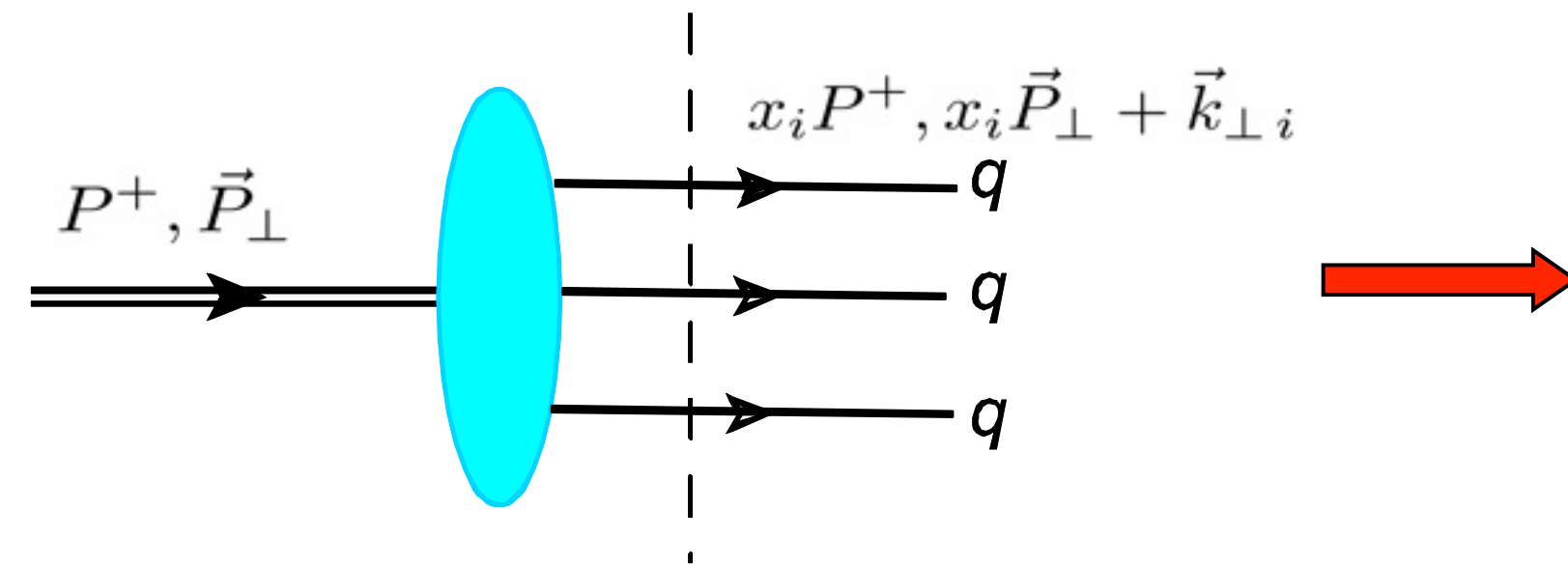
total OAM

$$l_z = \langle \hat{L}_z \rangle = \sum_{N,\beta} \sum_{i=1}^N l_z \rho_{N,\beta}^\Lambda$$

nucleon helicity

$$\Lambda = s_z + l_z$$

LFWF Overlap representation



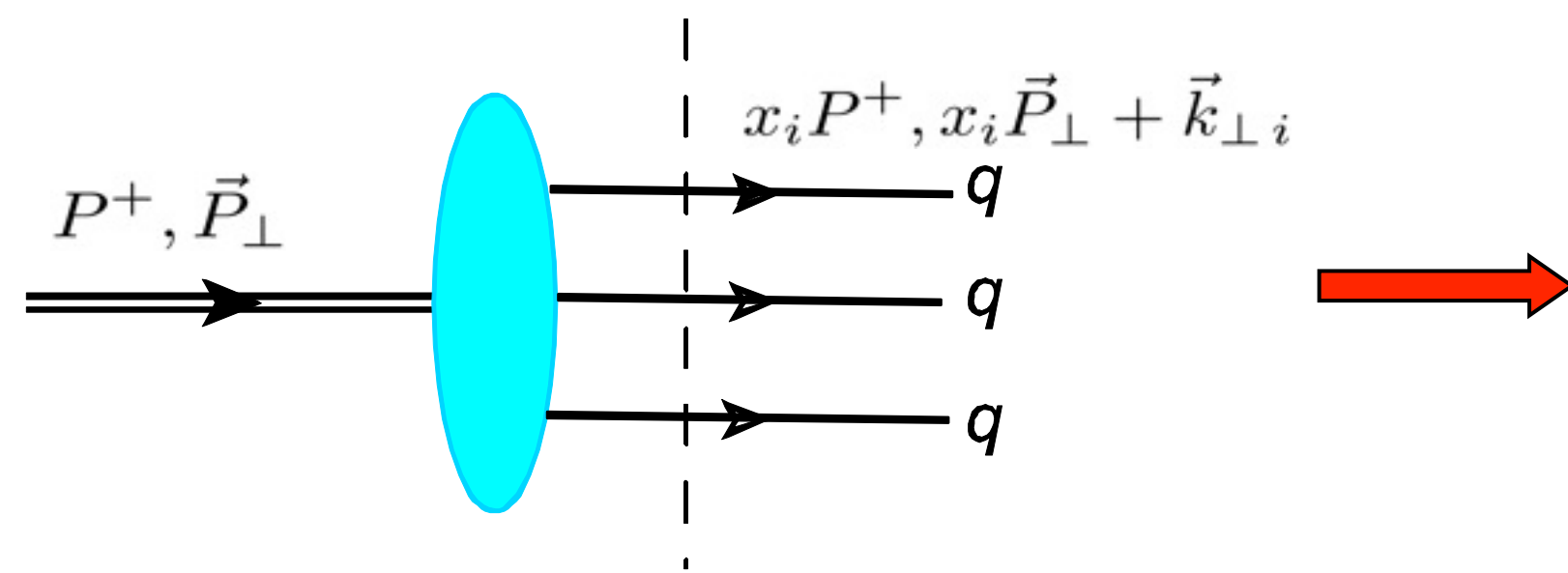
3q LFWF: $\Psi_{\lambda_1 \lambda_2 \lambda_3}^{\Lambda; q_1 q_2 q_3} (x_i, \vec{k}_{\perp,i})$

invariant under boost, independent of P^μ

internal variables: $\sum_{i=1}^3 x_i = 1, \sum_{i=1}^3 \vec{k}_{\perp,i} = \vec{0}_\perp$

[Brodsky, Pauli, Pinsky, '98]

LFWF Overlap representation



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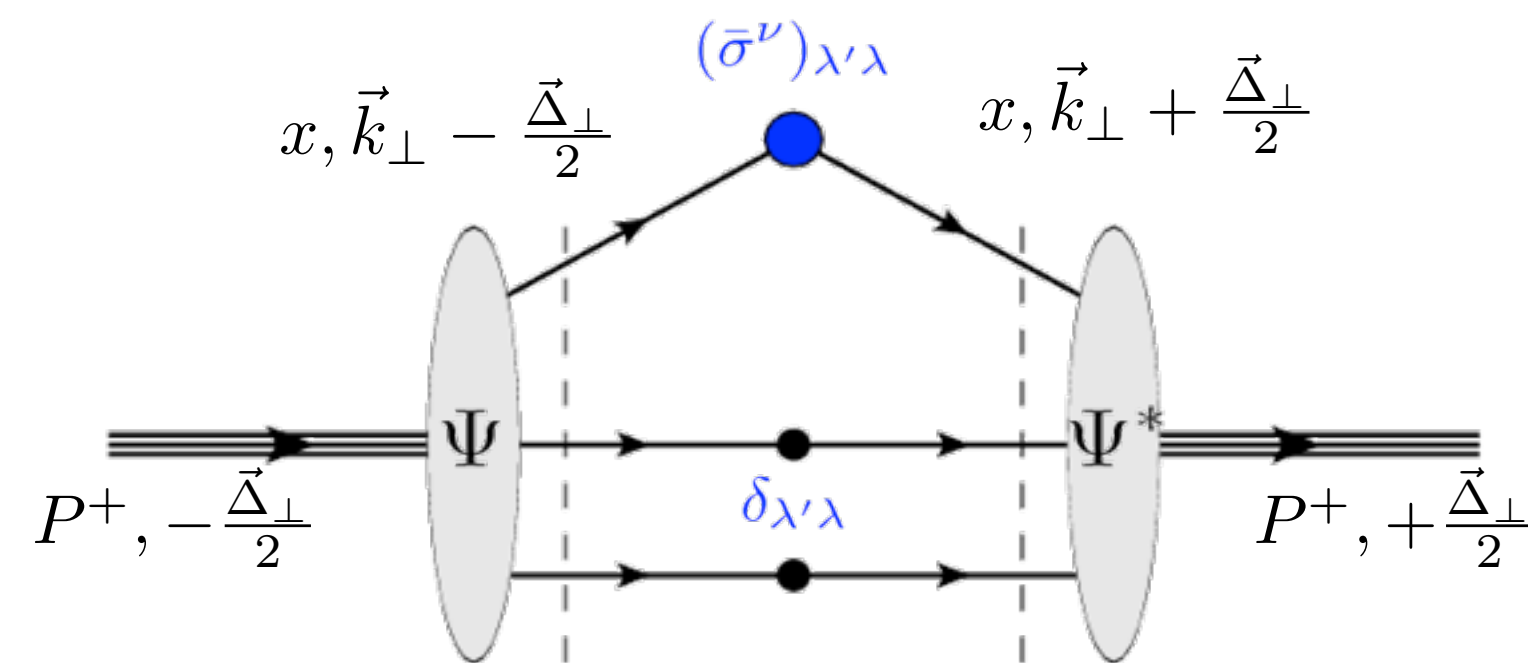
[Brodsky, Pauli, Pinsky, '98]

$A^+=0 \Rightarrow$ Wilson line equal to unit

$\Delta^+=0 \Rightarrow$ diagonal in the Fock-space

quark-quark correlator

$(\Delta^+ = 0)$



General formalism valid for

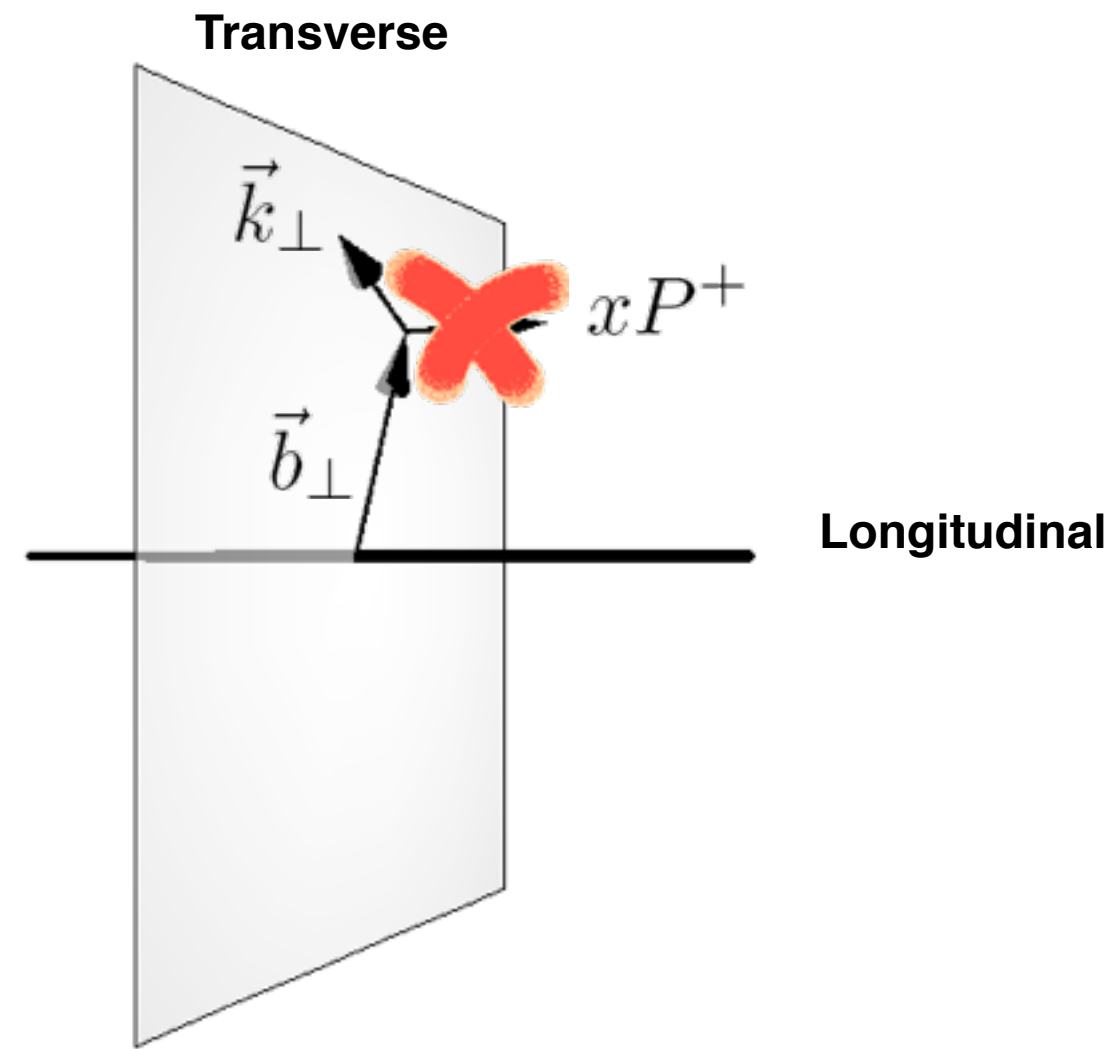
Bag Model, LF χ QSM, LFCQM, Quark-Diquark, Covariant Parton Models

Common assumptions :

- No gluons
- Independent quarks

[Lorce, BP, Vanderhaeghen, JHEP05 (2011)]

Quark Wigner Distributions



$$\rho(\vec{k}_\perp, \vec{b}_\perp) = \int dx \rho(x, \vec{k}_\perp, \vec{b}_\perp)$$

at fixed $\vec{k}_\perp \longrightarrow$ two-dimensional distributions in impact-parameter space

★ Twist-2 ~ LO in P^+

$$\Gamma_{\text{twist-2}} = \gamma^+, \gamma^+ \gamma_5, i\sigma^{j+} \gamma_5$$

quark polarization \longrightarrow U L T

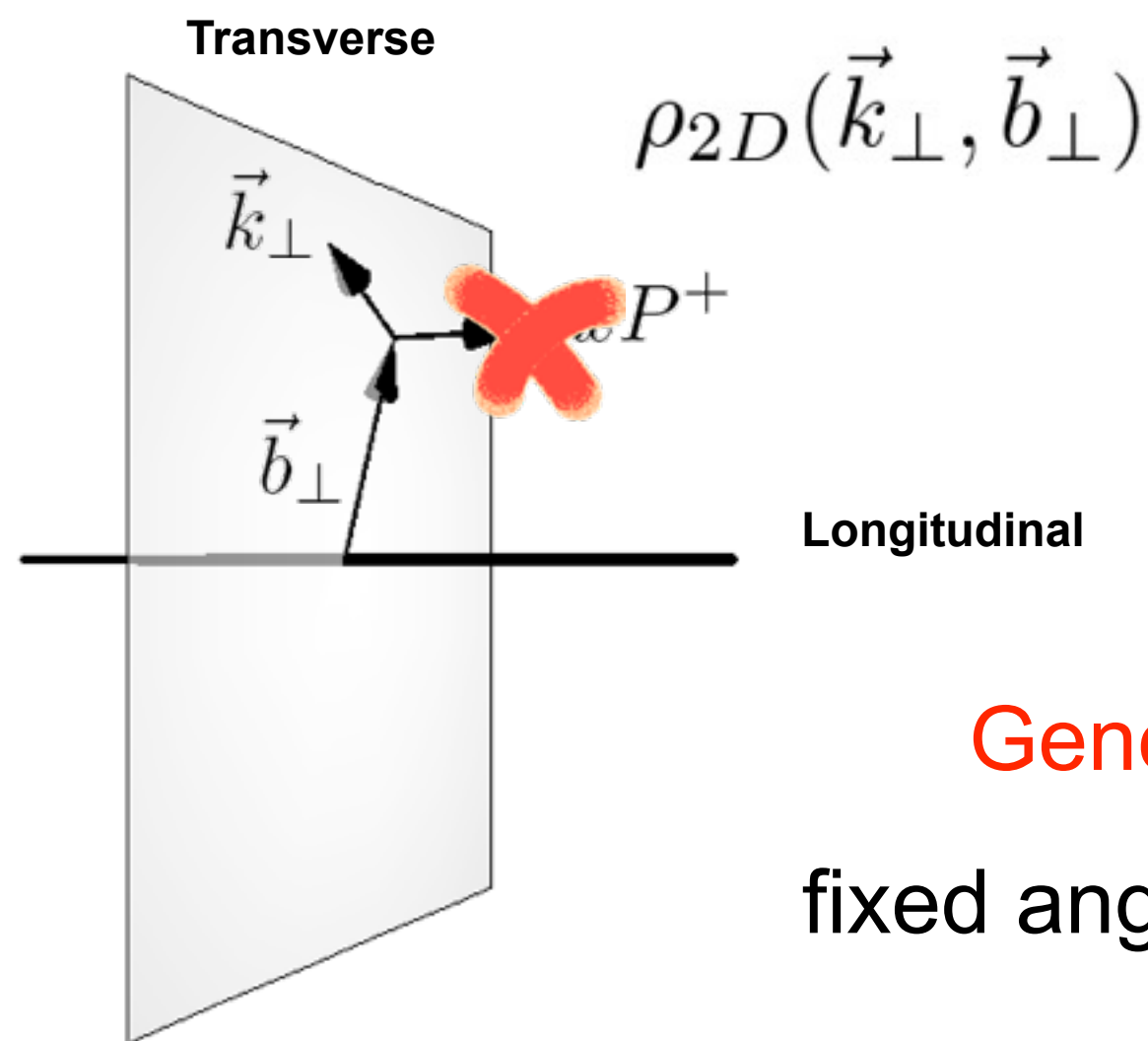
★ Nucleon polarization: U L T



16 independent Wigner distributions

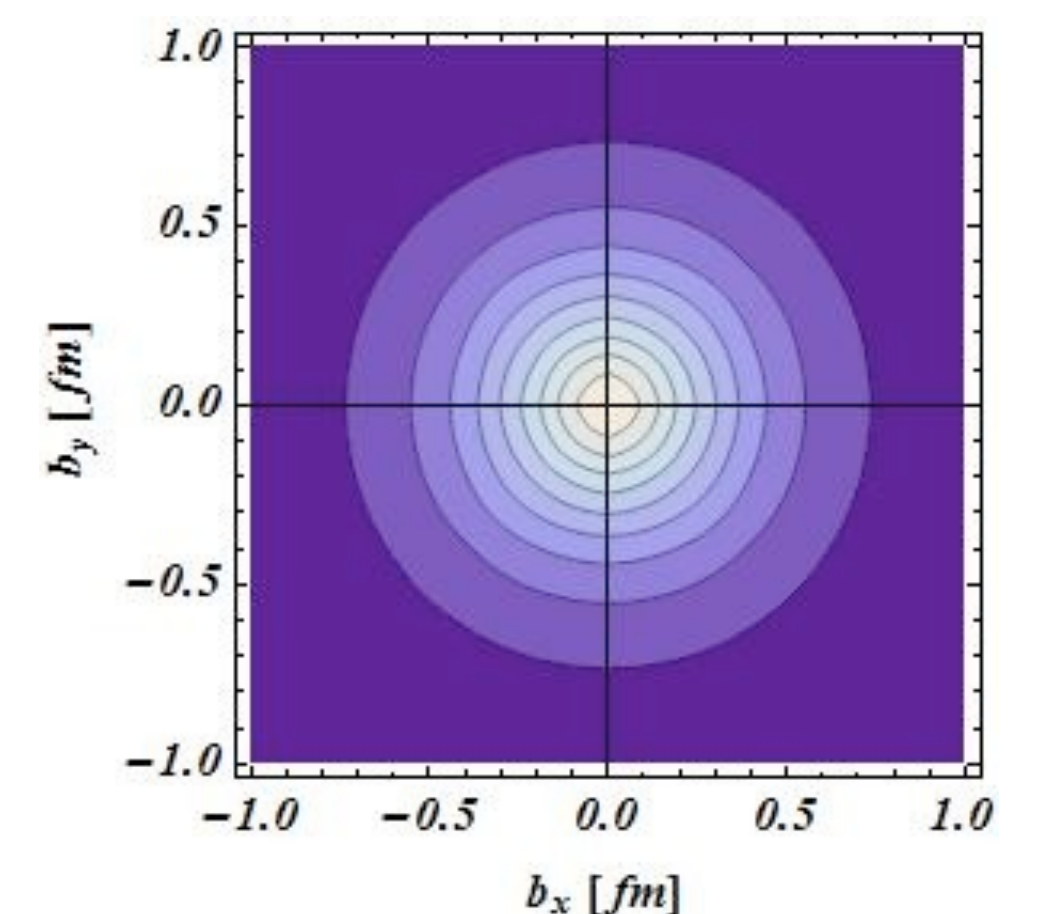
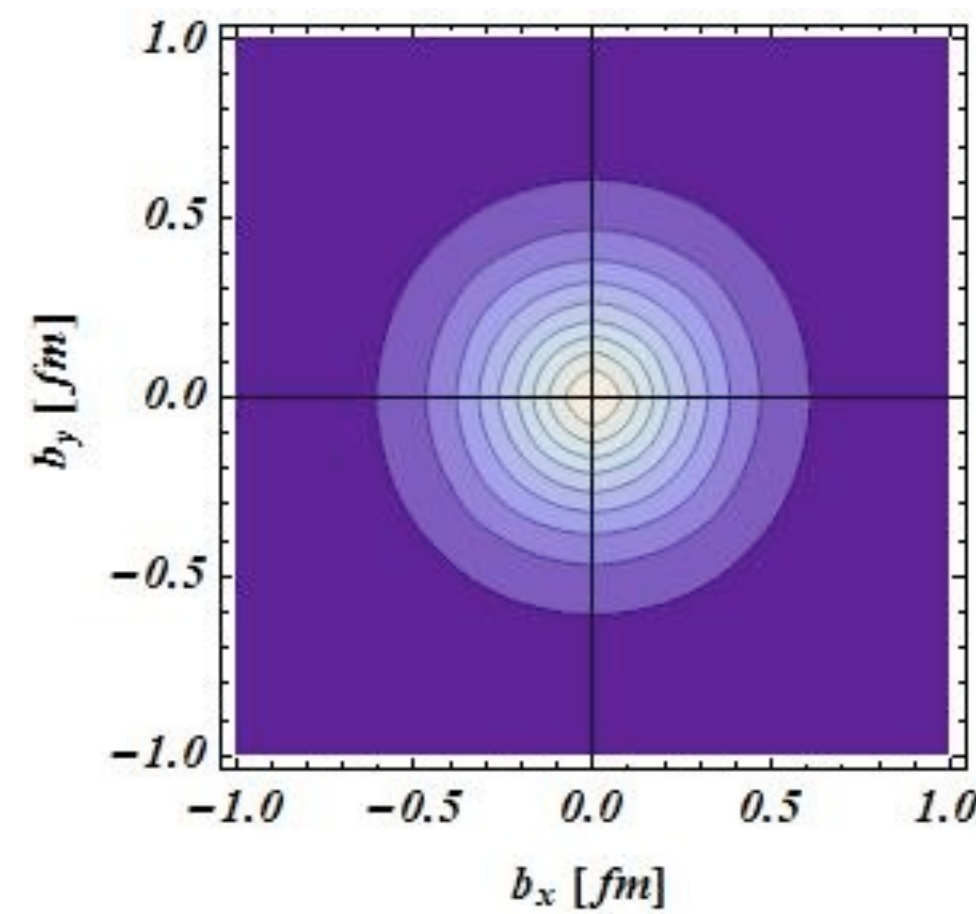
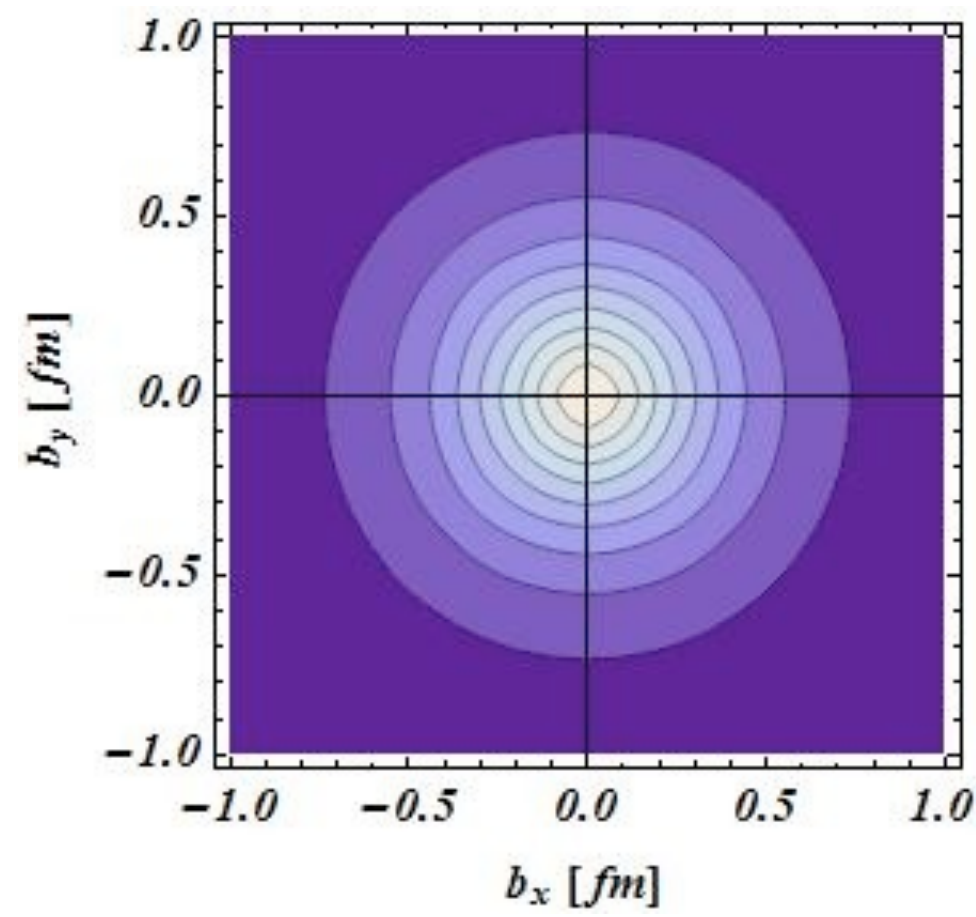
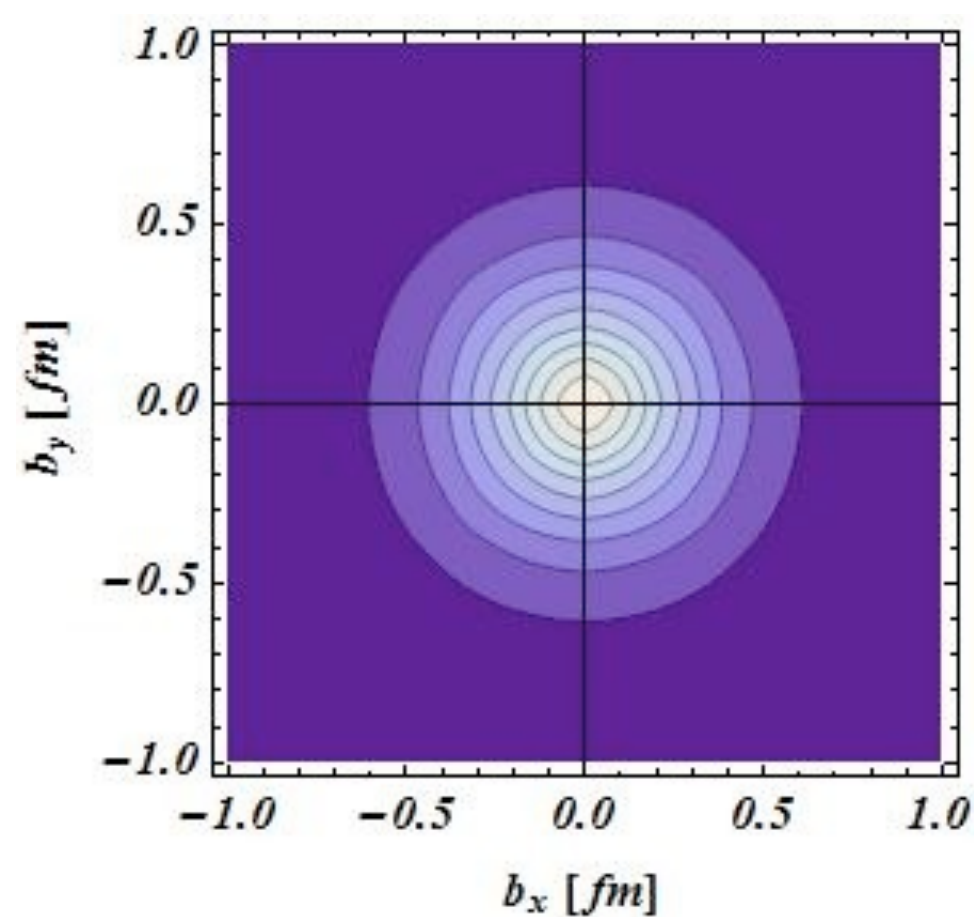
Unpol. up Quark in Unpol. Proton

[Lorce', BP, PRD84 (2011)]



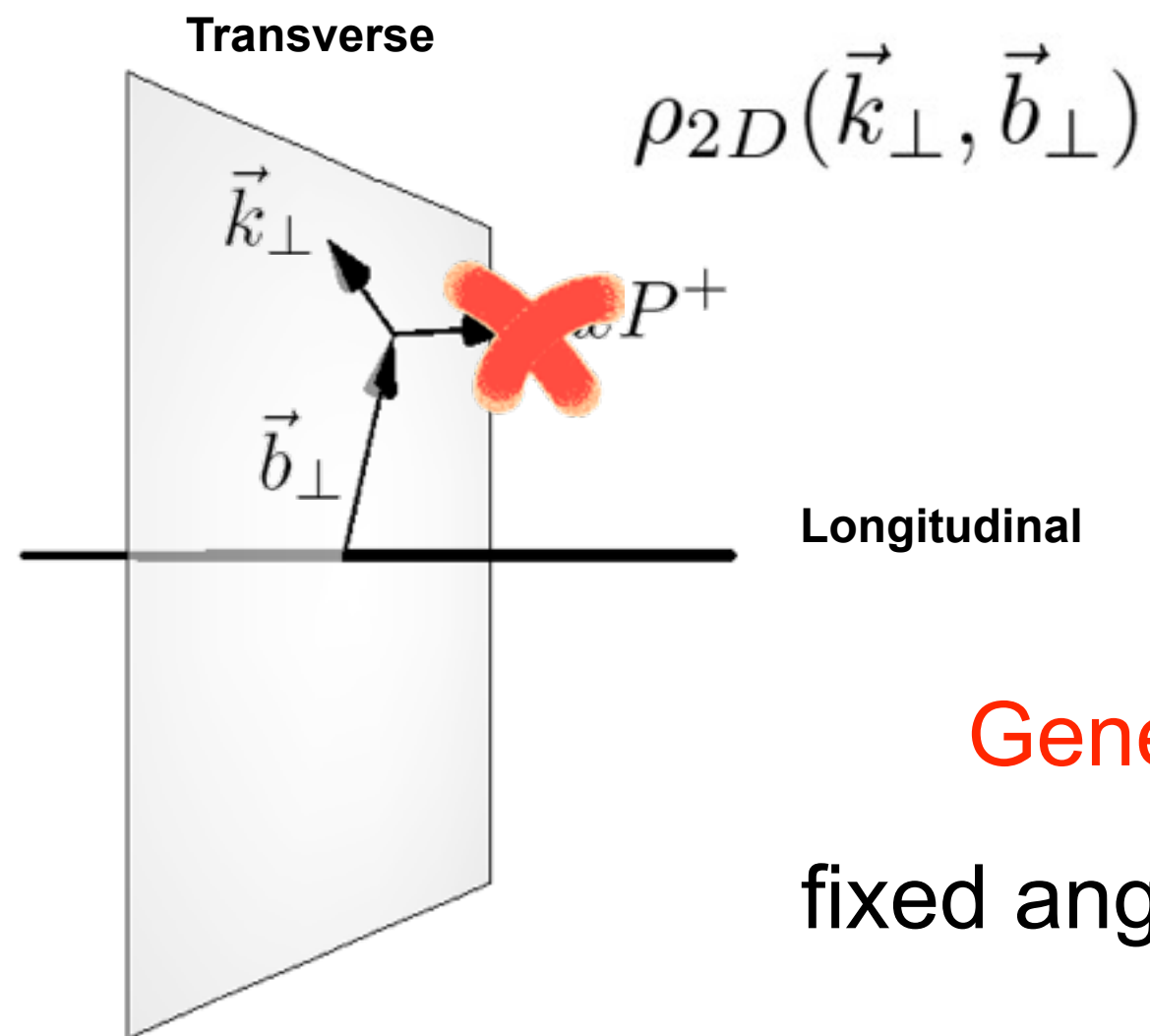
Generalized Transverse Charge Density

fixed angle between \vec{k}_\perp and \vec{b}_\perp and fixed value of $|\vec{k}_\perp|$



Unpol. up Quark in Unpol. Proton

[Lorce', BP, PRD84 (2011)]

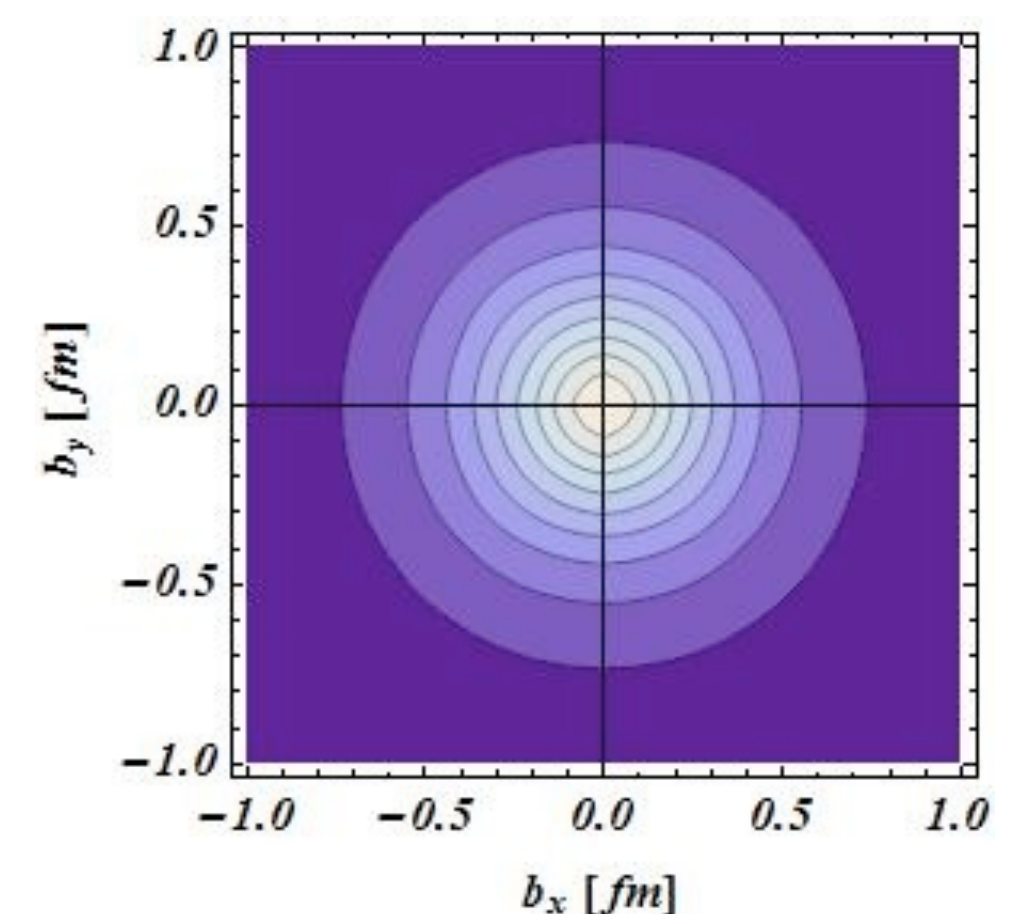
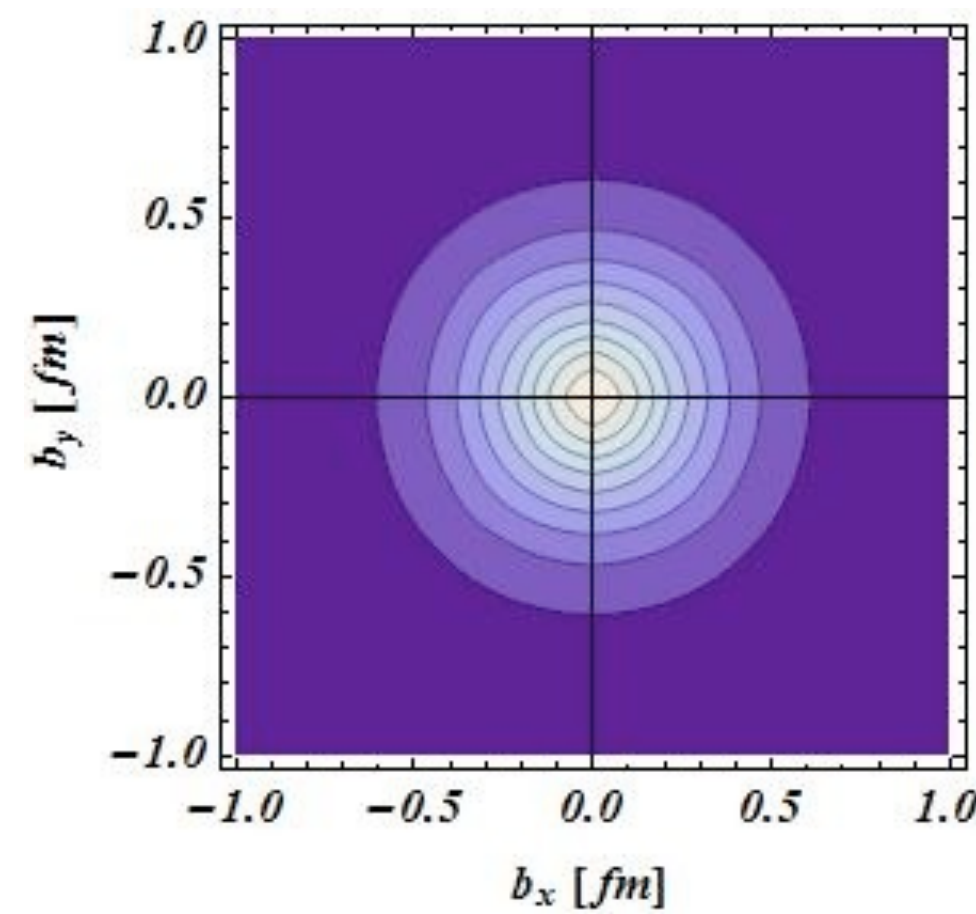
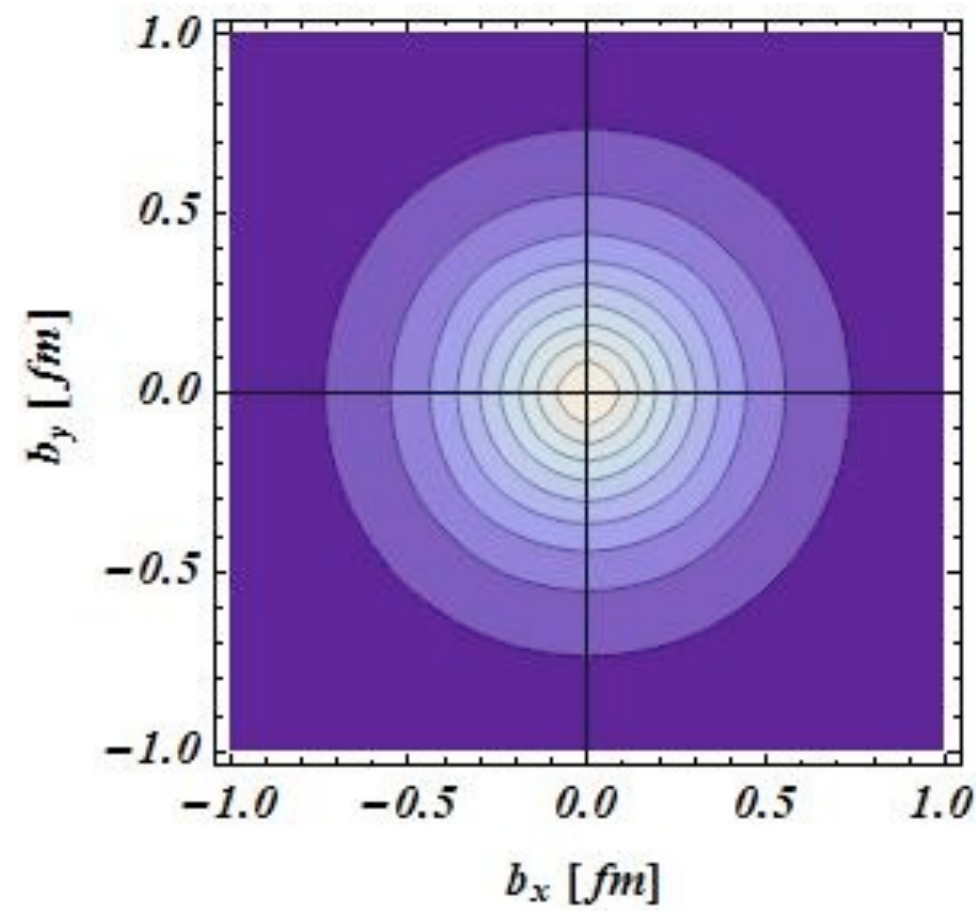
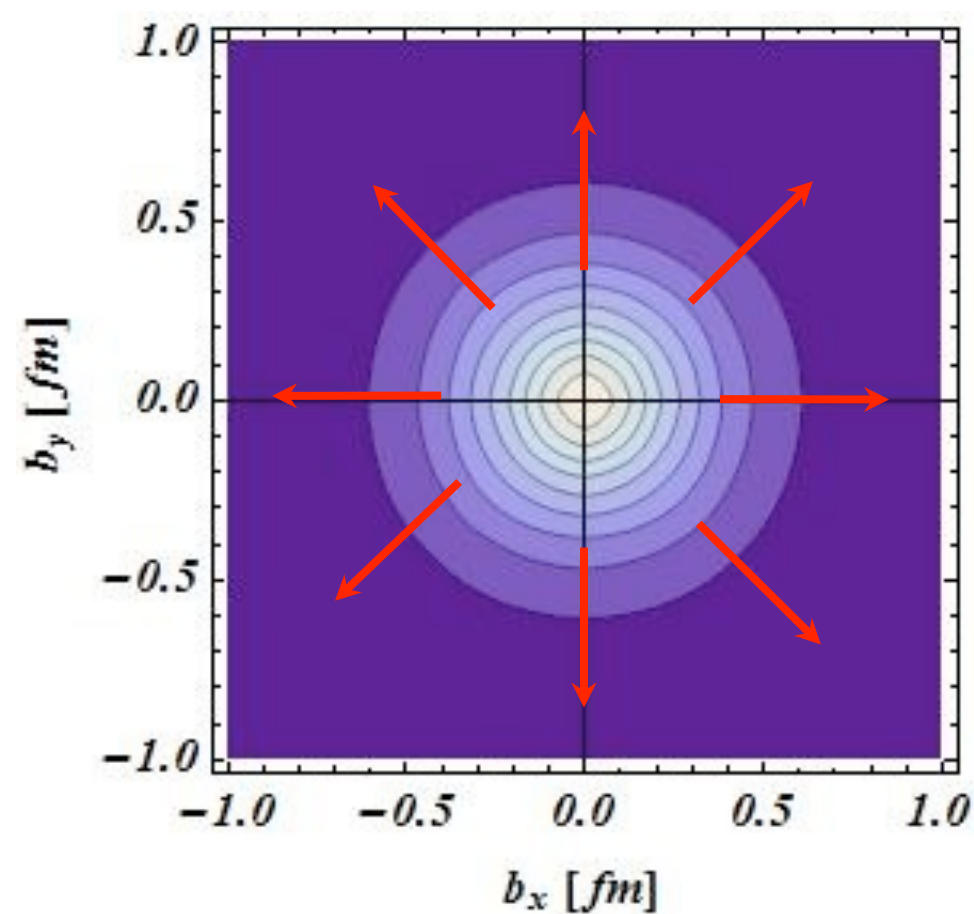
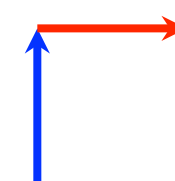
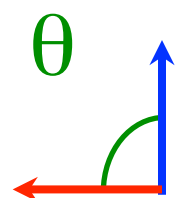


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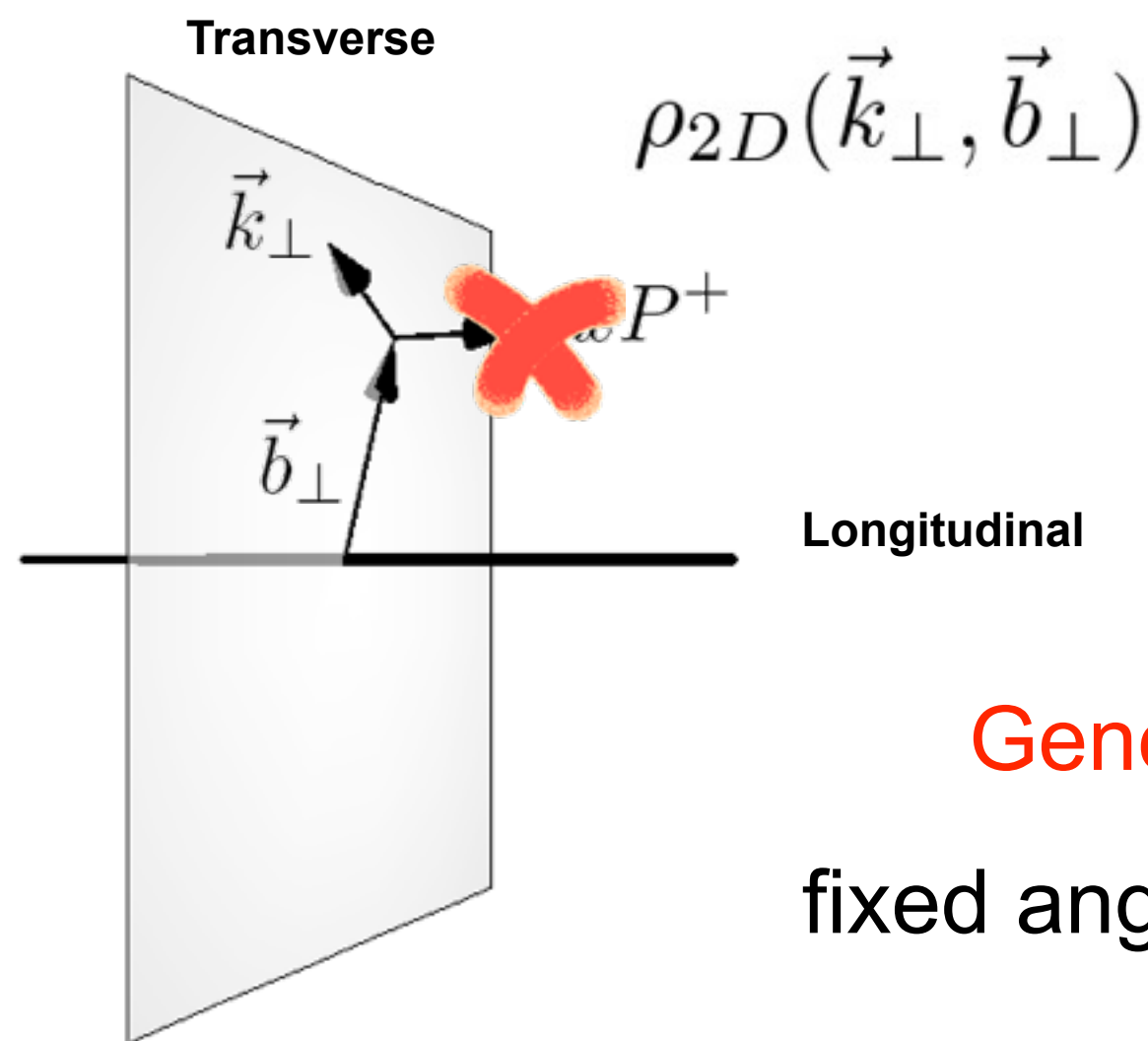
\vec{k}_{\perp}

\vec{b}_{\perp}



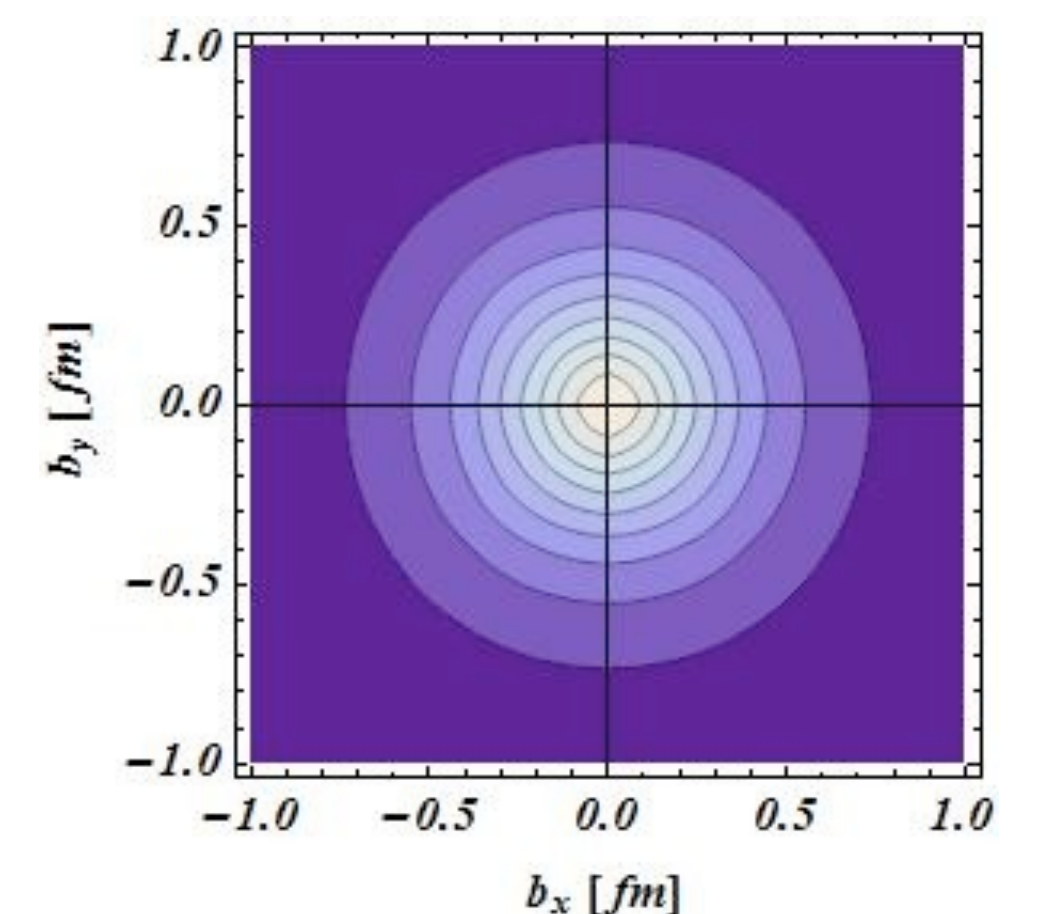
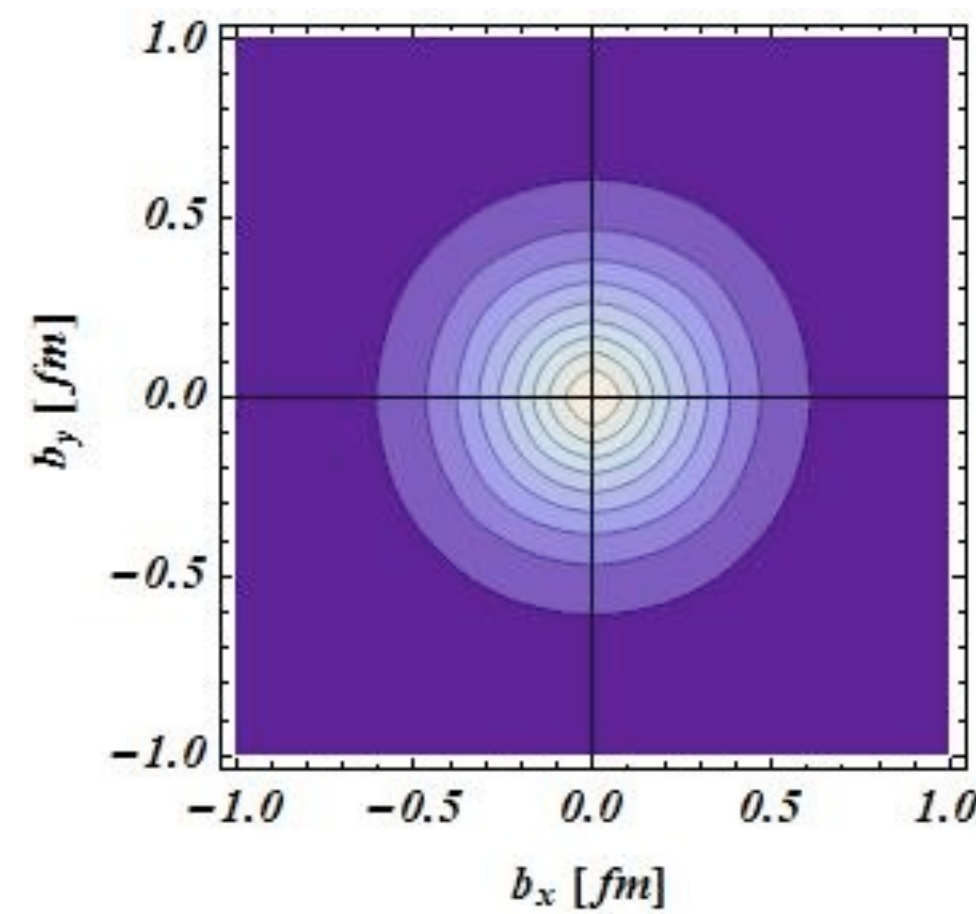
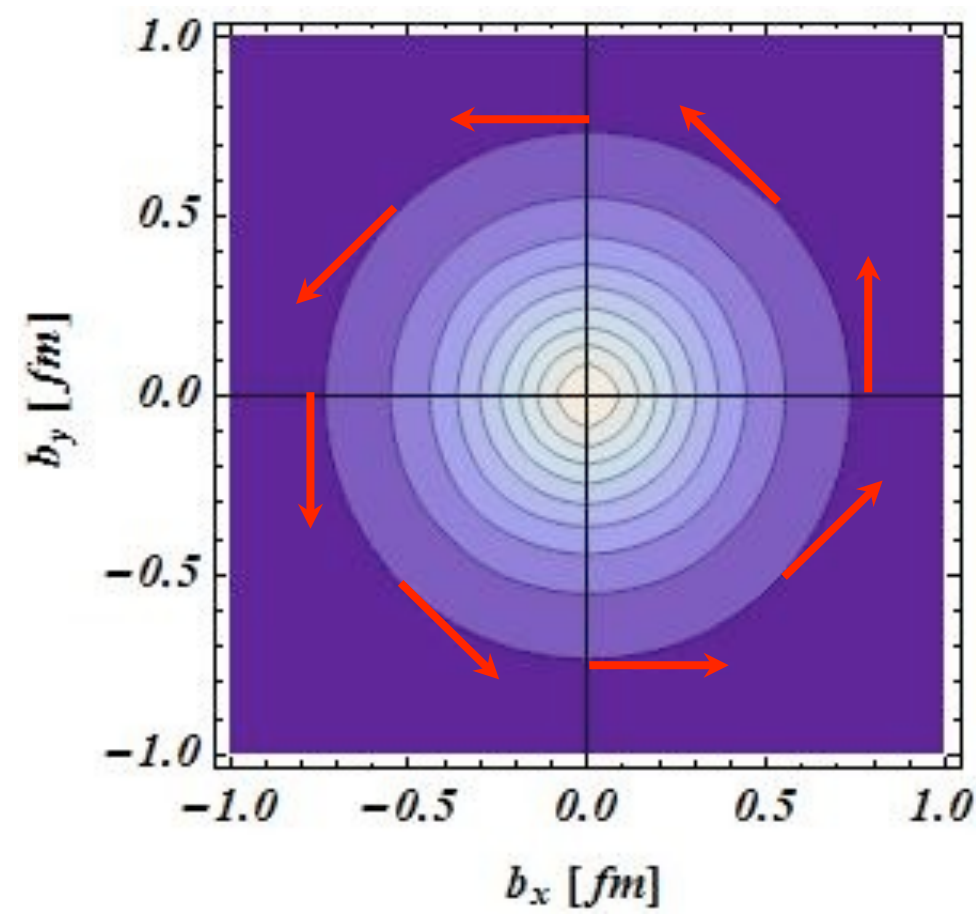
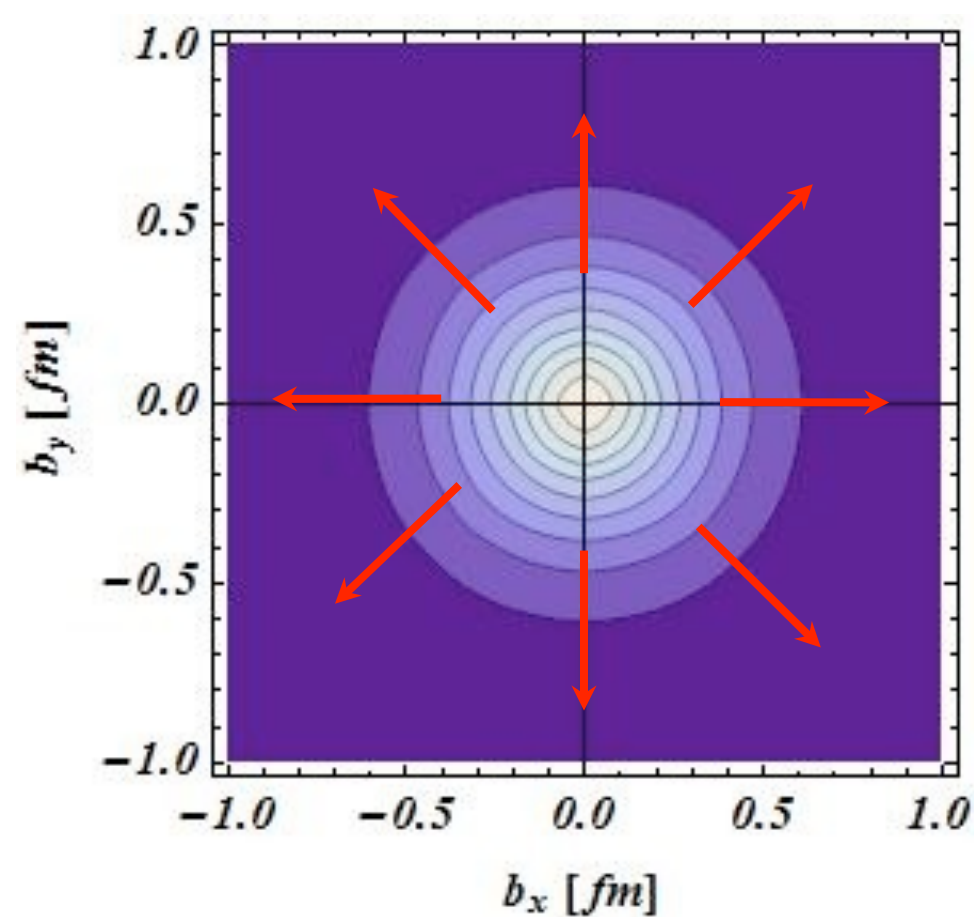
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[Lorce', BP, PRD84 (2011)]



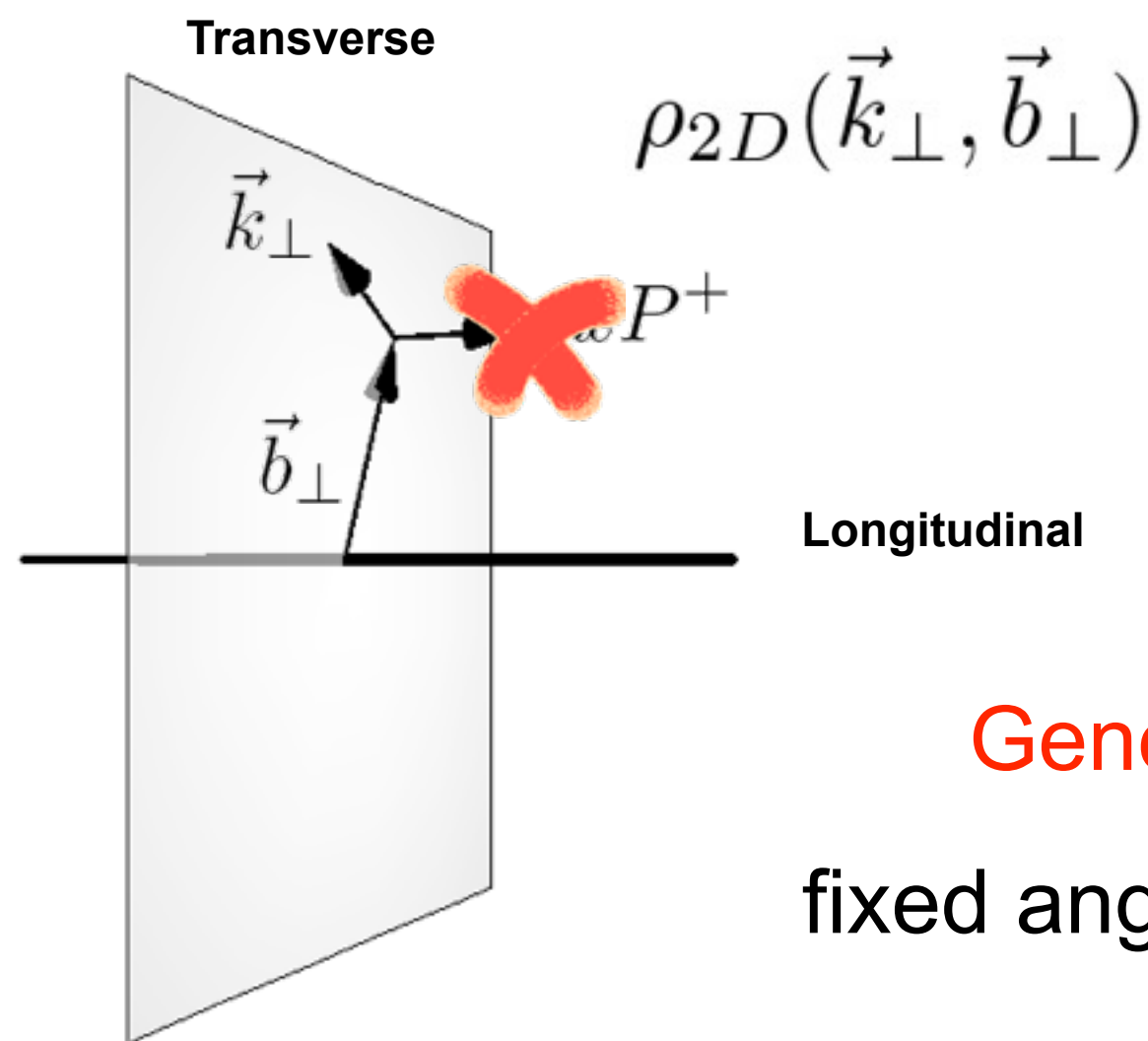
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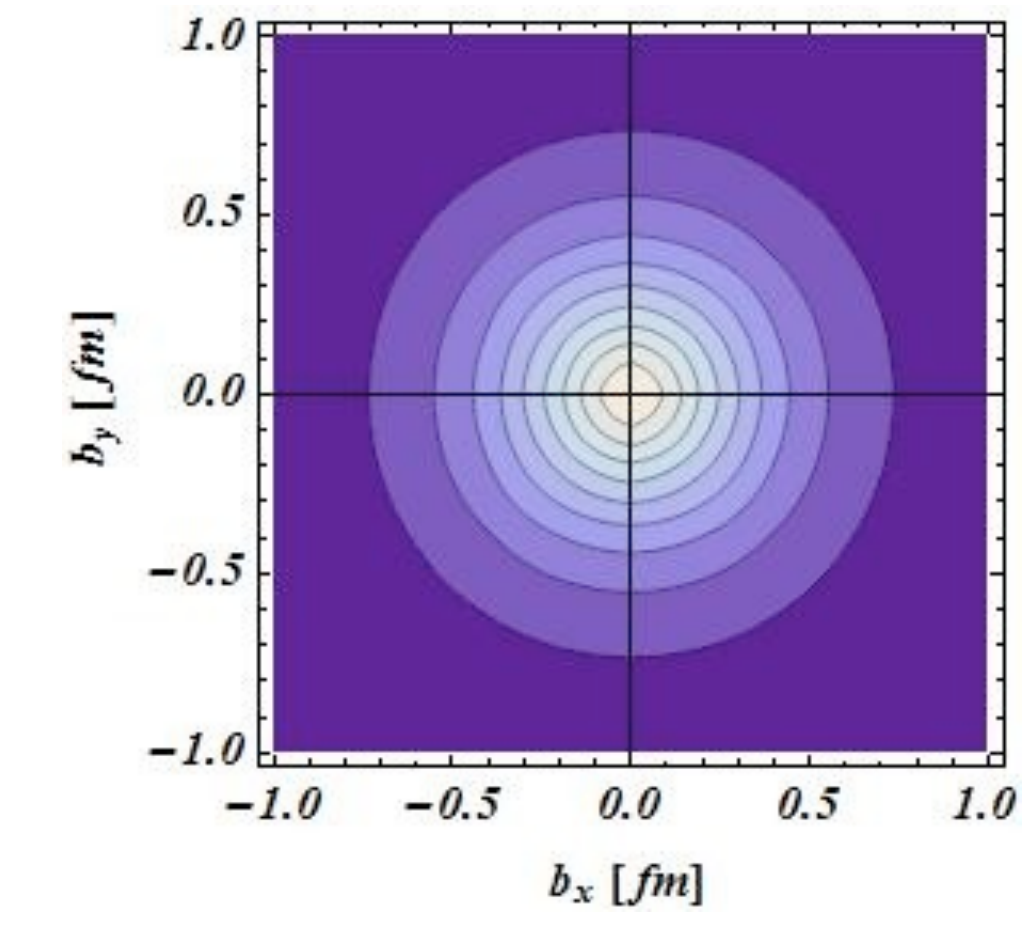
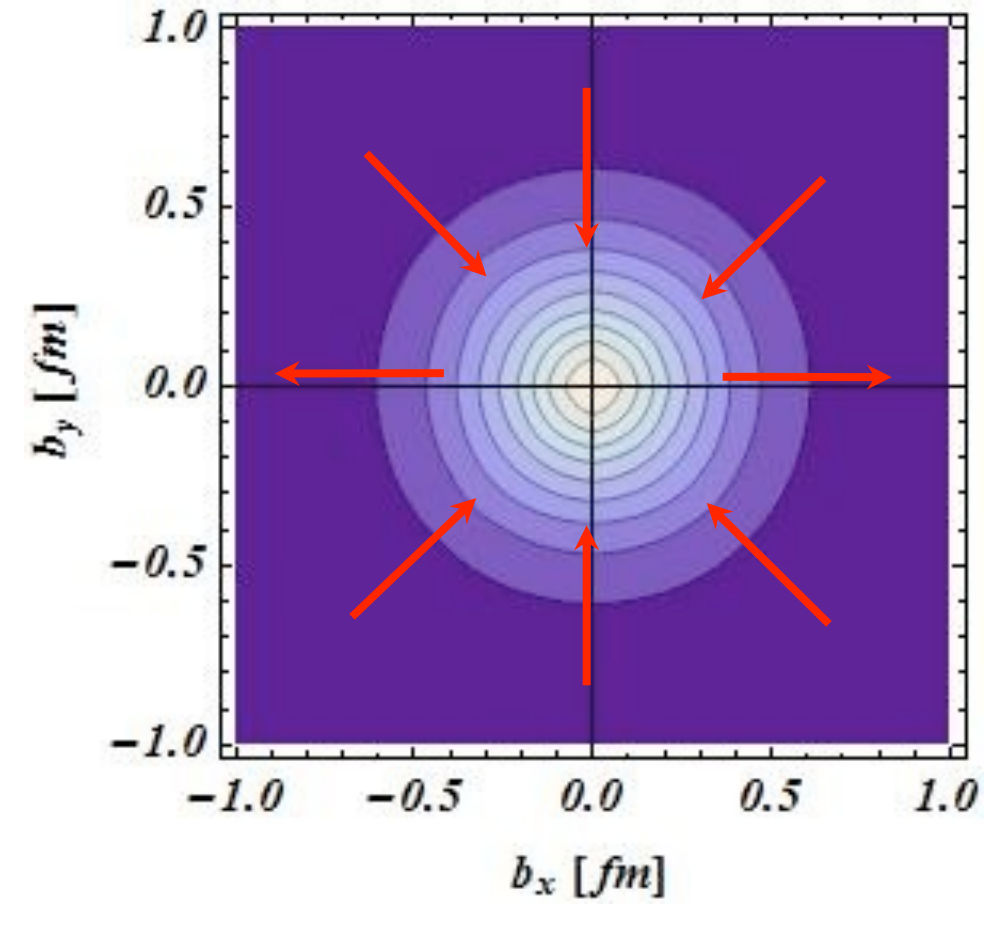
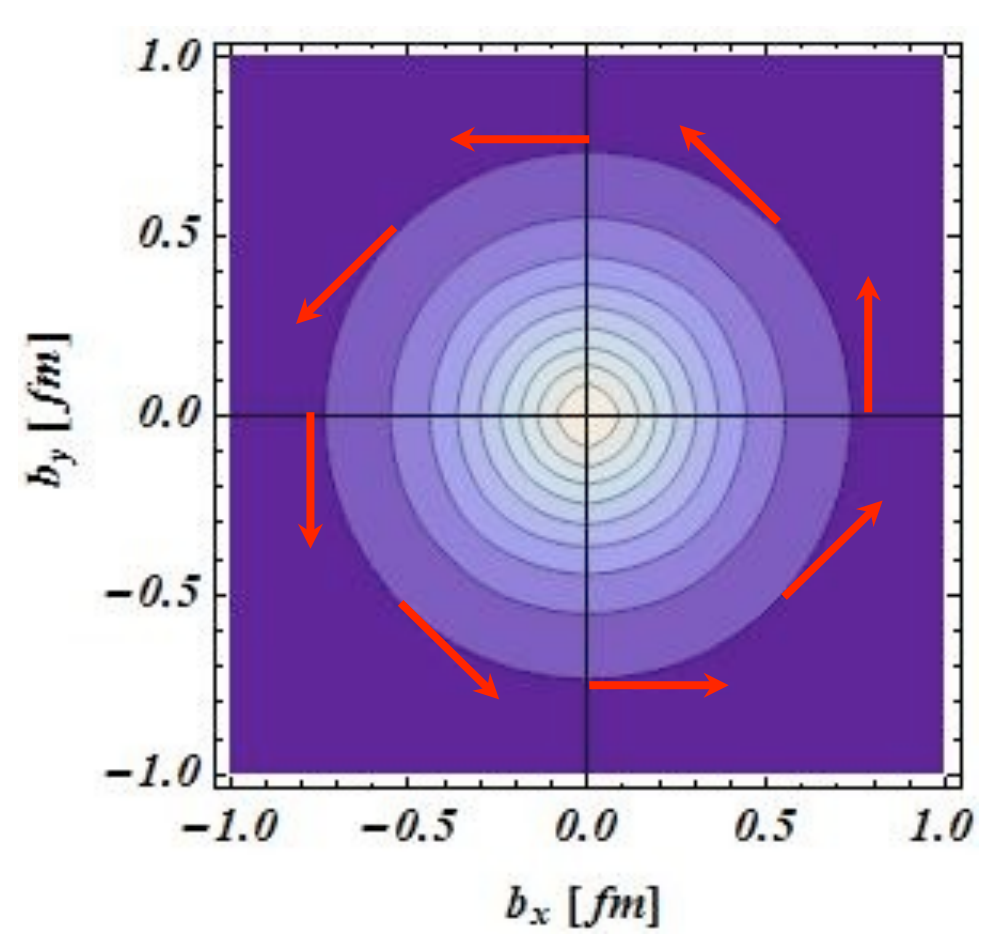
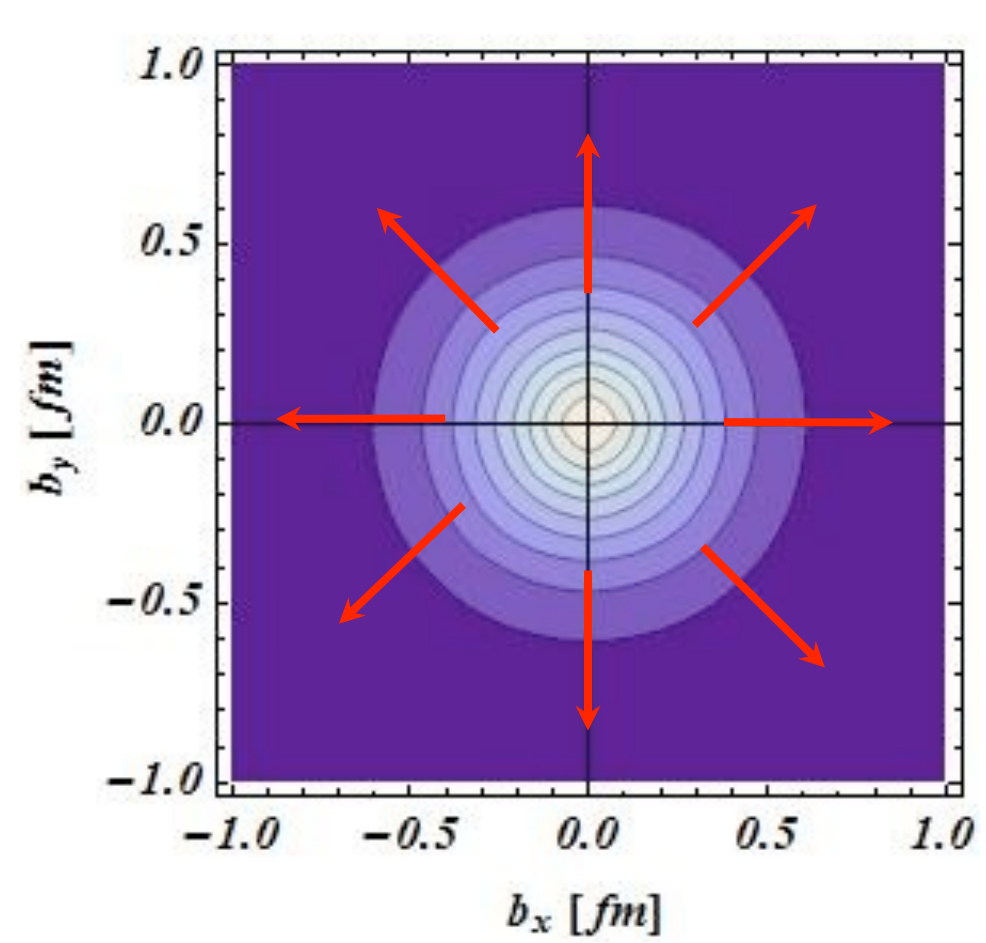
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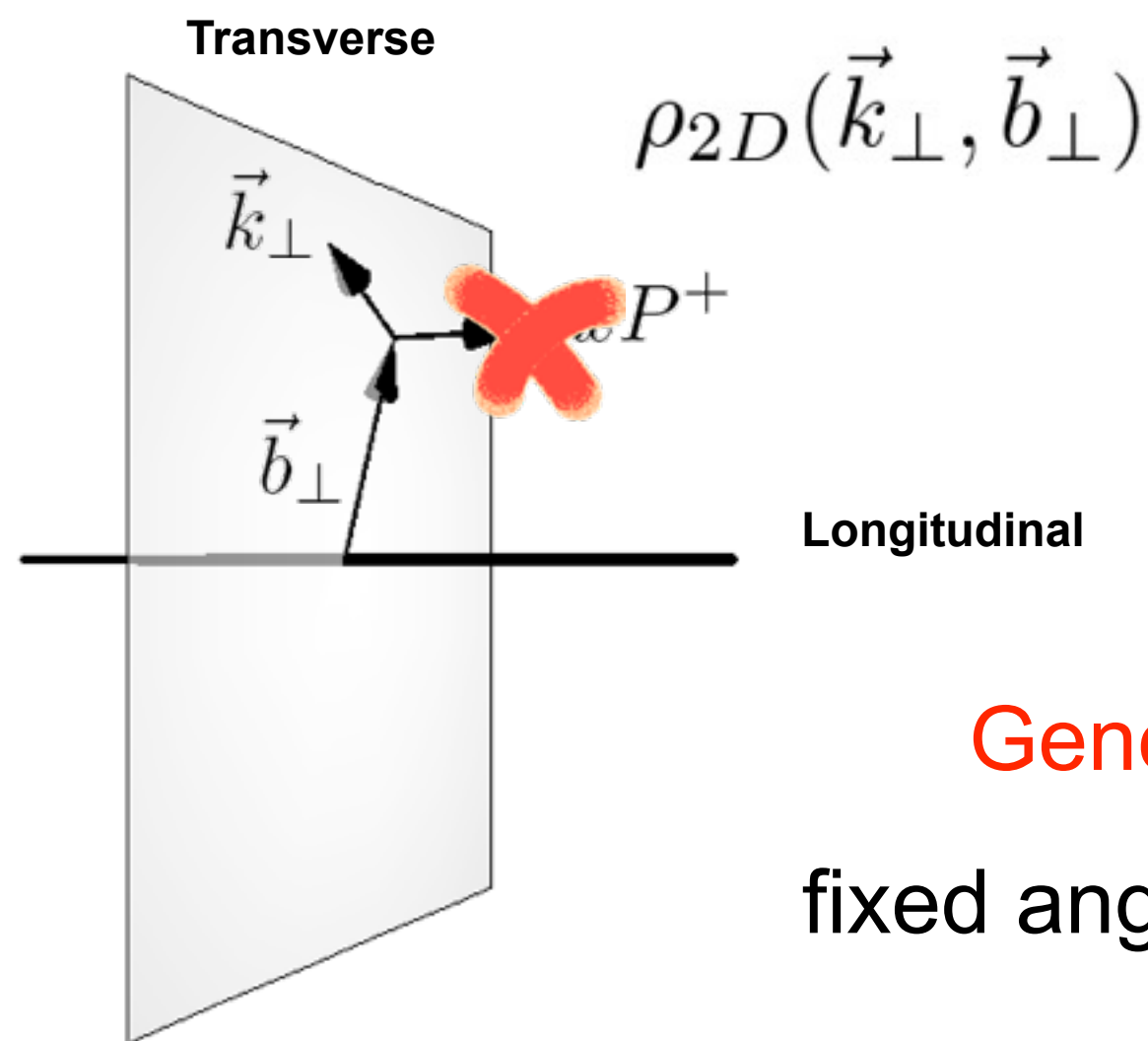
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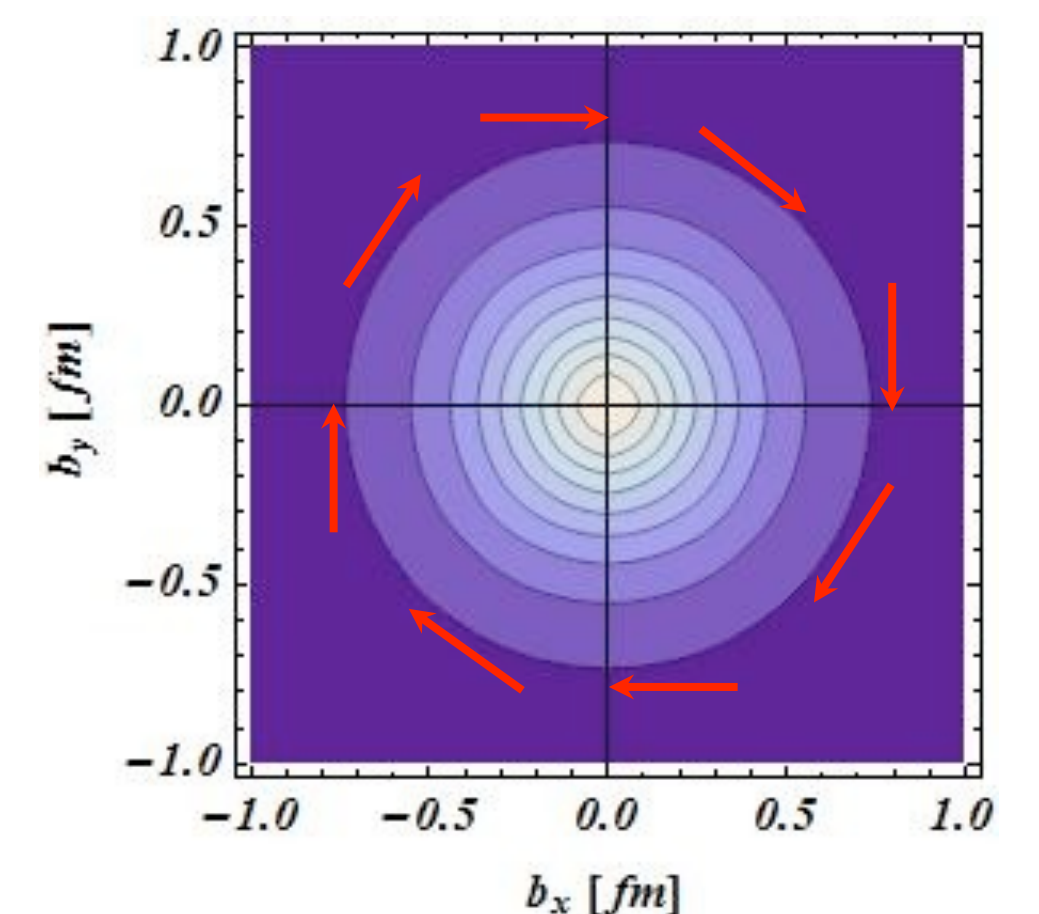
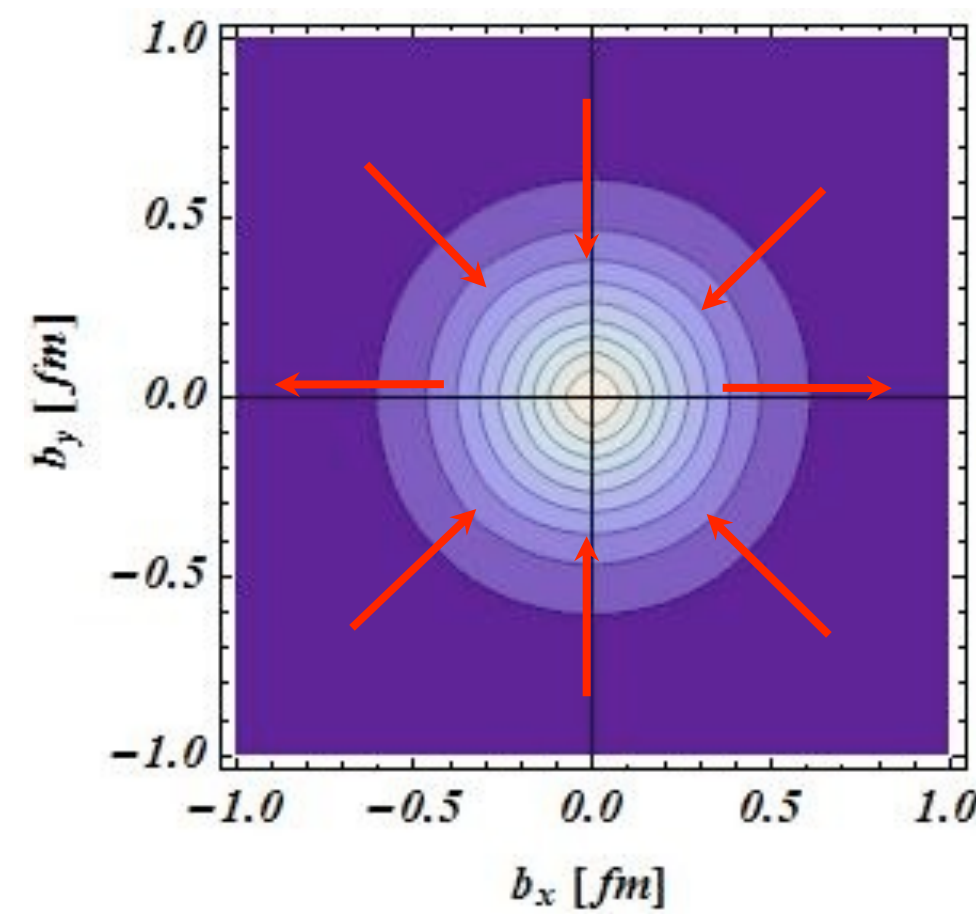
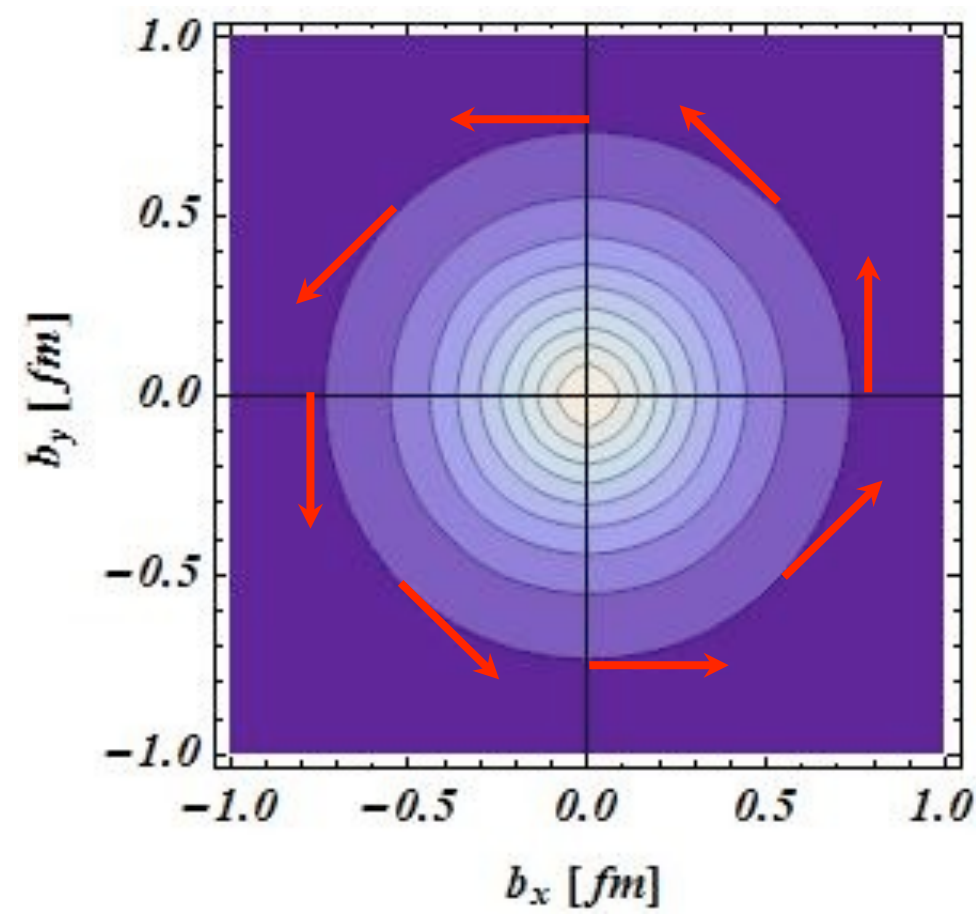
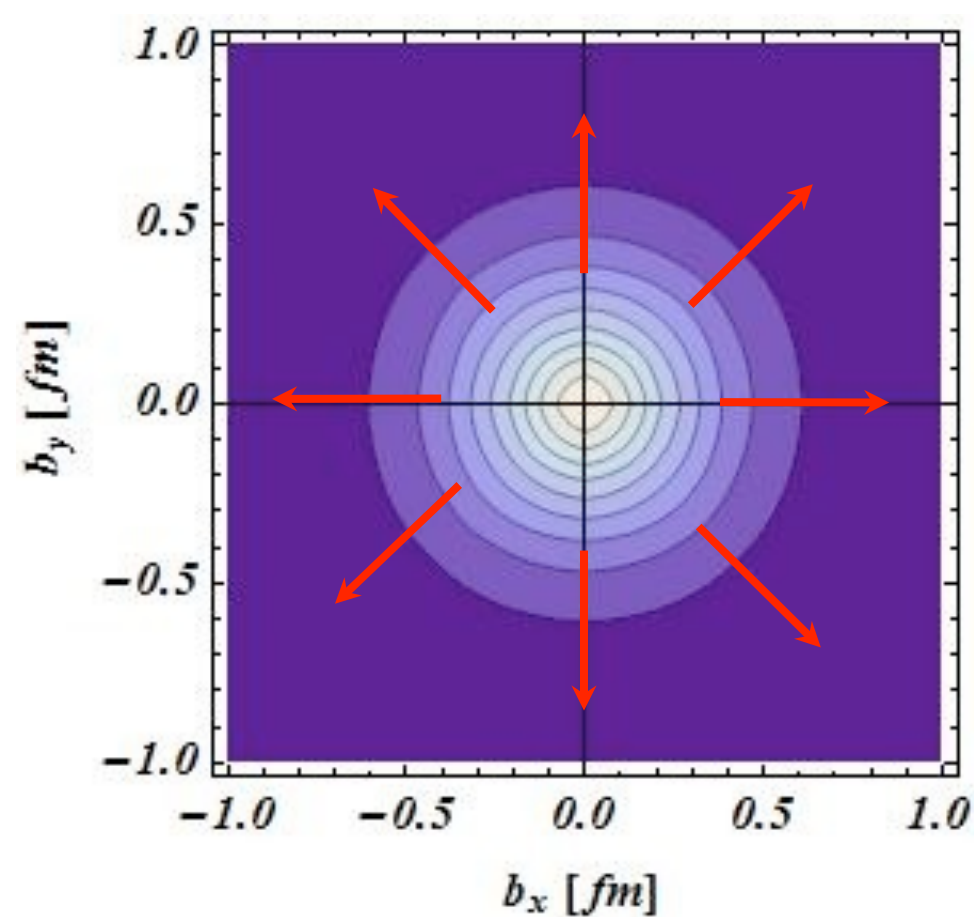
Unpol. up Quark in Unpol. Proton

[Lorce', BP, PRD84 (2011)]



Generalized Transverse Charge Density

fixed angle between \vec{k}_\perp and \vec{b}_\perp and fixed value of $|\vec{k}_\perp|$

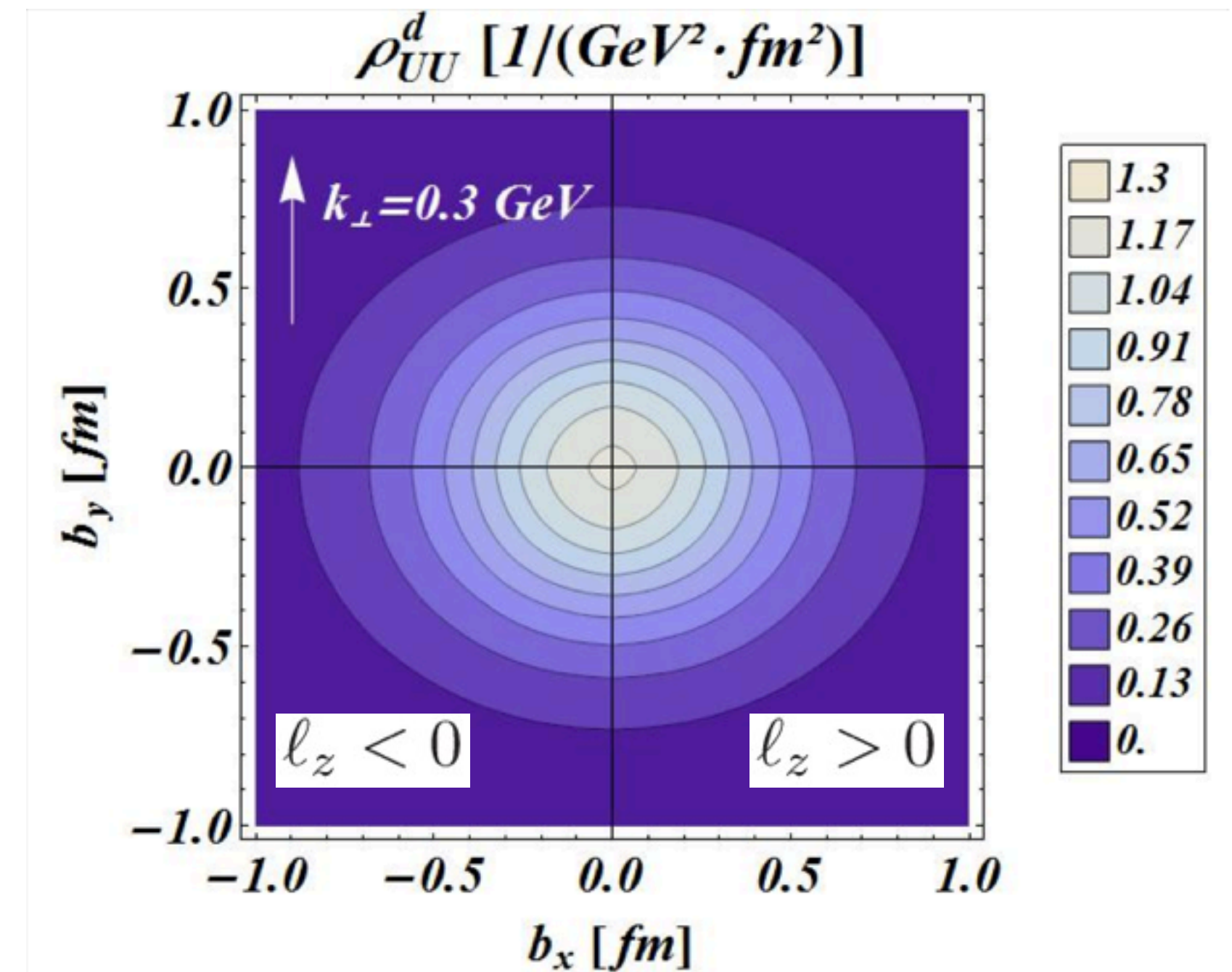
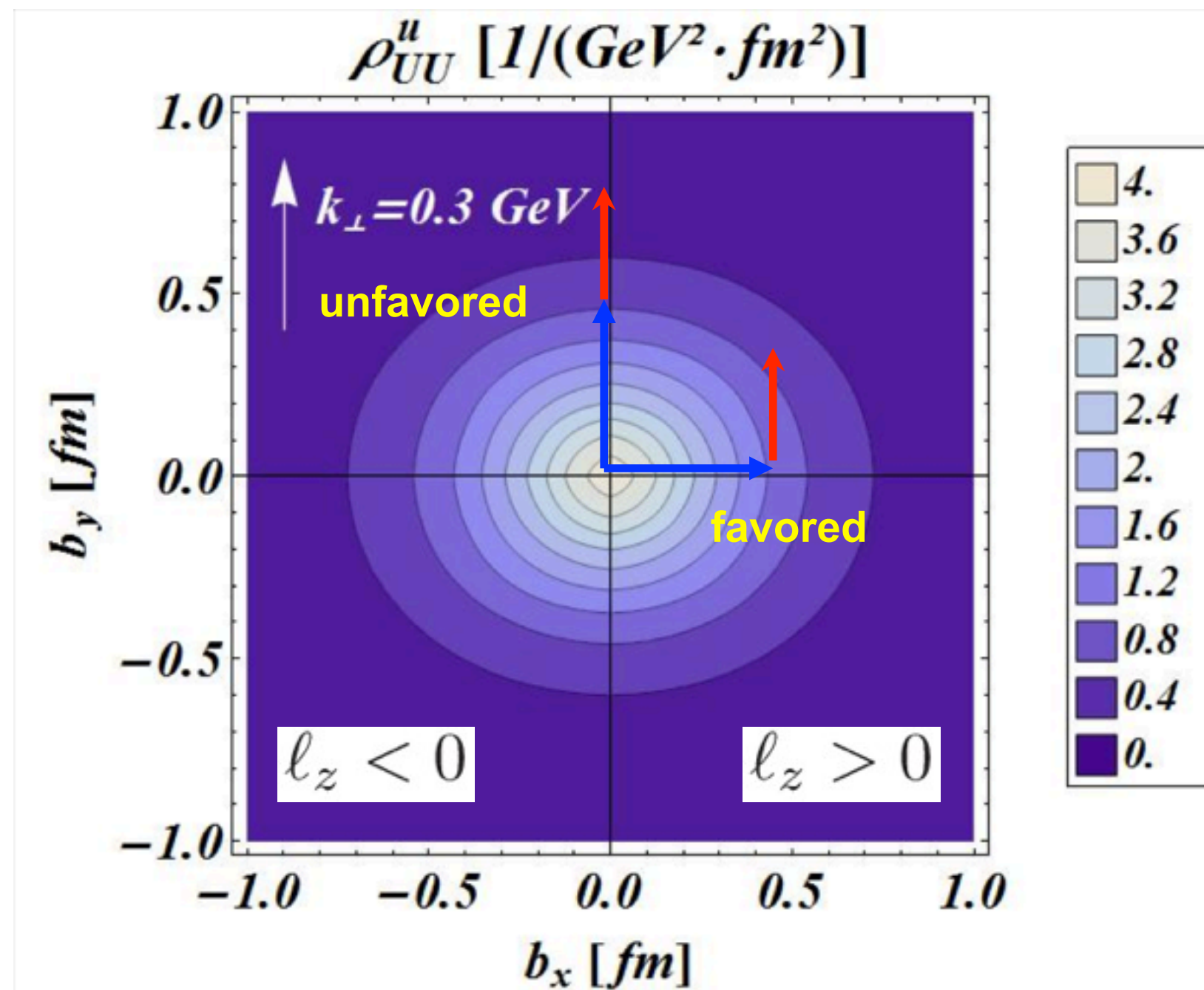


Unpol. up Quark in Unpol. Proton

up quark

down quark

fixed \vec{k}_\perp : \uparrow



Distortion due to correlations between \vec{k}_\perp and \vec{b}_\perp

\searrow absent in **GPDs** and **TMDs** !

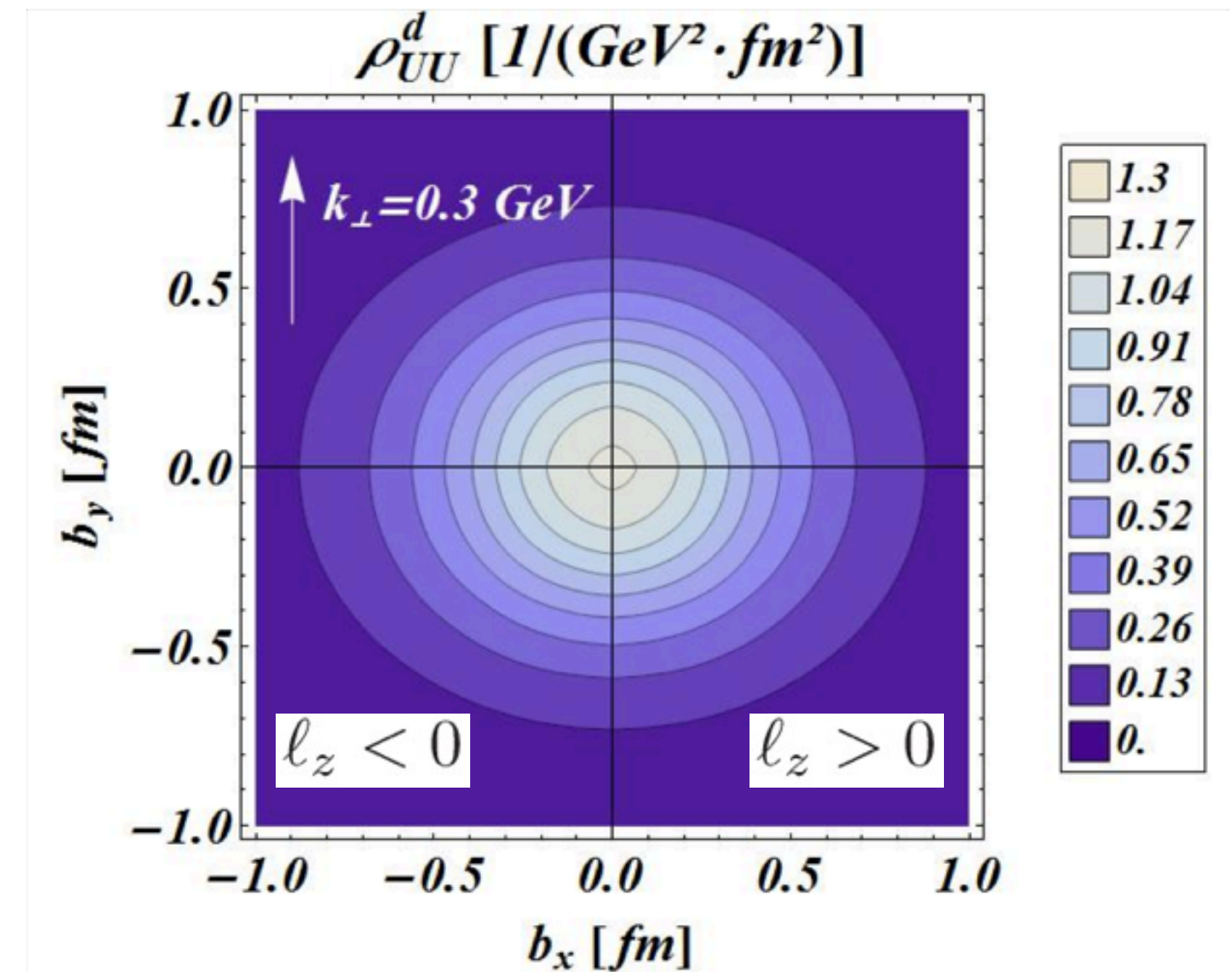
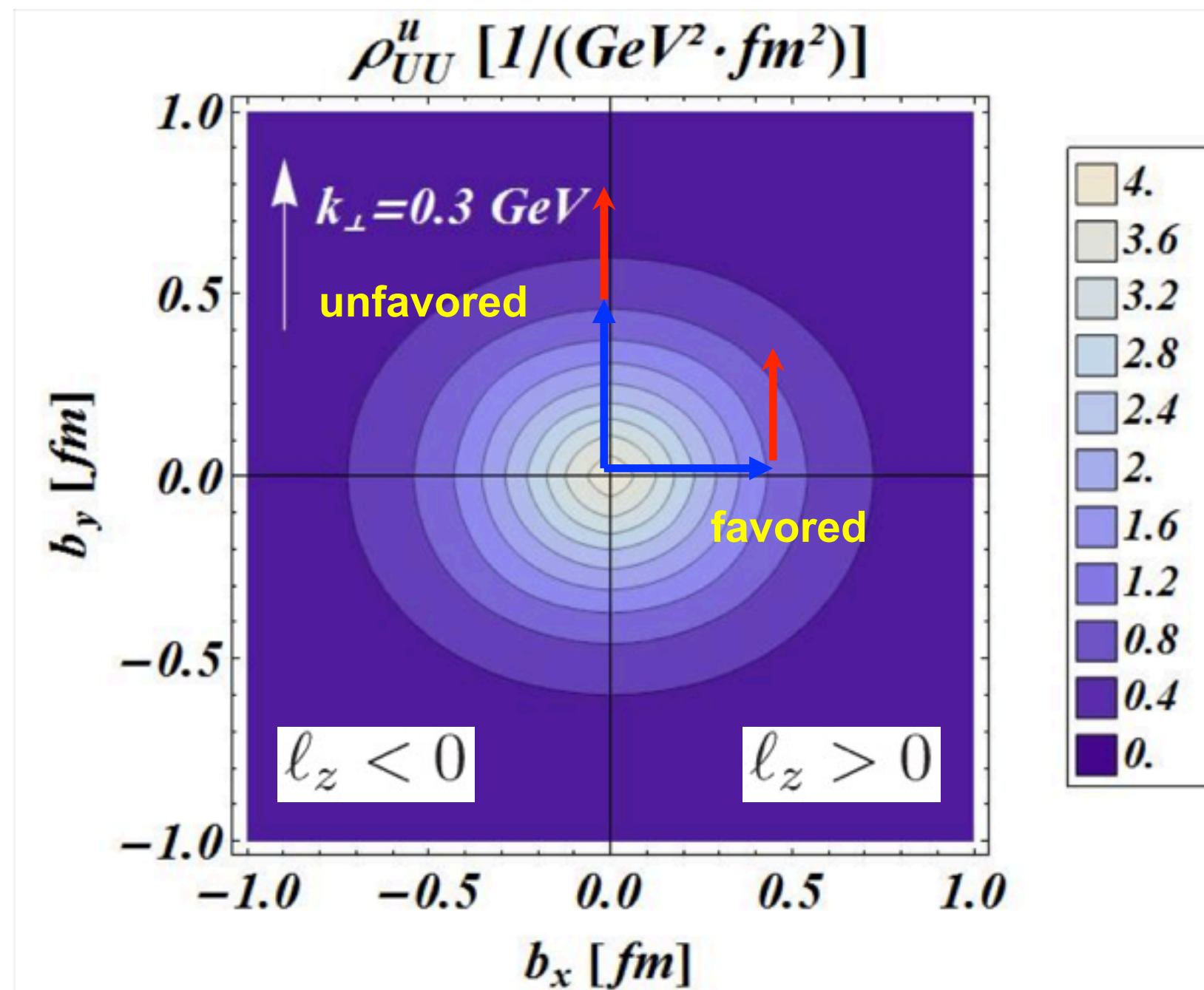
Left-right symmetry \longrightarrow no net quark OAM

Unpol. up Quark in Unpol. Proton

up quark

down quark

fixed \vec{k}_\perp : \uparrow



◆ integrating over \vec{b}_\perp \Rightarrow transverse-momentum density

$$f_1^q(k_\perp^2) = \int dx f_1^q(x, k_\perp^2)$$

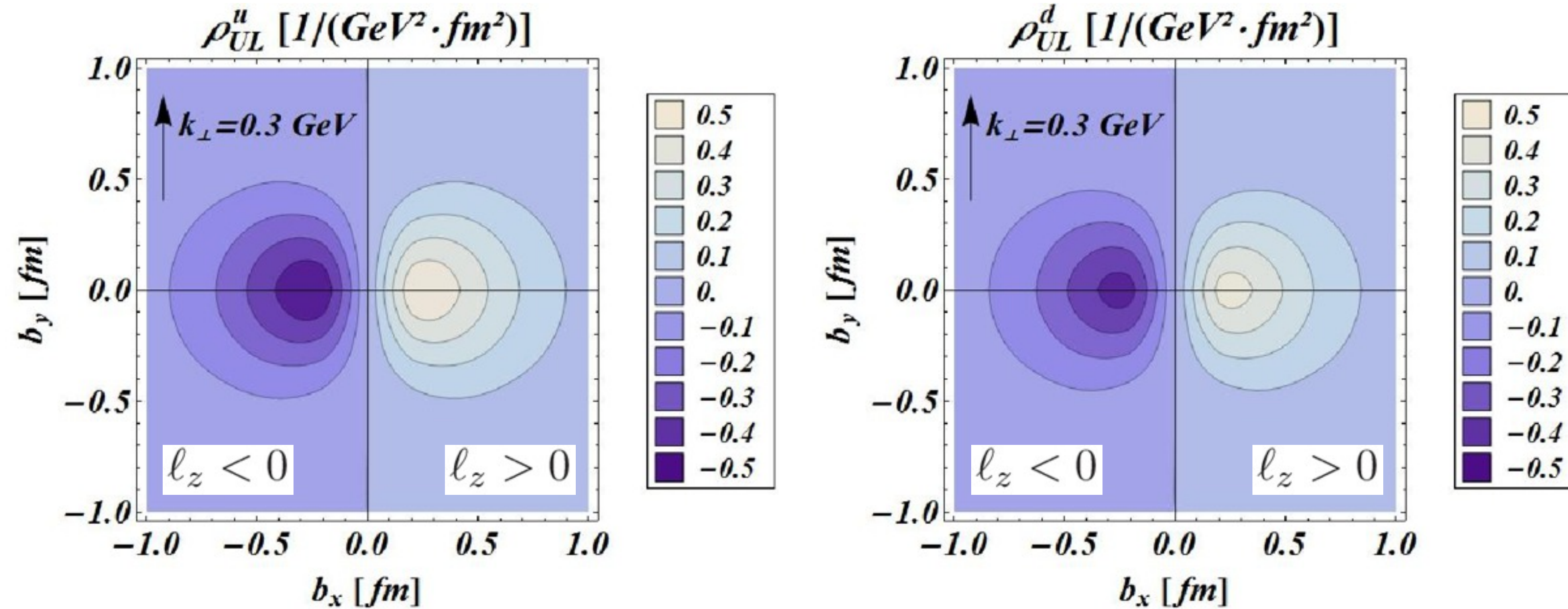
◆ integrating over \vec{k}_\perp \Rightarrow charge density in the transverse plane \vec{b}_\perp
 [Miller (2007); Burkardt (2007)]

$$\rho^q(b_\perp^2) = e^q \int d^2\Delta_\perp e^{-i\vec{b}_\perp \cdot \vec{\Delta}_\perp} F_1^q(\Delta_\perp^2)$$

**Monopole
Distributions**

Long. pol. quark in Unp. proton

fixed $\vec{k}_\perp \uparrow$

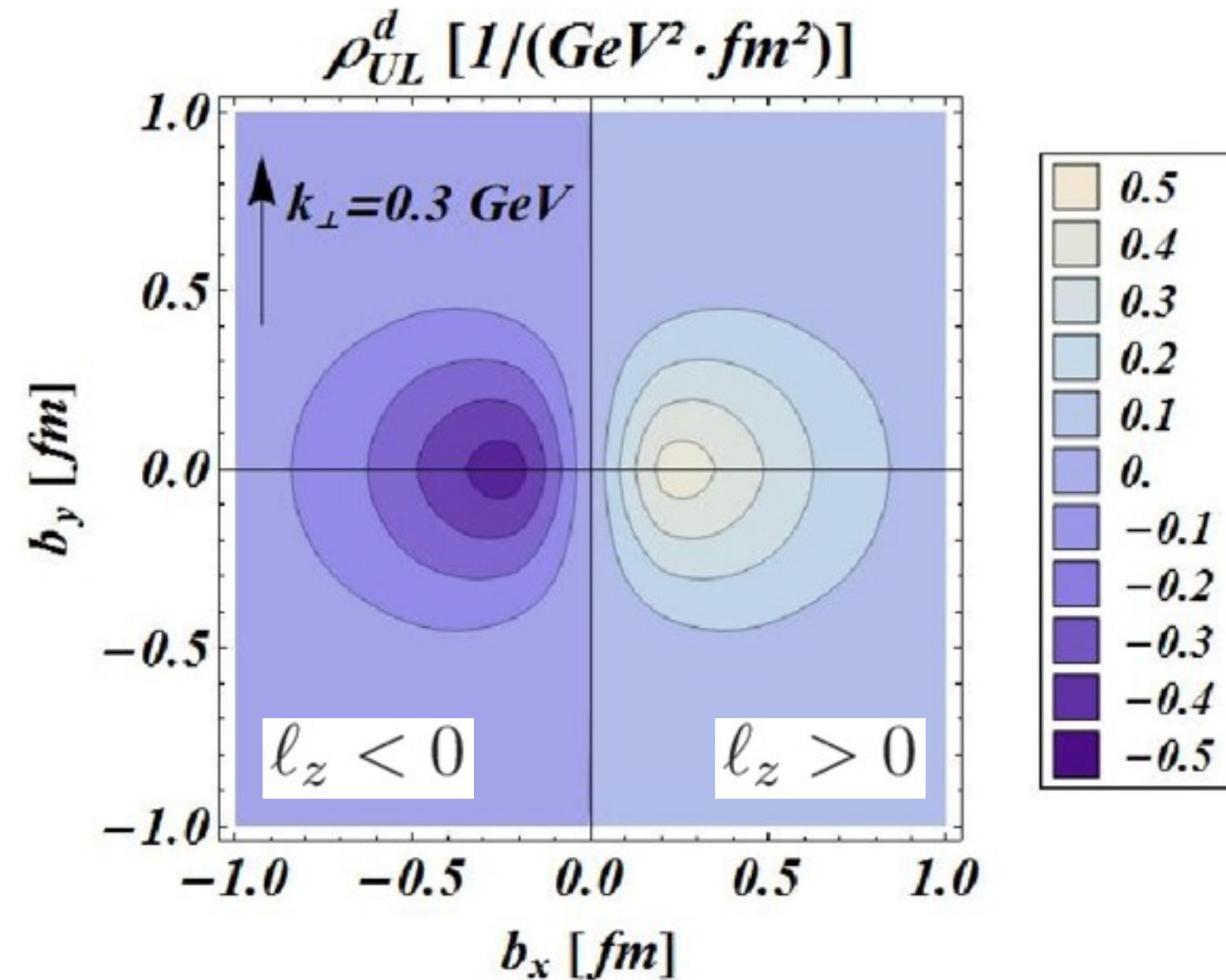
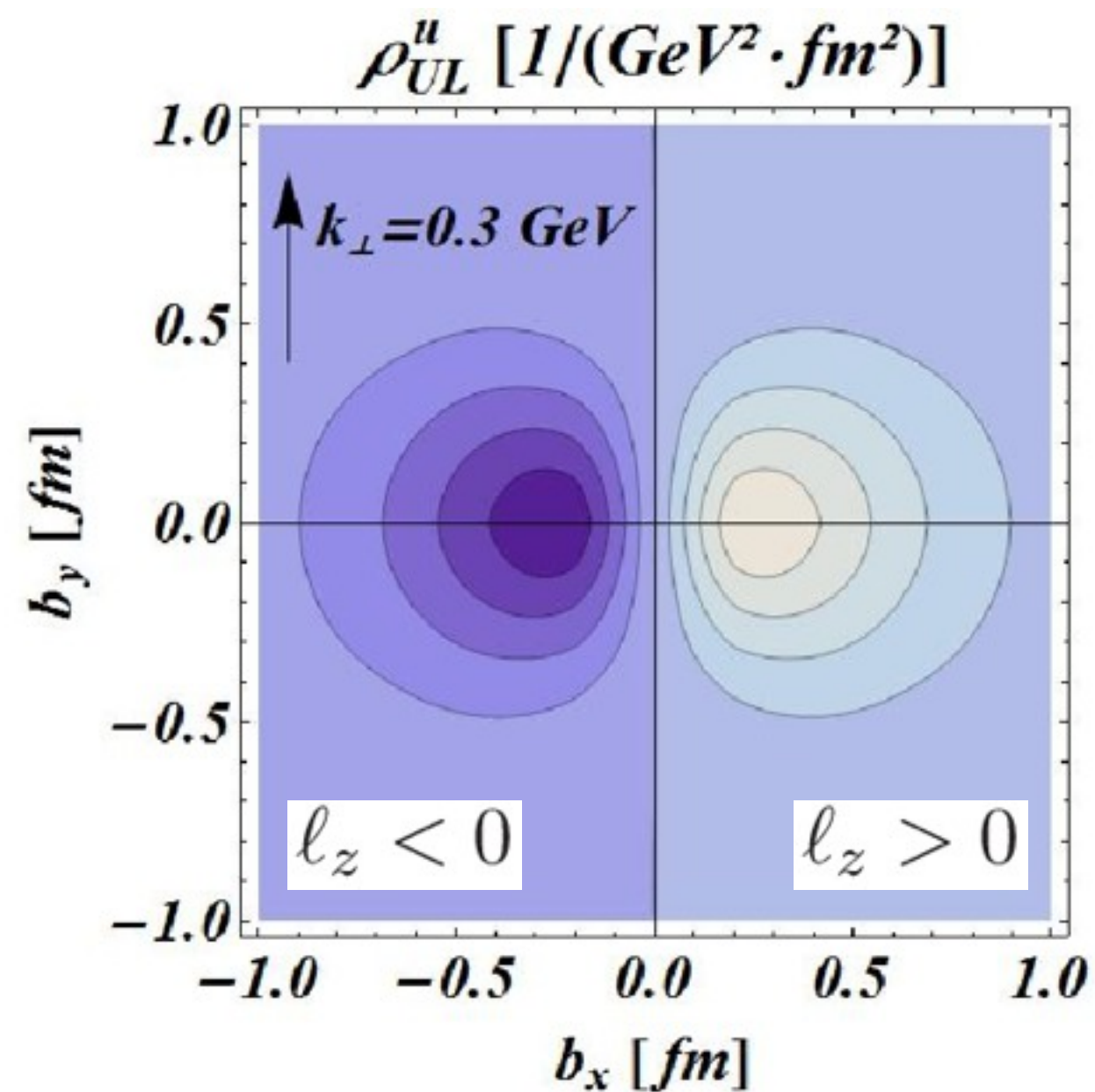


◆ projection to GPD and TMD is vanishing

➡ unique information on OAM from Wigner distributions

Long. pol. quark in Unp. proton

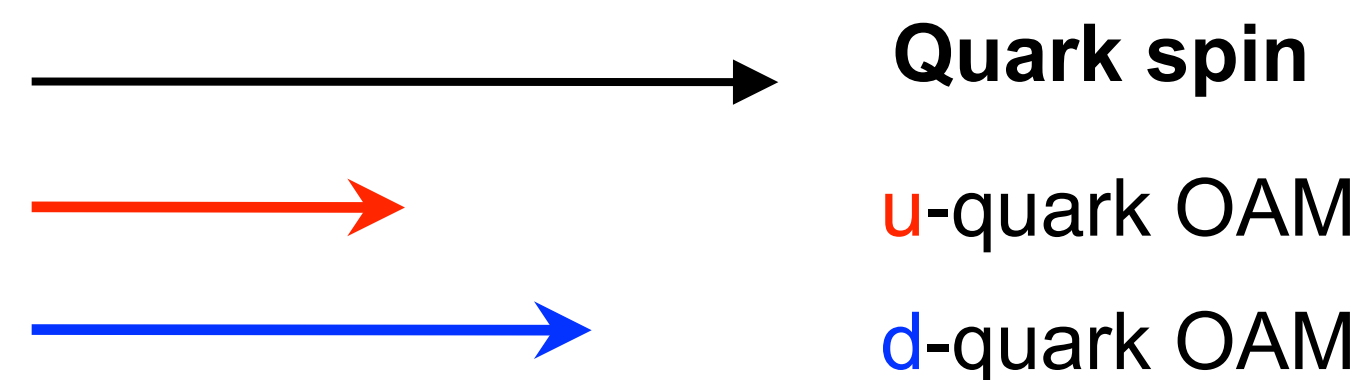
fixed $\vec{k}_\perp \uparrow$



correlation between quark spin and quark OAM

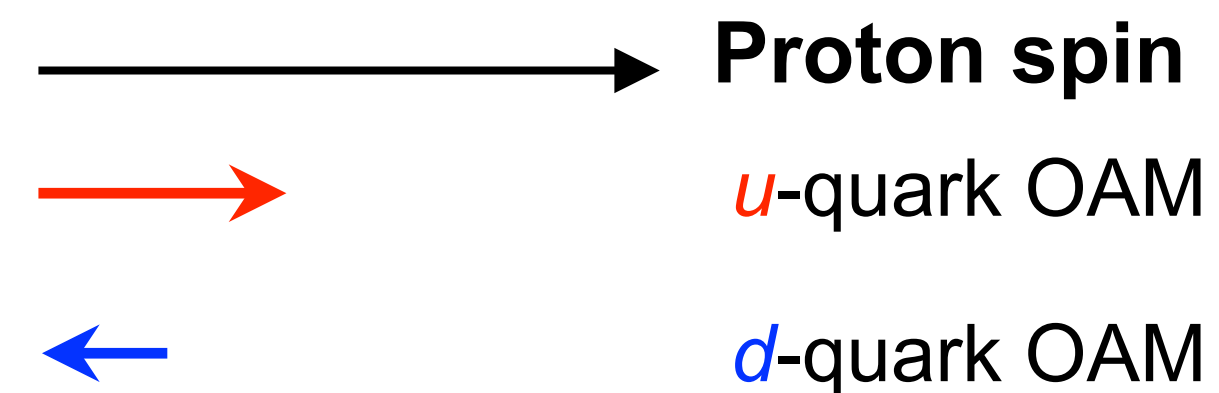
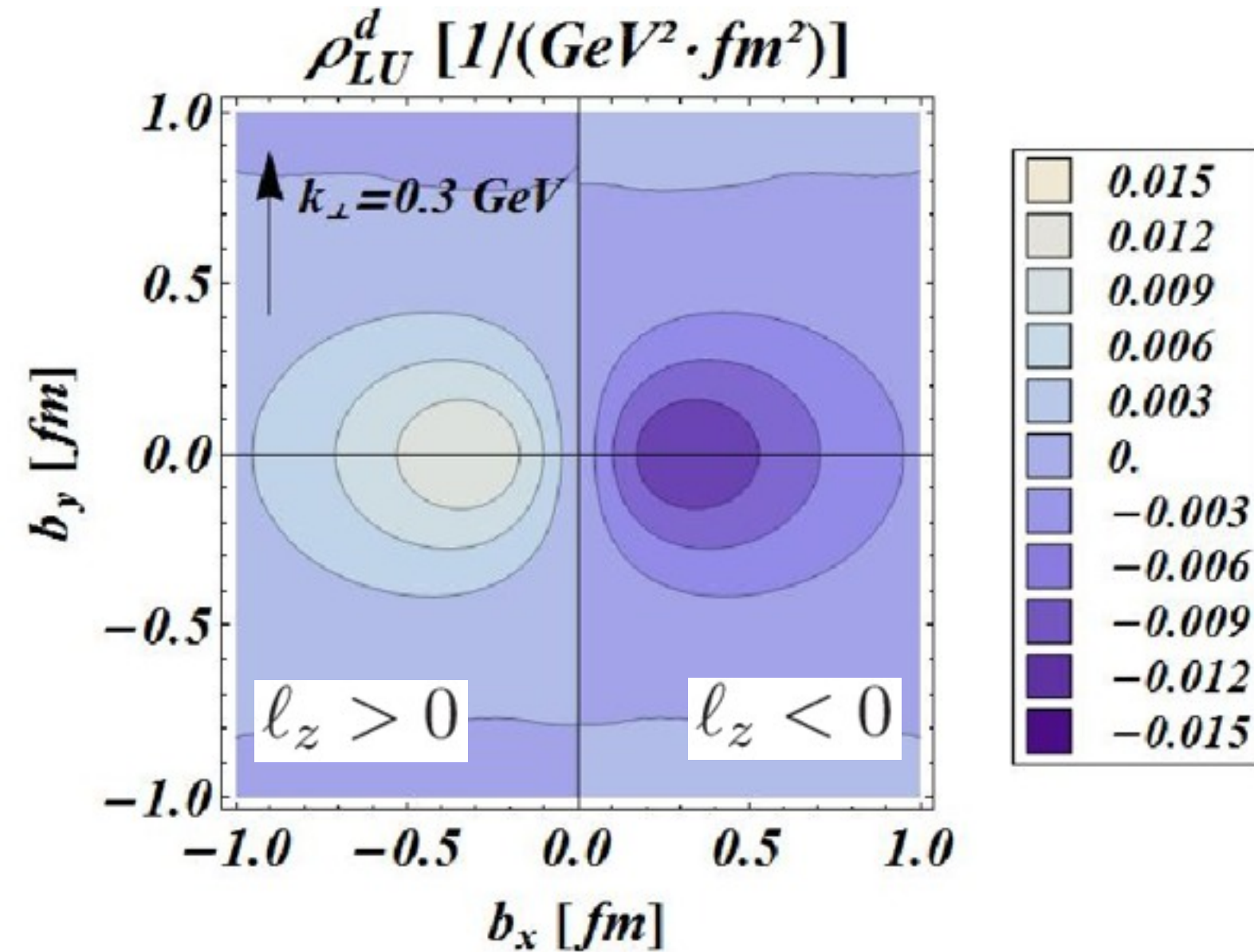
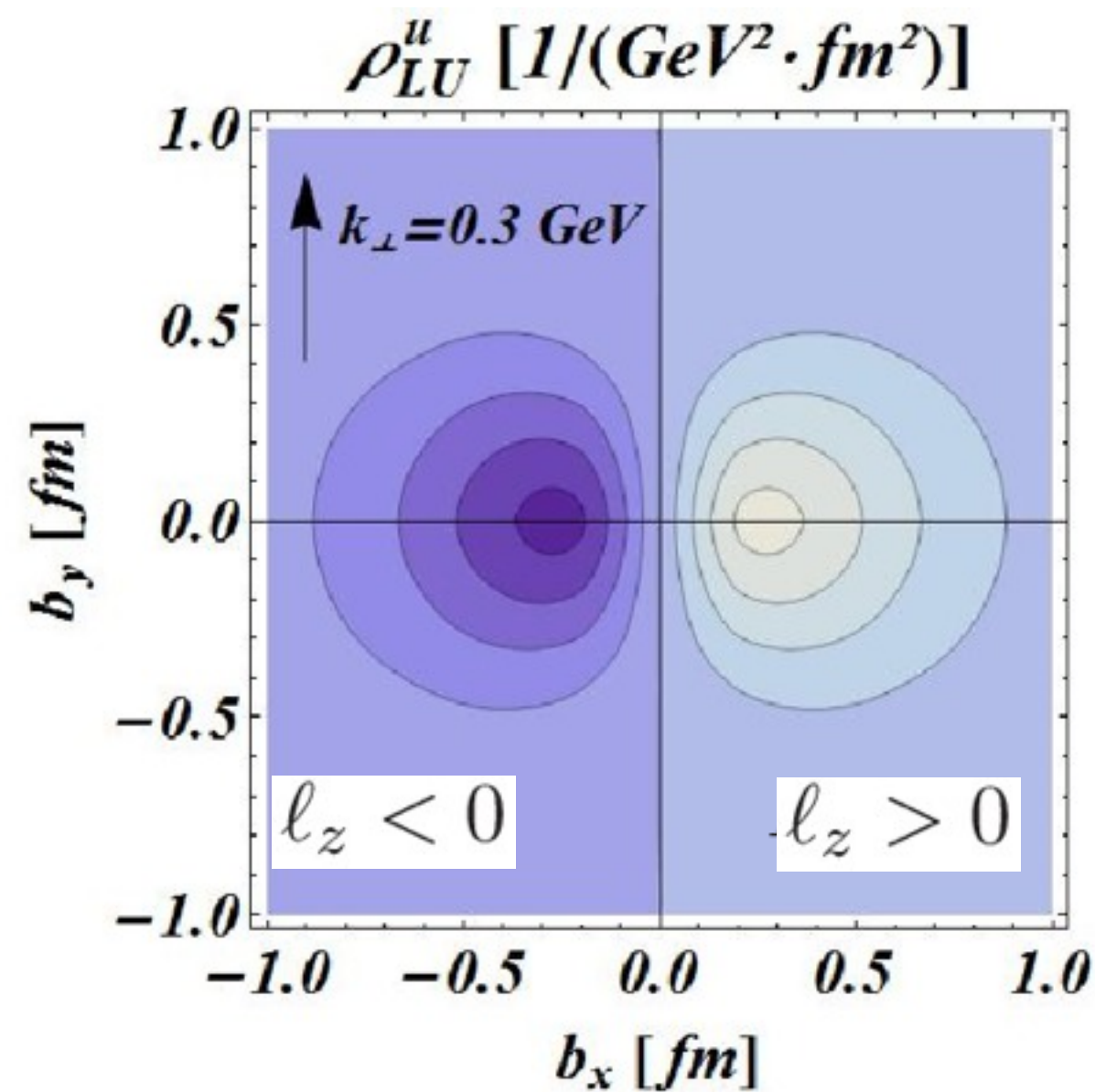
$$C_z^q = \int dx d\vec{k}_\perp d\vec{b}_\perp (\vec{b}_\perp \times \vec{k}_\perp) \rho_{UL}^q(x, \vec{k}_\perp, \vec{b}_\perp)$$

	u-quark	d-quark
C_z^q	0.23	0.19



Unpol. quark in long. pol. proton

fixed $\vec{k}_\perp \uparrow$

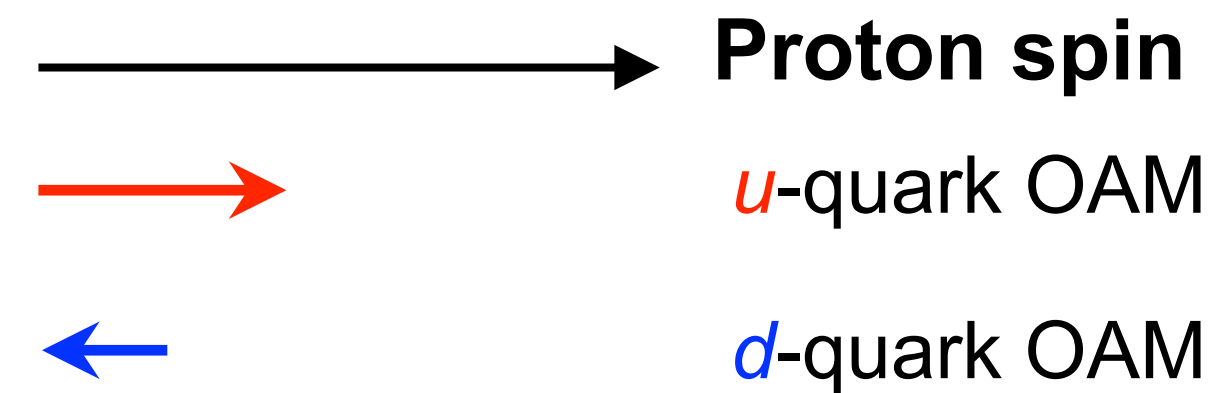
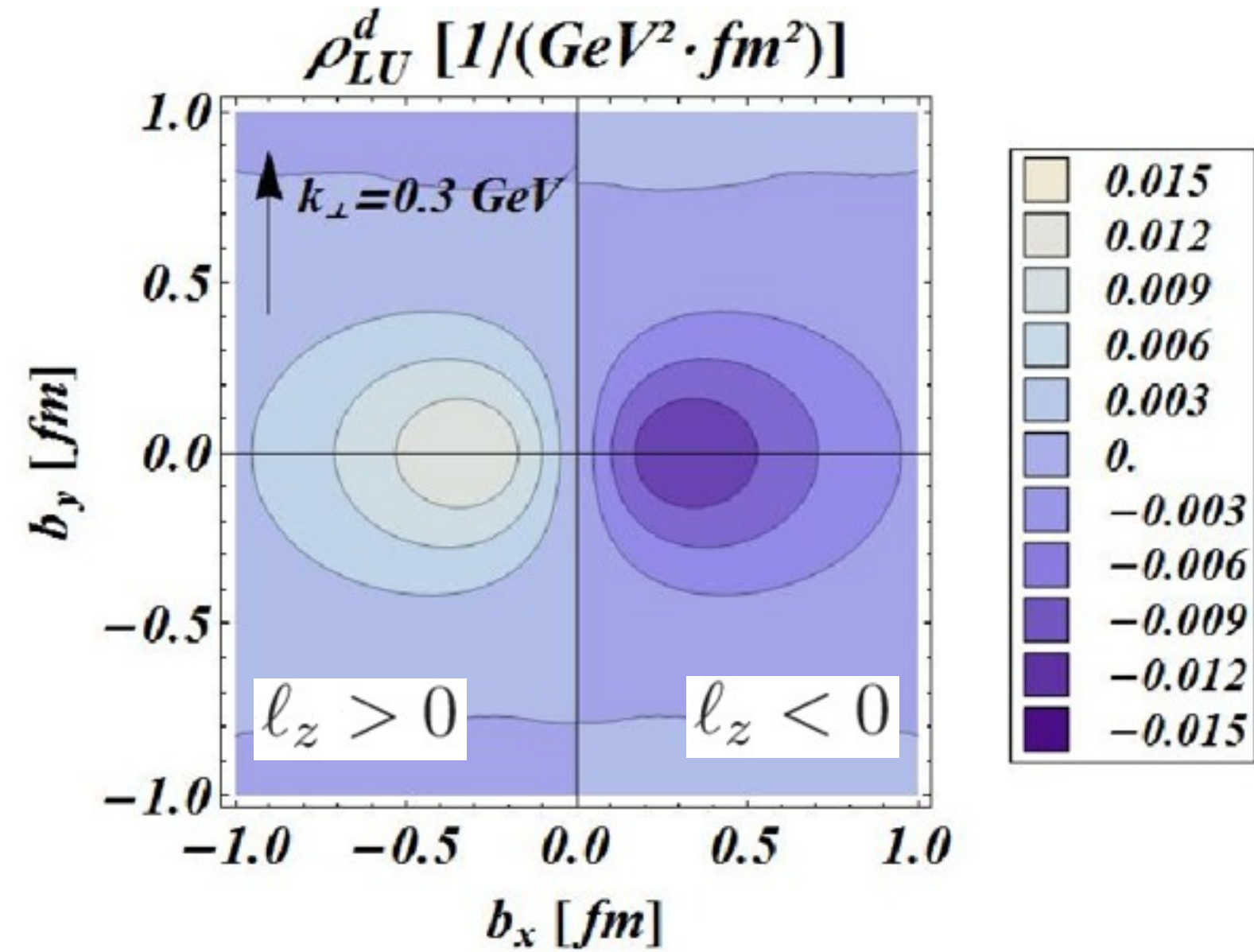
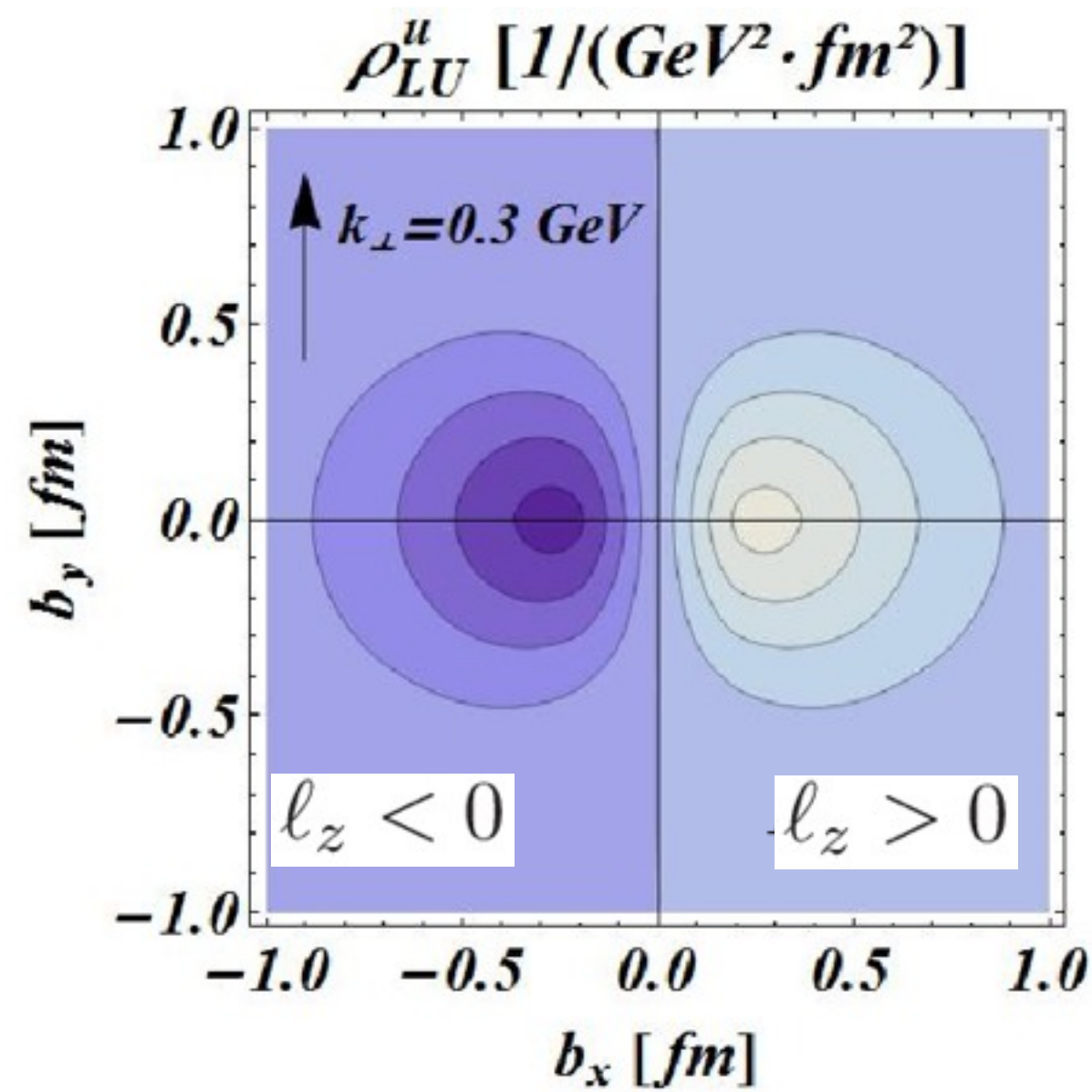


◆ projection to GPD and TMD is vanishing

\longrightarrow unique information on OAM from Wigner distributions

Unpol. quark in long. pol. proton

fixed $\vec{k}_\perp \uparrow$



Quark Orbital Angular Momentum

$$\mathcal{L}_z^q = \int dx d^2\vec{k}_\perp d^2\vec{b}_\perp (\vec{b}_\perp \times \vec{k}_\perp) \rho_{LU}^q(\vec{b}_\perp, \vec{k}_\perp, x)$$



Wigner distribution
for Unpolarized quark in a Longitudinally pol. nucleon

Lorce', BP (11)
Hatta (12)
Ji, Xiong, Yuan (12)

Quark Orbital Angular Momentum

$$\begin{aligned}\mathcal{L}_z^q &= \int dx d^2\vec{k}_\perp d^2\vec{b}_\perp (\vec{b}_\perp \times \vec{k}_\perp) \rho_{LU}^q(\vec{b}_\perp, \vec{k}_\perp, x) \\ &= \int d^2\vec{b}_\perp \vec{b}_\perp \times \langle \vec{k}_\perp^q \rangle \longrightarrow \langle \vec{k}_\perp^q \rangle = \int dx d^2\vec{k}_\perp \vec{k}_\perp \rho_{LU}^q(\vec{b}_\perp, \vec{k}_\perp, x)\end{aligned}$$

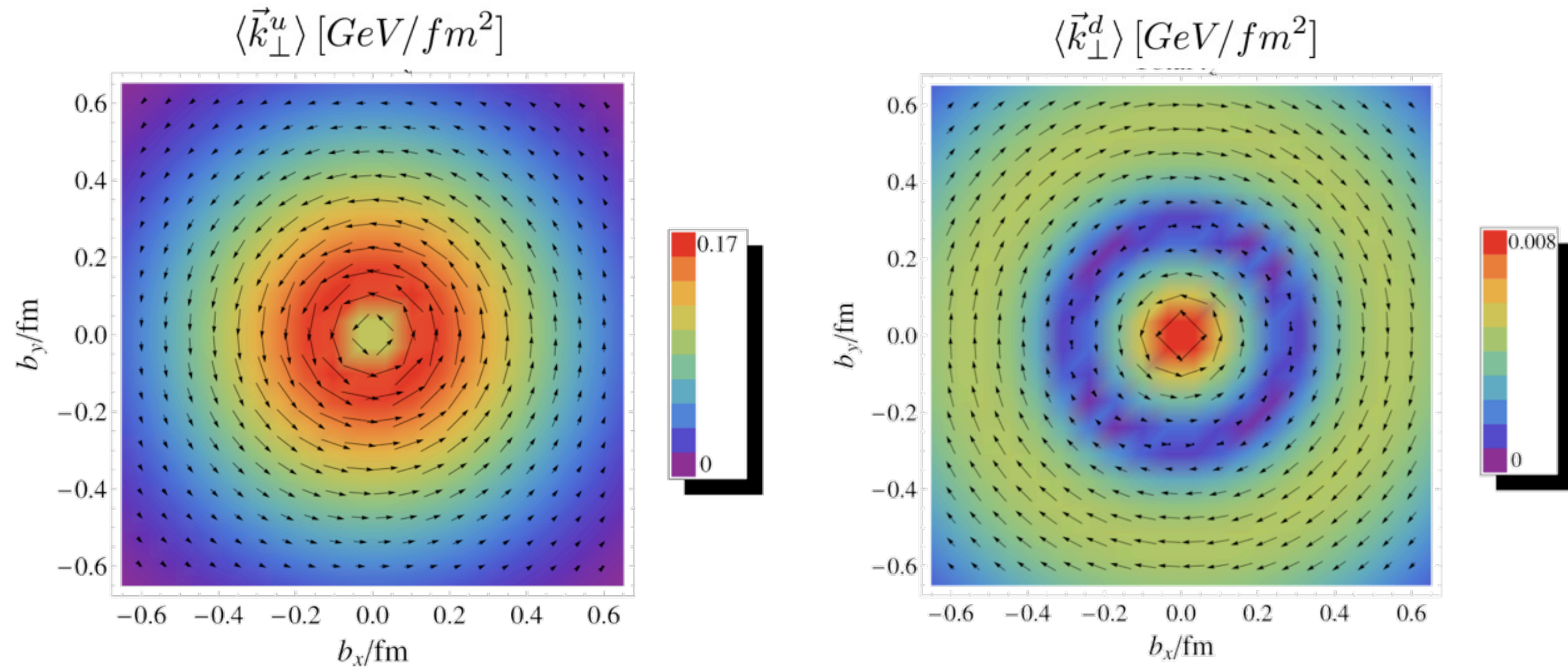
Lorce', BP (11)
Hatta (12)
Ji, Xiong, Yuan (12)




Quark Orbital Angular Momentum

$$\mathcal{L}_z^q = \int dx d^2\vec{k}_\perp d^2\vec{b}_\perp (\vec{b}_\perp \times \vec{k}_\perp) \rho_{LU}^q(\vec{b}_\perp, \vec{k}_\perp, x)$$

$$= \int d^2\vec{b}_\perp \vec{b}_\perp \times \langle \vec{k}_\perp^q \rangle \longrightarrow \langle \vec{k}_\perp^q \rangle = \int dx d^2\vec{k}_\perp \vec{k}_\perp \rho_{LU}^q(\vec{b}_\perp, \vec{k}_\perp, x)$$

Lorce', BP (11)
Hatta (12)
Ji, Xiong, Yuan (12)



 Proton spin
 *u*-quark OAM
 *d*-quark OAM

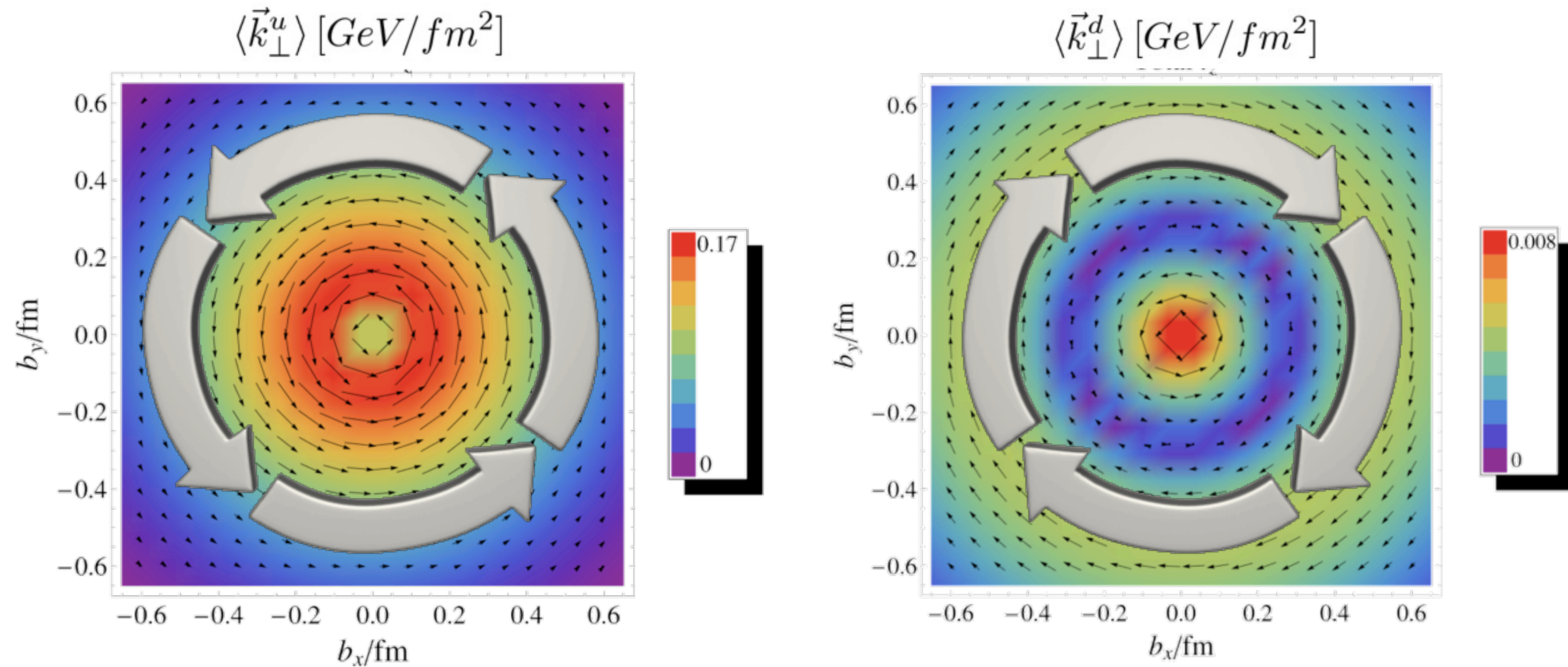
Results in a light-front constituent quark model:
Lorce', BP, Xiong, Yuan, PRD85 (2012)




Quark Orbital Angular Momentum

$$\mathcal{L}_z^q = \int dx d^2\vec{k}_\perp d^2\vec{b}_\perp (\vec{b}_\perp \times \vec{k}_\perp) \rho_{LU}^q(\vec{b}_\perp, \vec{k}_\perp, x)$$

$$= \int d^2\vec{b}_\perp \vec{b}_\perp \times \langle \vec{k}_\perp^q \rangle \longrightarrow \langle \vec{k}_\perp^q \rangle = \int dx d\vec{k}_\perp \vec{k}_\perp \rho_{LU}^q(\vec{b}_\perp, \vec{k}_\perp, x)$$

Lorce', BP (11)
 Hatta (12)
 Ji, Xiong, Yuan (12)



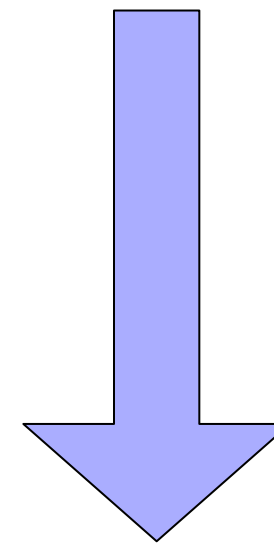
 Proton spin
 *u*-quark OAM
 *d*-quark OAM

Results in a light-front constituent quark model:
 Lorce', BP, Xiong, Yuan, PRD85 (2012)

$$\mathcal{L}_z^q = \int dx d^2\vec{k}_\perp d^2\vec{b}_\perp (\vec{b}_\perp \times \vec{k}_\perp) W(\vec{b}_\perp, \vec{k}_\perp, x)$$

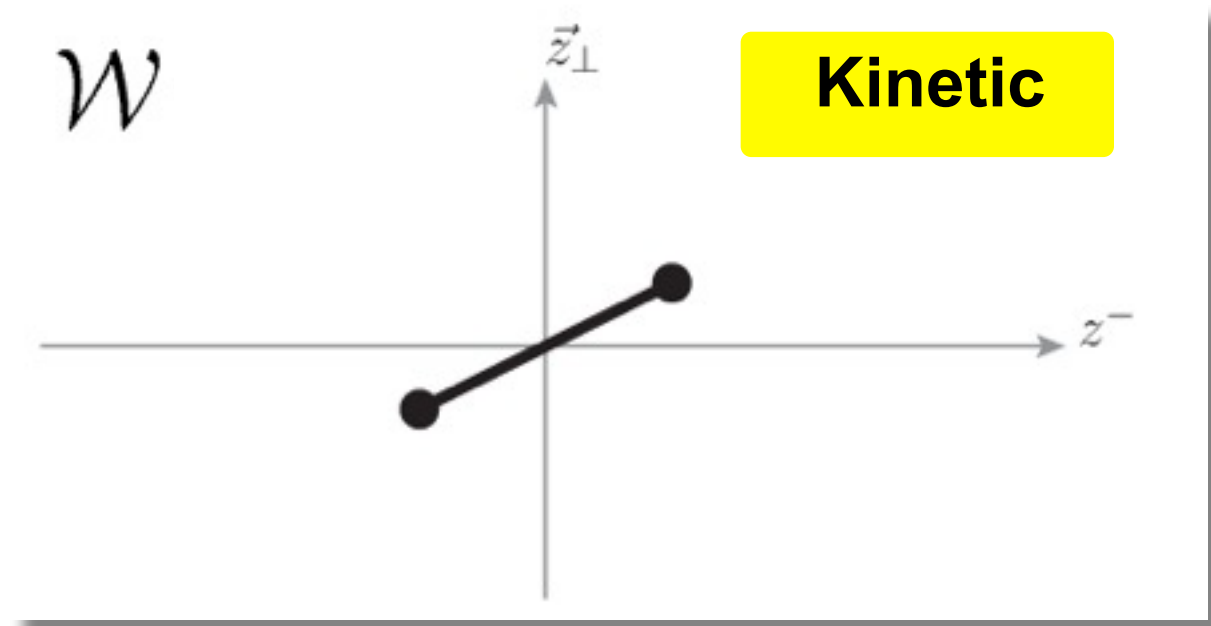
[C.L., Pasquini (2011)]
[C.L., Pasquini, Xiong, Yuan(2011)]

Light-cone gauge $A^+ = 0$
not gauge invariant, but with simple partonic interpretation

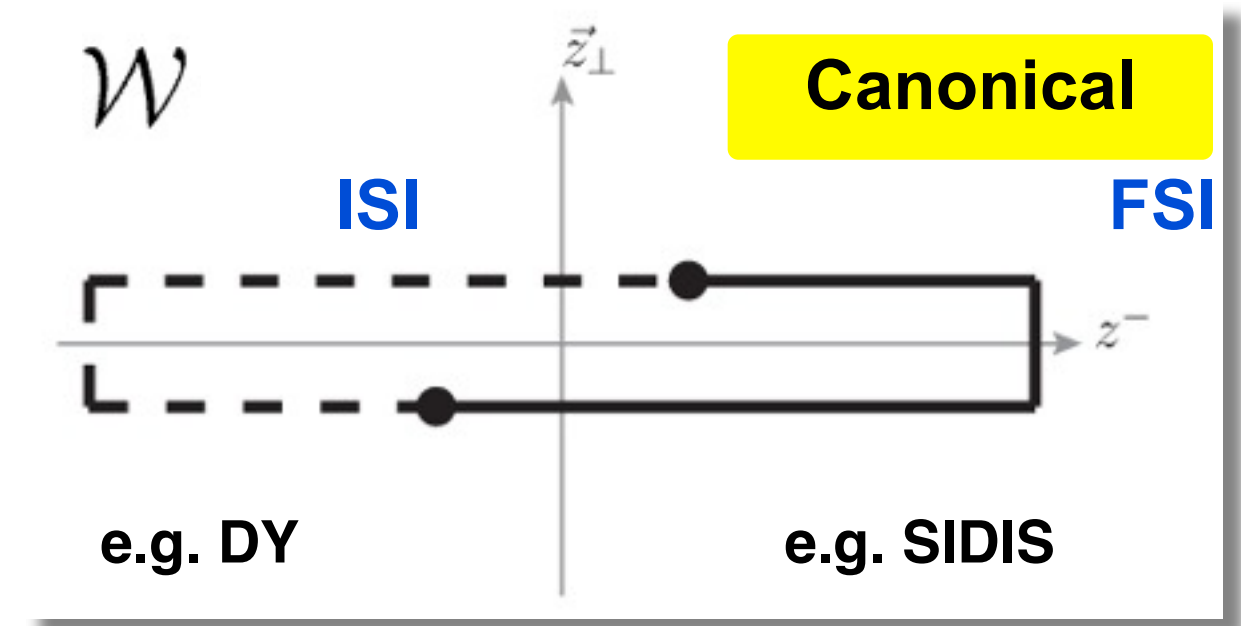


Gauge-invariant extension

$$W \rightarrow W^{\mathcal{W}}$$

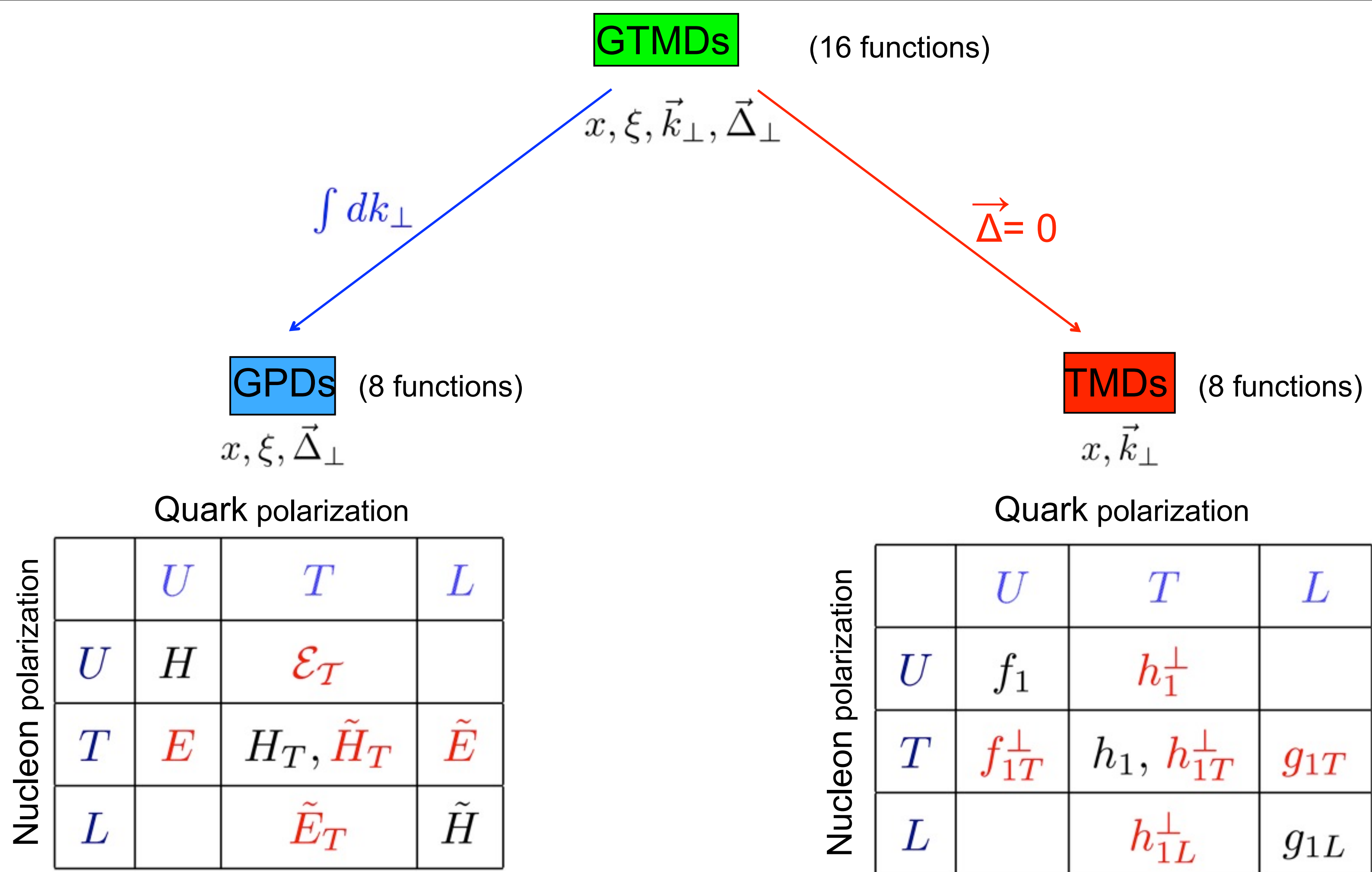


[Ji, Xiong, Yuan (2012)]
[Burkardt (2012)]



[Hatta (2012)]

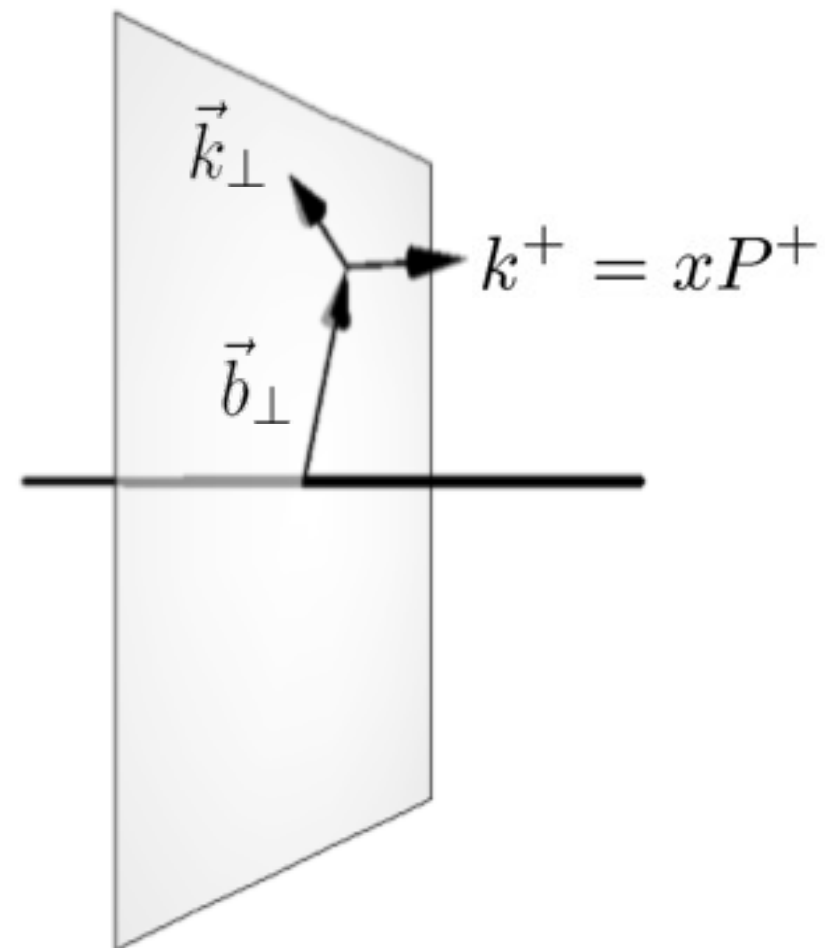
relations between the two gauge-invariant definitions
→ talk Burkardt



- ◆ almost all distributions (in red) vanish if there is no quark orbital angular momentum
- ◆ quark GPDs (at $\xi=0$) and TMDs given by the same overlap of LCWFs but in different kinematics
 - ⇒ each distribution contains unique information
 - ⇒ no model-independent relations between GPDs and TMDs

Wigner Distributions

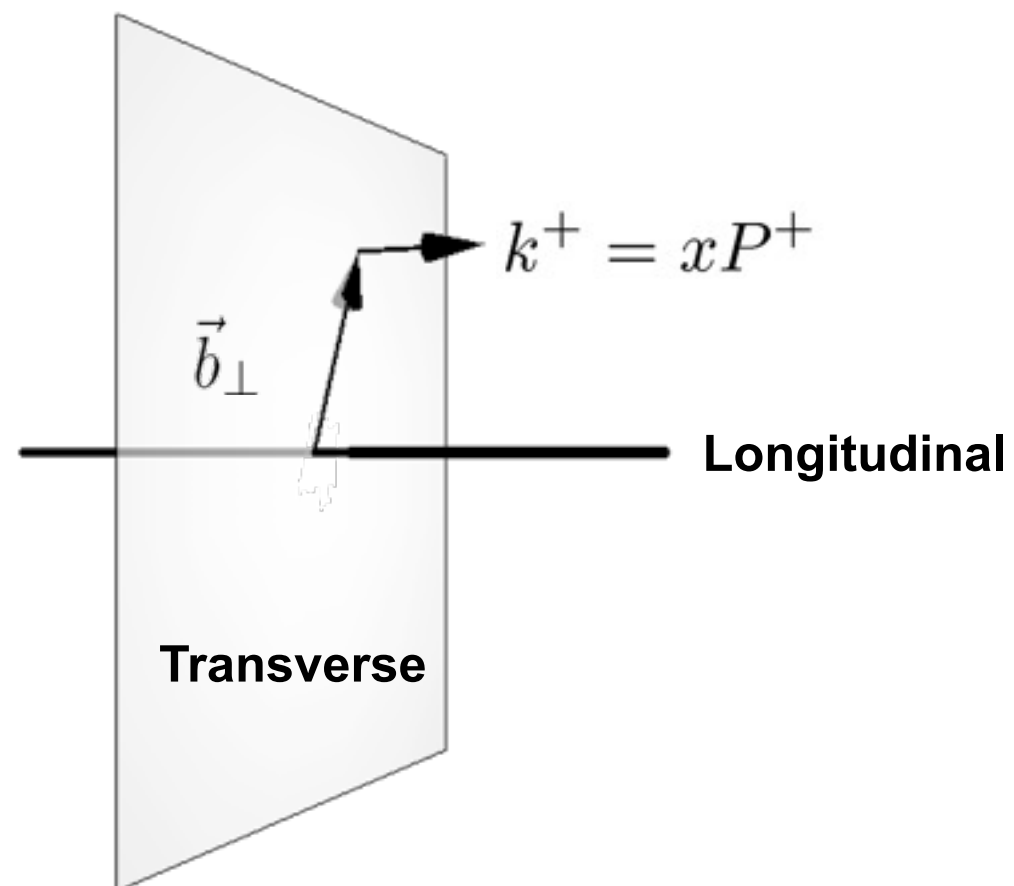
$$\rho(x, \vec{b}_\perp, \vec{k}_\perp)$$



quasi-probabilistic interpretation

GPDs

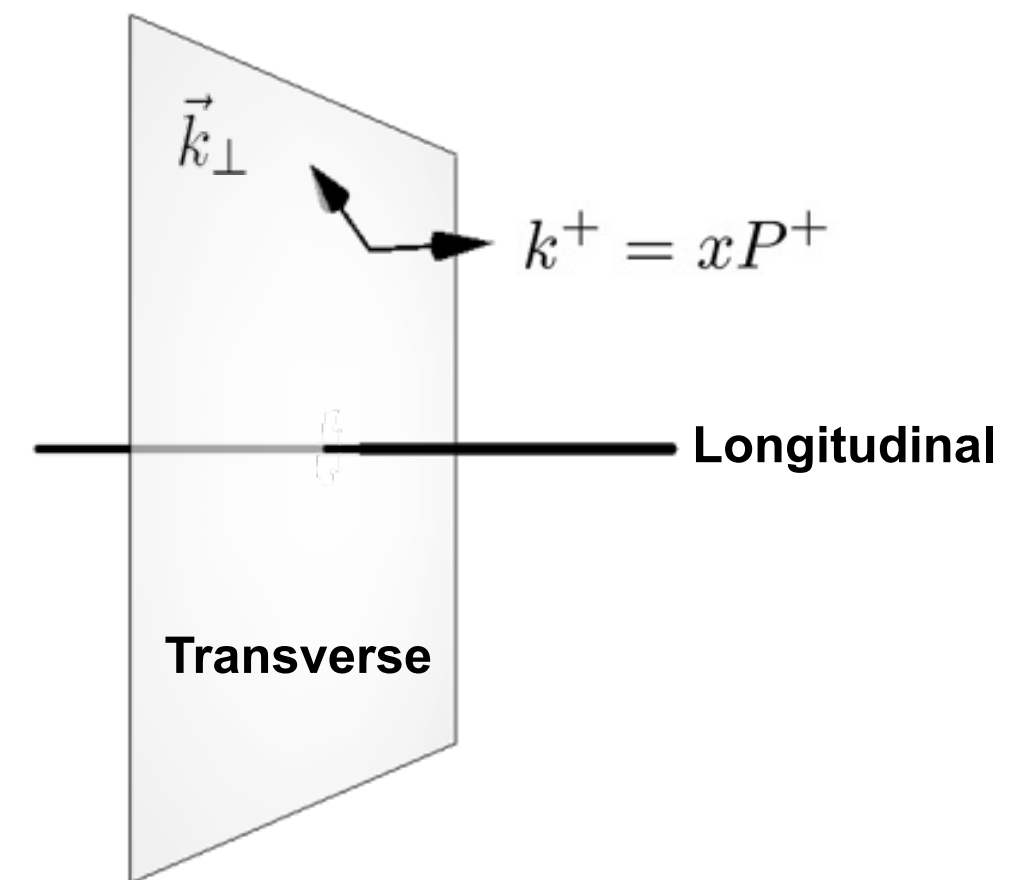
$$\rho(x, \vec{b}_\perp)$$



probabilistic interpretation

TMDs

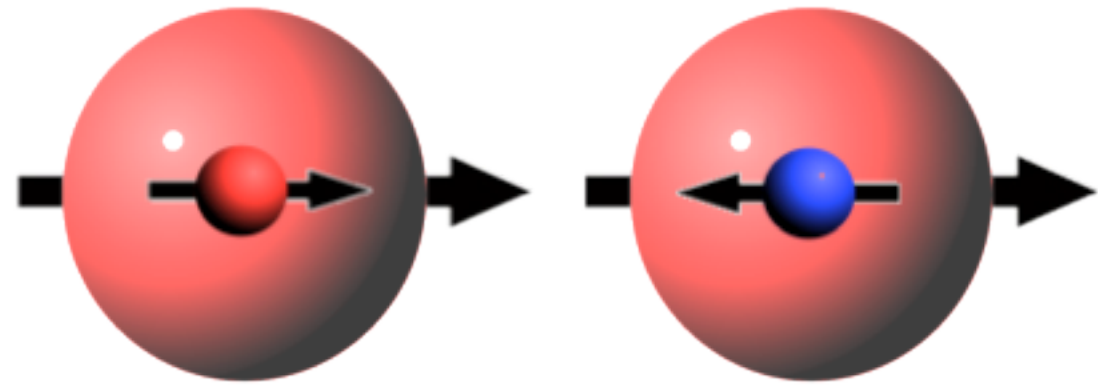
$$\rho(x, \vec{k}_\perp)$$



Longitudinal

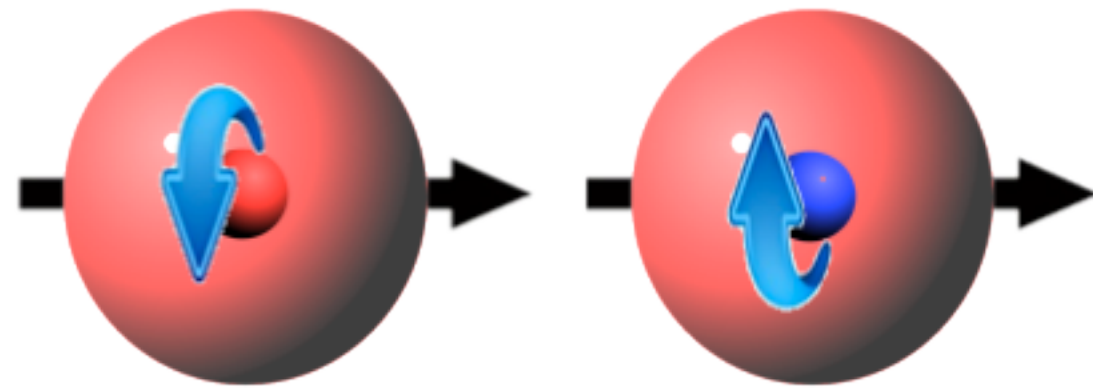
$$\vec{k}_\perp \quad \vec{b}_\perp$$

$$g_{1L}^q \leftrightarrow \tilde{\mathcal{H}}^q$$



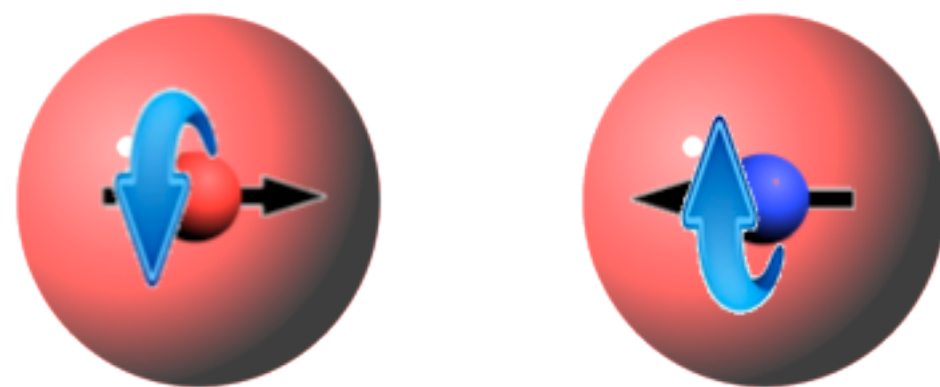
$$\vec{b}_\perp, \vec{k}_\perp$$

$$l_z^q \leftrightarrow \tilde{\mathcal{F}}_{14}^q$$



$$\vec{b}_\perp, \vec{k}_\perp$$

$$C_z^q \leftrightarrow \tilde{\mathcal{G}}_{11}^q$$

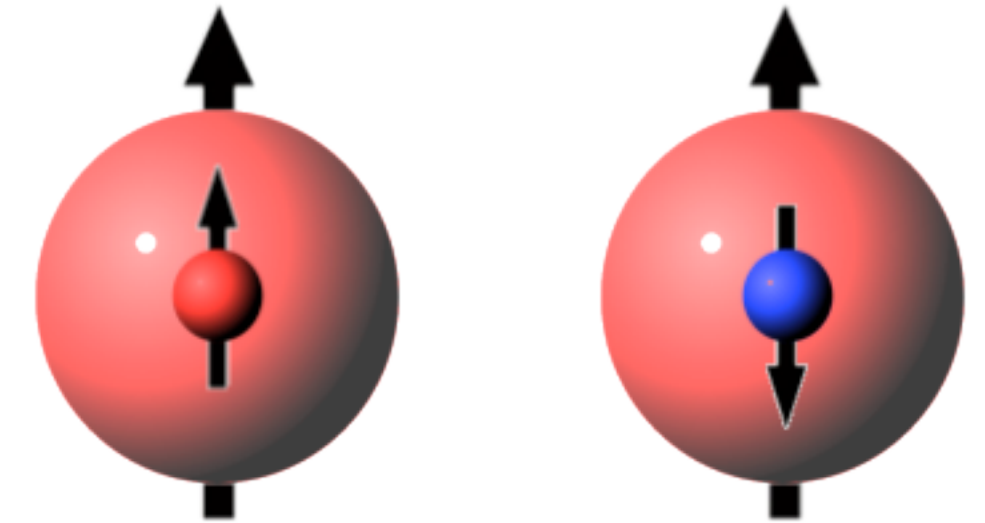


[C.L., Pasquini (2011)]

Transverse

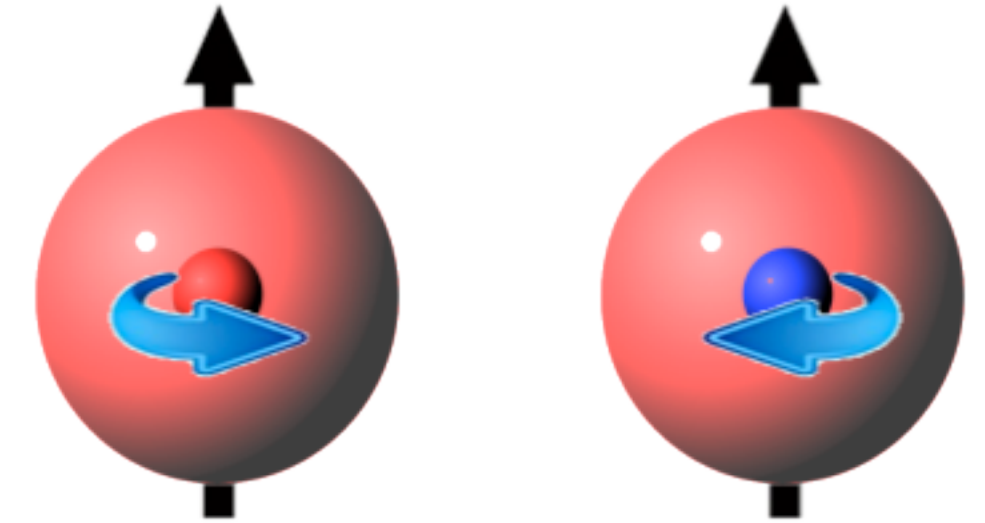
$$\vec{k}_\perp \quad \vec{b}_\perp$$

$$h_1^q \leftrightarrow \mathcal{H}_T^q$$



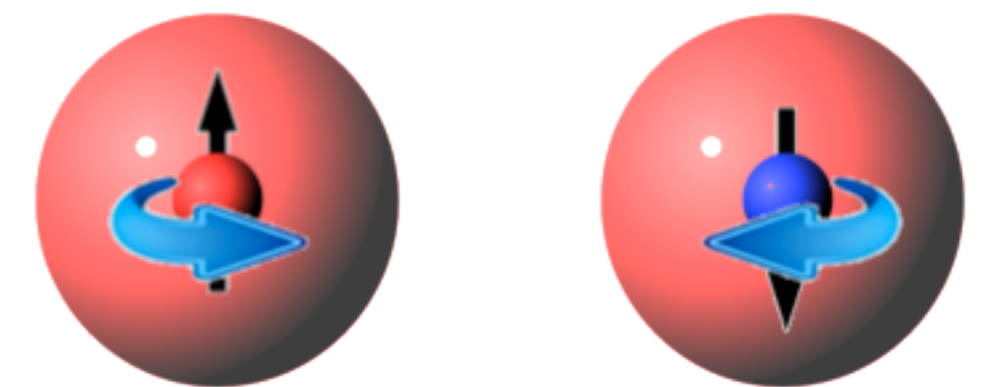
$$\vec{k}_\perp \quad \vec{b}_\perp$$

$$f_{1T}^{\perp q} \leftrightarrow \mathcal{E}^q$$



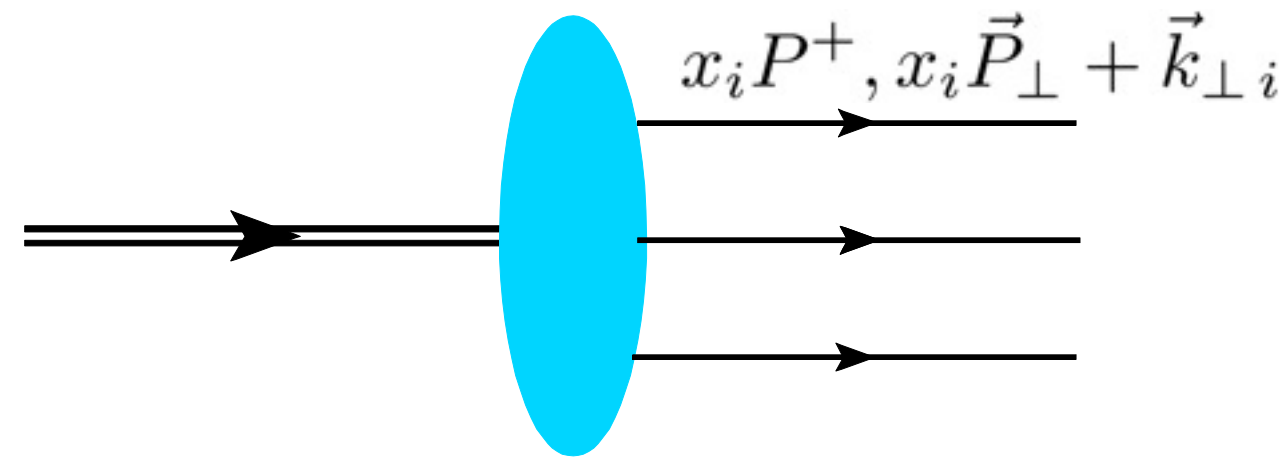
$$\vec{k}_\perp \quad \vec{b}_\perp$$

$$h_1^{\perp q} \leftrightarrow \mathcal{E}_T^q$$



[Burkardt (2005)]
[Barone et al. (2008)]

Quark OAM: Partial-Wave Decomposition



$$|P, \Lambda\rangle = \int d[1]d[2]d[3] \Psi_{\lambda_1 \lambda_2 \lambda_3}^\Lambda(x_i, \vec{k}_{\perp, i}) \frac{\epsilon^{ijk}}{\sqrt{6}} u_{i\lambda_1}^\dagger(1) u_{j\lambda_2}^\dagger(2) d_{k\lambda_3}^\dagger(3) |0\rangle$$

LCWF: eigenstate of OAM
(gauge $A^+=0 \rightarrow$ Jaffe-Manohar)

$$J_z^q \rightarrow (\uparrow\uparrow\uparrow)_{LC} = \frac{3}{2} \quad (\uparrow\uparrow\downarrow)_{LC} = \frac{1}{2} \quad (\uparrow\downarrow\downarrow)_{LC} = -\frac{1}{2} \quad (\downarrow\downarrow\downarrow)_{LC} = -\frac{3}{2}$$

$$L_z^q = \frac{1}{2} - J_z^q \rightarrow L_z^q = -1 \quad L_z^q = 0 \quad L_z^q = 1 \quad L_z^q = 2$$

$L_z \langle P, \uparrow | P, \uparrow \rangle^{L_z}$: probability to find the proton in a state with eigenvalue of OAM L_z

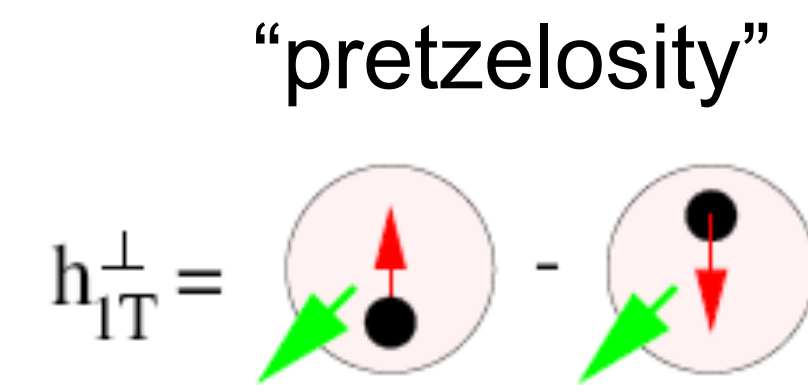
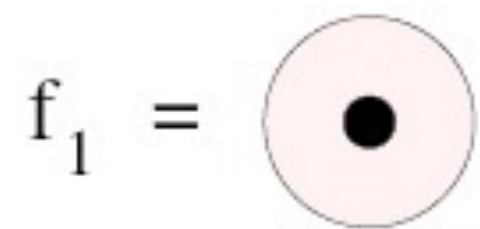
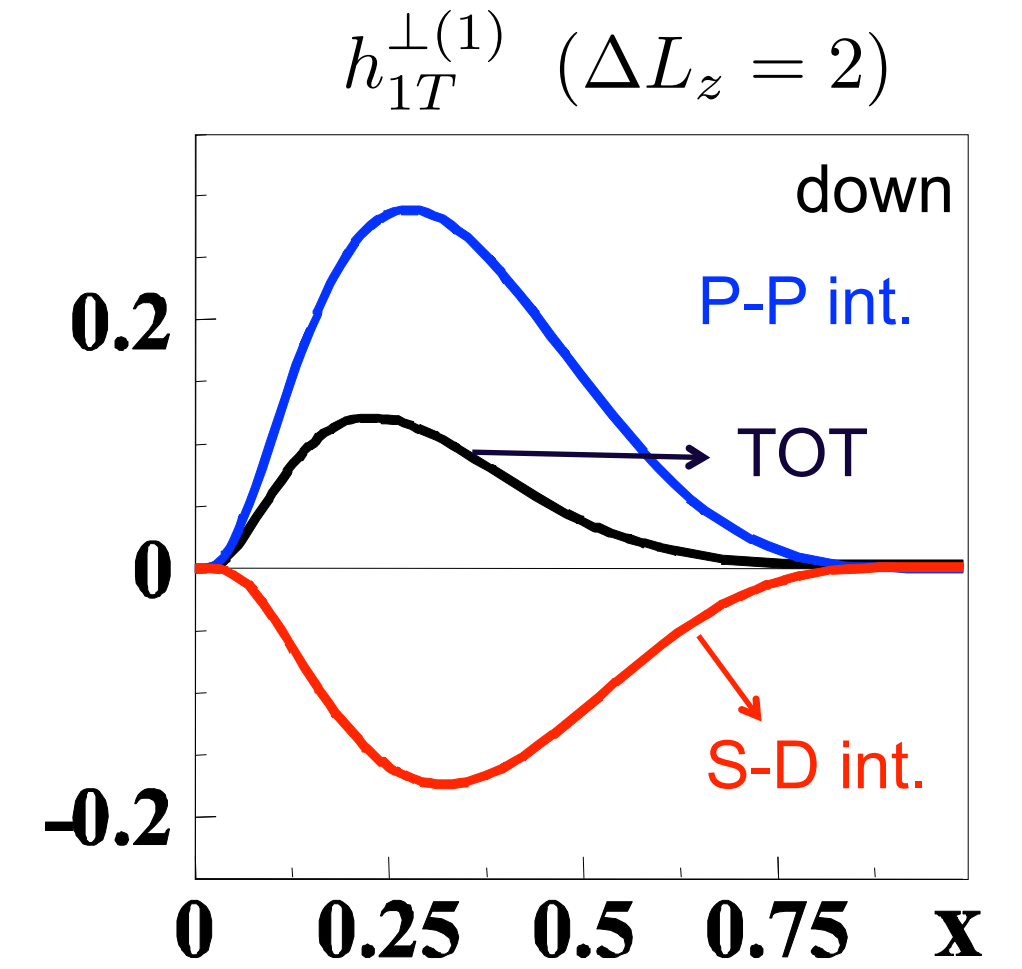
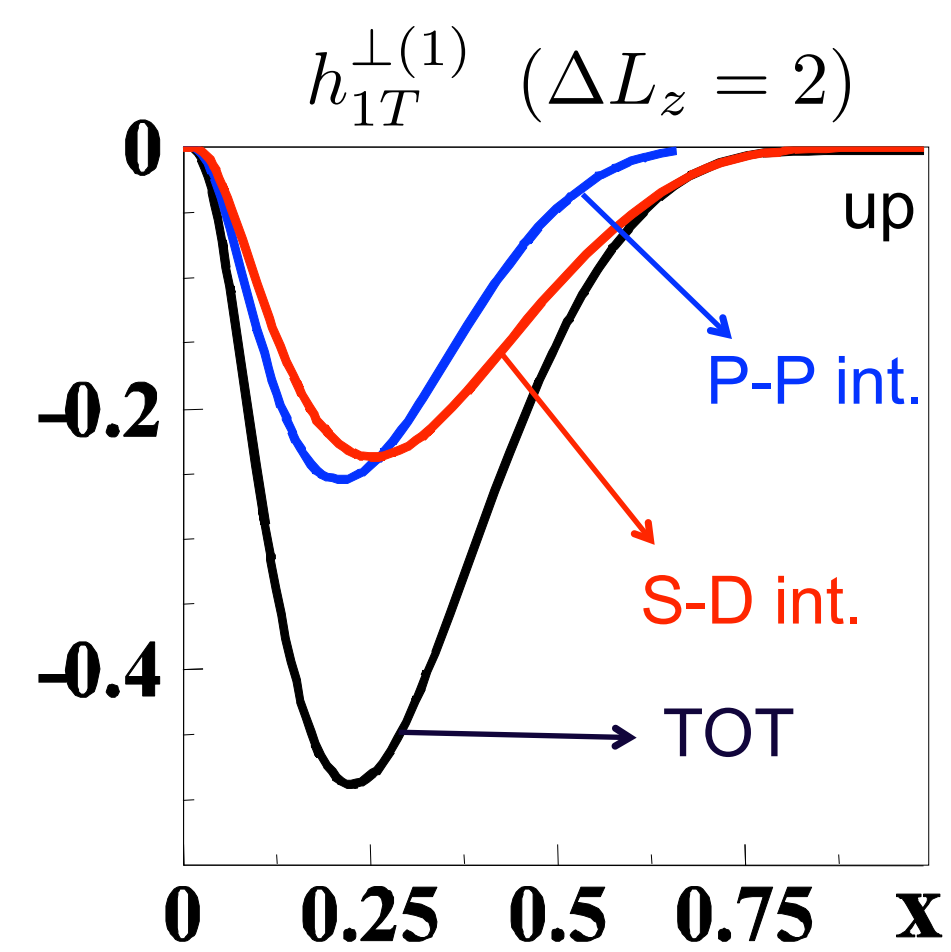
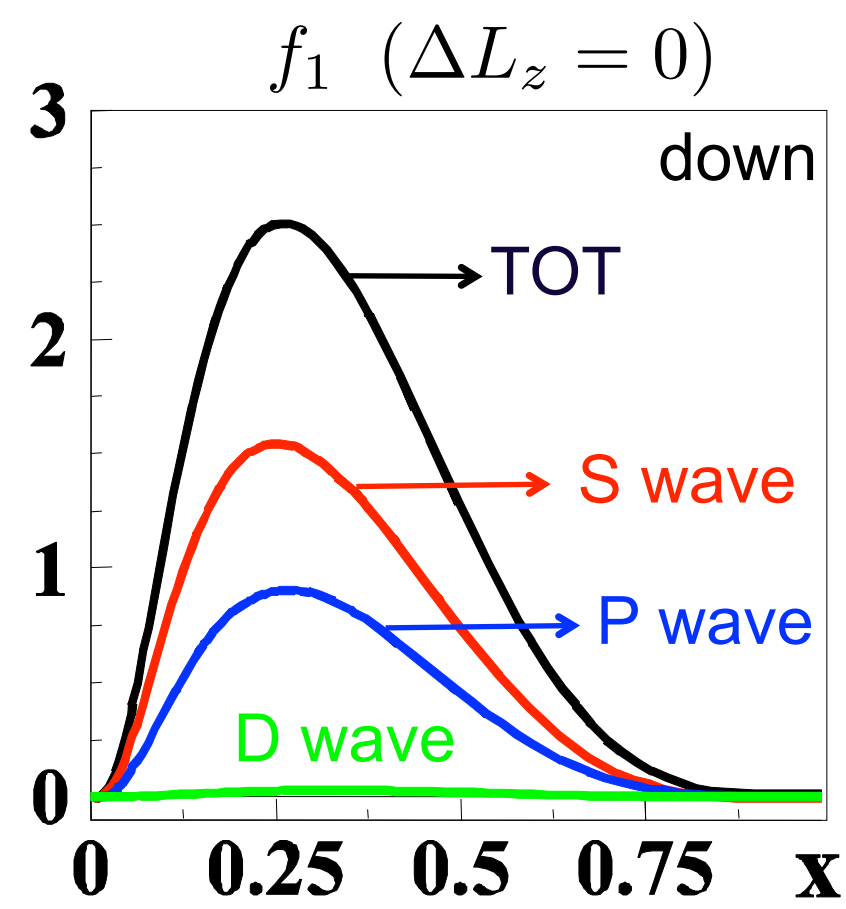
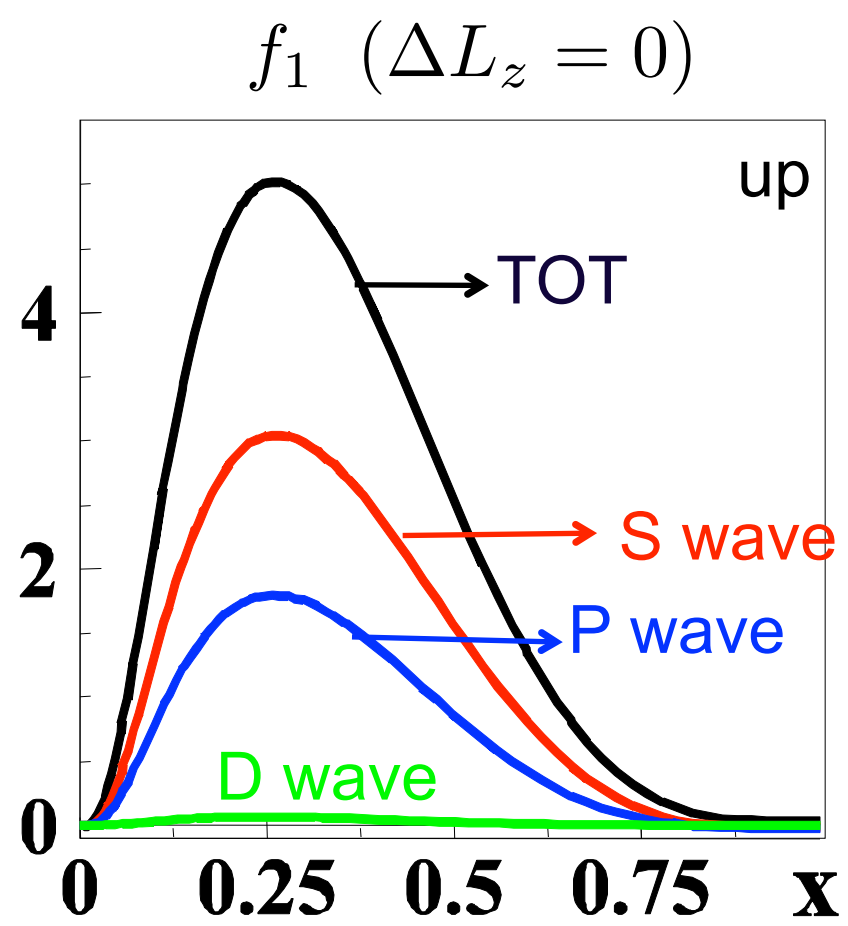


$$\mathcal{L}_z = \sum_{L_z} L_z \langle P, \uparrow | P, \uparrow \rangle^{L_z}$$

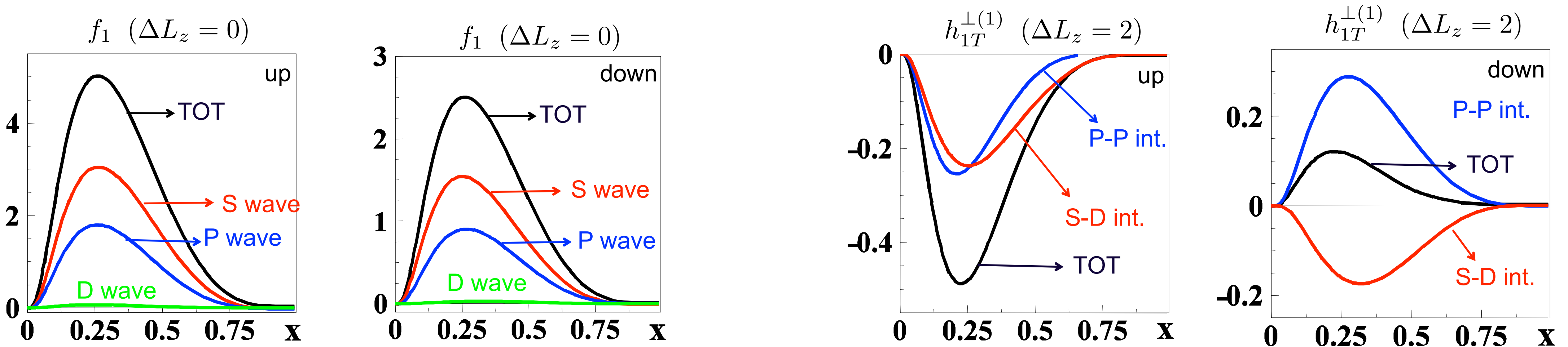


squared of LCWFs

◆ Orbital angular momentum content of TMDs (light-front constituent quark model)

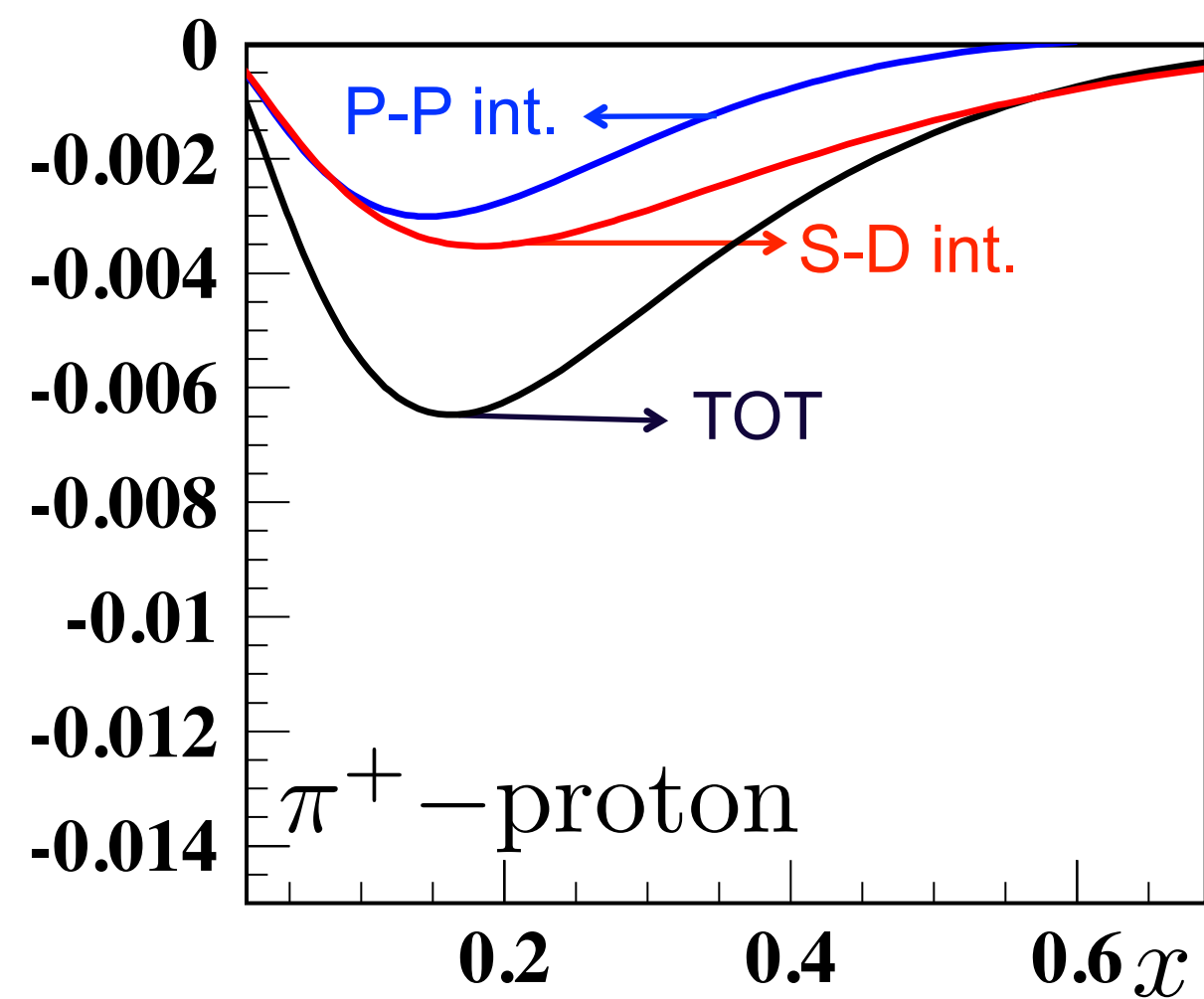


◆ Orbital angular momentum content of TMDs (light-front constituent quark model)

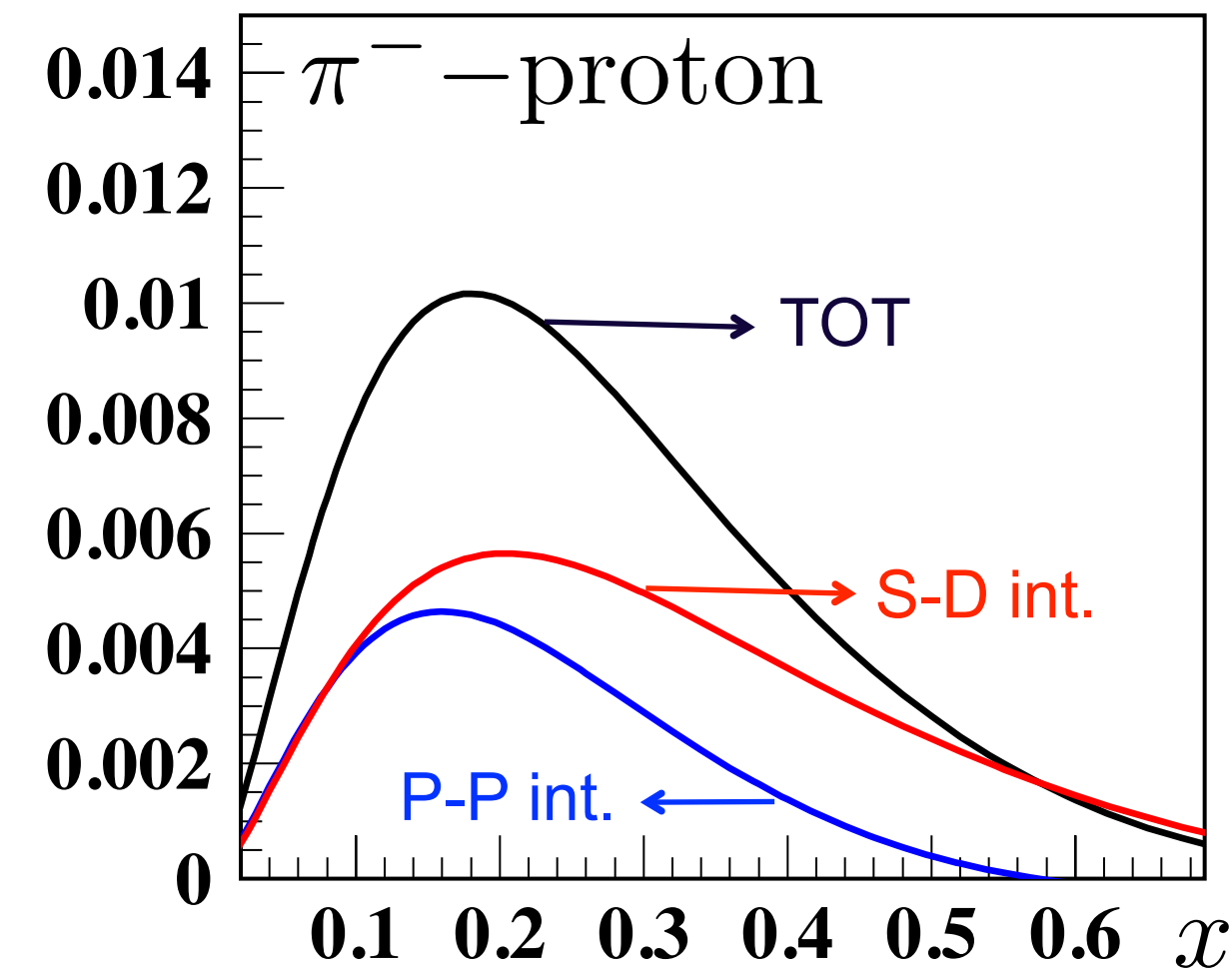


◆ Effects on SIDIS observables

$$A_{UT}^{\sin(3\phi - \phi_S)} \sim \frac{h_{1T}^{\perp} \otimes H_1}{f_1 \otimes D_1}$$



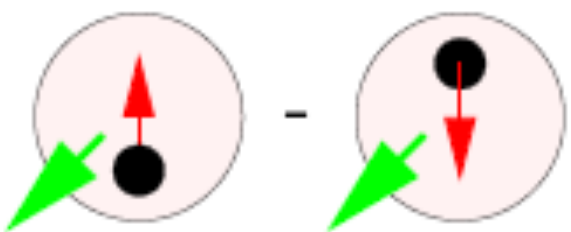
$\langle Q^2 \rangle = 2.5 \text{ GeV}^2$



Boffi, Efremov, BP, Schweitzer, PRD79(2009)

→ Efremov's lecture on TMDs

Quark OAM from Pretzelosity

$$h_{1T}^\perp = \text{[diagram]} - \text{[diagram]} \quad \text{“pretzelosity”}$$


model-dependent relation

$$\mathcal{L}_z = - \int dx d^2\vec{k}_\perp \frac{k_\perp^2}{2M^2} h_{1T}^\perp(x, k_\perp^2)$$

first derived in LC-diquark model and bag model

[She, Zhu, Ma, 2009; Avakian, Efremov, Schweitzer, Yuan, 2010]

Quark OAM from Pretzelosity

$$h_{1T}^\perp = \text{[diagram]} - \text{[diagram]} \quad \text{“pretzelosity”}$$

The diagram shows two circles representing quarks. The left circle has a red arrow pointing up and a black dot at the bottom. The right circle has a red arrow pointing down and a black dot at the top. A minus sign is between them.

model-dependent relation

$$\mathcal{L}_z = - \int dx d^2\vec{k}_\perp \frac{k_\perp^2}{2M^2} h_{1T}^\perp(x, k_\perp^2)$$

first derived in LC-diquark model and bag model

[She, Zhu, Ma, 2009; Avakian, Efremov, Schweitzer, Yuan, 2010]

\mathcal{L}_z
chiral even and charge even

h_{1T}^\perp
chiral odd and charge odd

$$\Delta L_z = 0$$

$$|\Delta L_z| = 2$$

no operator identity
relation at level of matrix elements of
operators

Quark OAM from Pretzelosity

$$h_{1T}^\perp = \text{[diagram: two circles with arrows]} \quad \text{“pretzelosity”}$$

model-dependent relation

$$\mathcal{L}_z = - \int dx d^2\vec{k}_\perp \frac{k_\perp^2}{2M^2} h_{1T}^\perp(x, k_\perp^2)$$

first derived in LC-diquark model and bag model

[She, Zhu, Ma, 2009; Avakian, Efremov, Schweitzer, Yuan, 2010]

\mathcal{L}_z	h_{1T}^\perp
chiral even and charge even	chiral odd and charge odd
$\Delta L_z = 0$	$ \Delta L_z = 2$

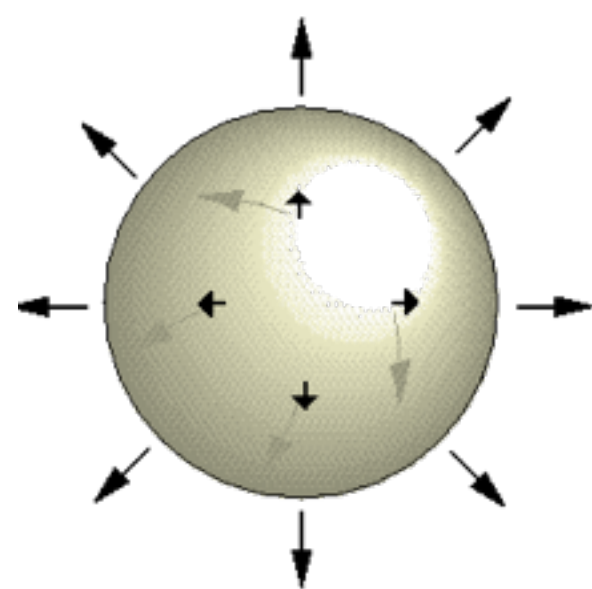
no operator identity
relation at level of matrix elements of
operators



valid in all quark models with spherical symmetry in the rest frame

[Lorce', BP, PLB (2012)]

→ Efremov's lecture on TMDs



Quark spin and OAM

GTMDs

Quark spin (from DIS)

$$S_z^q = \frac{1}{2} \int dx d^2k_\perp G_{14}^q(x, 0, \vec{k}_\perp, \vec{0}_\perp)$$

polarized PDF
inclusive DIS

$$\ell_z^q = - \int dx d^2k_\perp \frac{\vec{k}_\perp^2}{M^2} F_{14}^q(x, 0, \vec{k}_\perp, \vec{0}_\perp)$$

[Lorce, BP(2011)]
[Hatta (2011)]
[Lorce',BP, et al. (2012)]

$$\ell_{iz}^{\text{int}} = \vec{b}_{i\perp} \times \vec{k}_{i\perp}$$

TMDs

Quark spin

$$S_z^q = \frac{1}{2} \int dx d^2k_\perp g_{1L}^q(x, \vec{k}_\perp)$$

polarized PDF
inclusive DIS

$$\mathcal{L}_z^q(x, \vec{k}_\perp) = -\frac{\vec{k}_\perp^2}{2M^2} h_{1T}^{\perp q}(x, \vec{k}_\perp)$$

[Burkardt (2007)]
[Efremov et al. (2008,2010)]
[She, Zhu, Ma (2009)]
[Avakian et al. (2010)]
[Lorce', BP (2011)]



- Model-dependent
- Not intrinsic!

$$\mathcal{L}_{iz} = \vec{r}_{i\perp} \times \vec{k}_{i\perp}$$

Quark spin and OAM

GTMDs

Quark spin (from DIS)

$$S_z^q = \frac{1}{2} \int dx d^2k_\perp G_{14}^q(x, 0, \vec{k}_\perp, \vec{0}_\perp)$$

polarized PDF
inclusive DIS

$$\ell_z^q = - \int dx d^2k_\perp \frac{\vec{k}_\perp^2}{M^2} F_{14}^q(x, 0, \vec{k}_\perp, \vec{0}_\perp)$$

[Lorce, BP(2011)]
[Hatta (2011)]
[Lorce',BP, et al. (2012)]

$$\ell_{iz}^{\text{int}} = \vec{b}_{i\perp} \times \vec{k}_{i\perp}$$

TMDs

Quark spin

$$S_z^q = \frac{1}{2} \int dx d^2k_\perp g_{1L}^q(x, \vec{k}_\perp)$$

polarized PDF
inclusive DIS

$$\mathcal{L}_z^q(x, \vec{k}_\perp) = -\frac{\vec{k}_\perp^2}{2M^2} h_{1T}^{\perp q}(x, \vec{k}_\perp^2)$$

[Burkardt (2007)]
[Efremov et al. (2008,2010)]
[She, Zhu, Ma (2009)]
[Avakian et al. (2010)]
[Lorce', BP (2011)]



- Model-dependent
- Not intrinsic!

$$\mathcal{L}_{iz} = \vec{r}_{i\perp} \times \vec{k}_{i\perp}$$

GPDs

Quark spin (from DIS)

$$S_z^q = \frac{1}{2} \int dx \tilde{H}^q(x, 0, 0)$$

polarized PDF
inclusive DIS

Ji sum rule

$$J^q = \frac{1}{2} \int dx x [H^q(x, 0, 0) + E^q(x, 0, 0)]$$

$$L^q = J^q - S_z^q$$

[Ji (1997)]

Twist-3

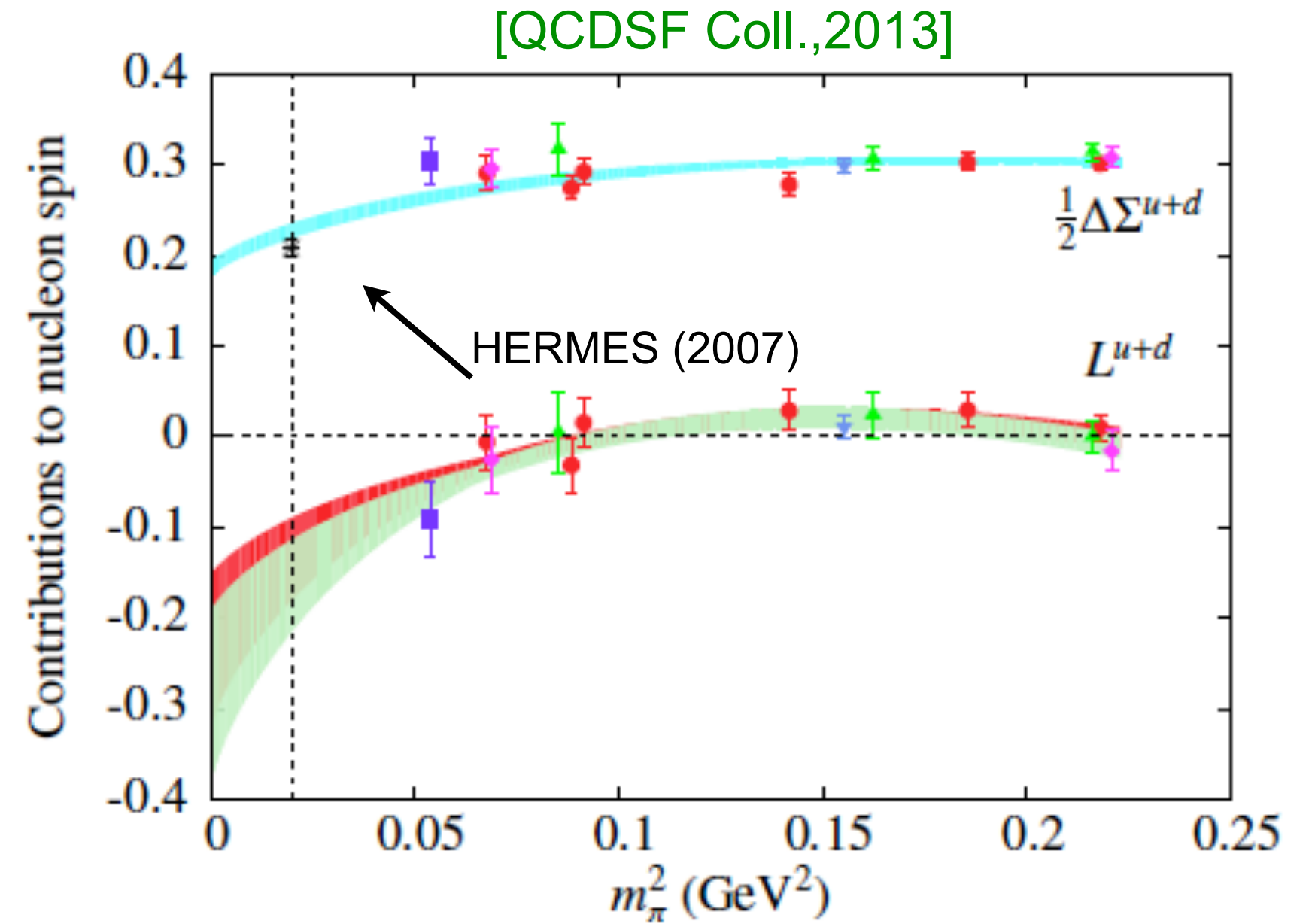
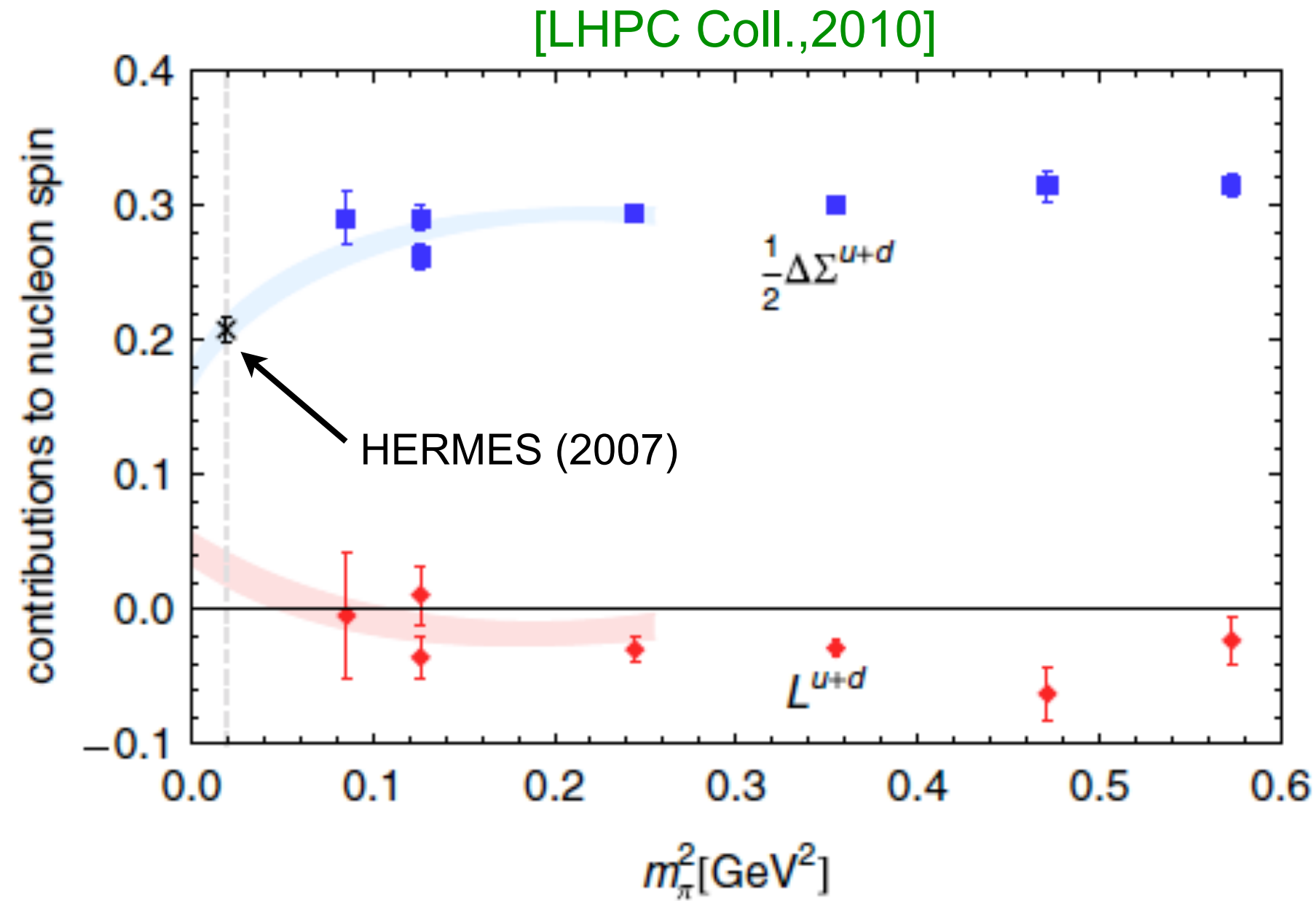
$$L_z^q + 2S_z^q = - \int dx x \tilde{E}_{2T}^q(x, 0, 0)$$

$$L_z^q = - \int dx x G_2^q(x, 0, 0)$$

Pure twist-3!

[Penttinen et al. (2000)]

Lattice Results



- “disconnected diagrams” not included
- Error bands: chiral extrapolation in m_π and extrapolation to $t=0$

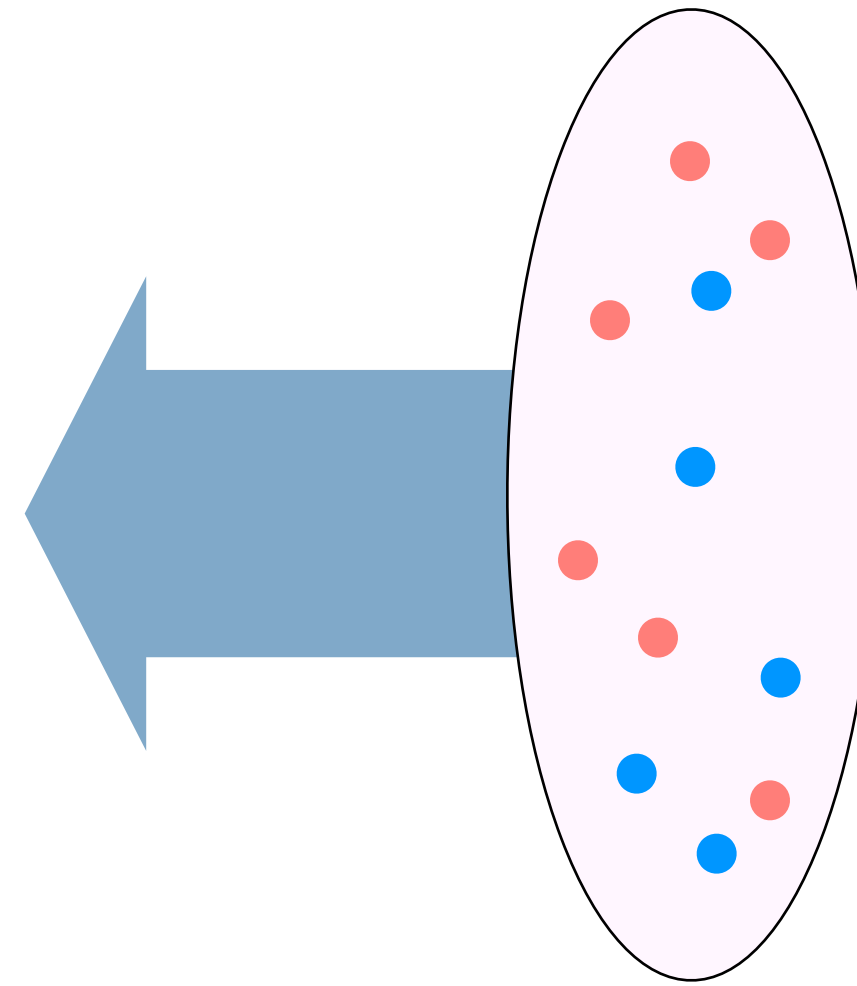
Lattice results ($\mu = 2 \text{ GeV}$)

$J^{u+d}=0.264(6)$	$\Delta \Sigma^{u+d} / 2 = 0.208(10)$	$L^{u+d} = 0.056(11)$	[LHPC Coll.,2010]
$J^{u+d}=0.168(42)$	$\Delta \Sigma^{u+d} / 2 = 0.225(8)$	$L^{u+d} = -0.141(27)$	[QCDSF Coll.,2013]

↓
cancellation between $L^u < 0$ and $L^d > 0$

Constraining quark OAM with Sivers function

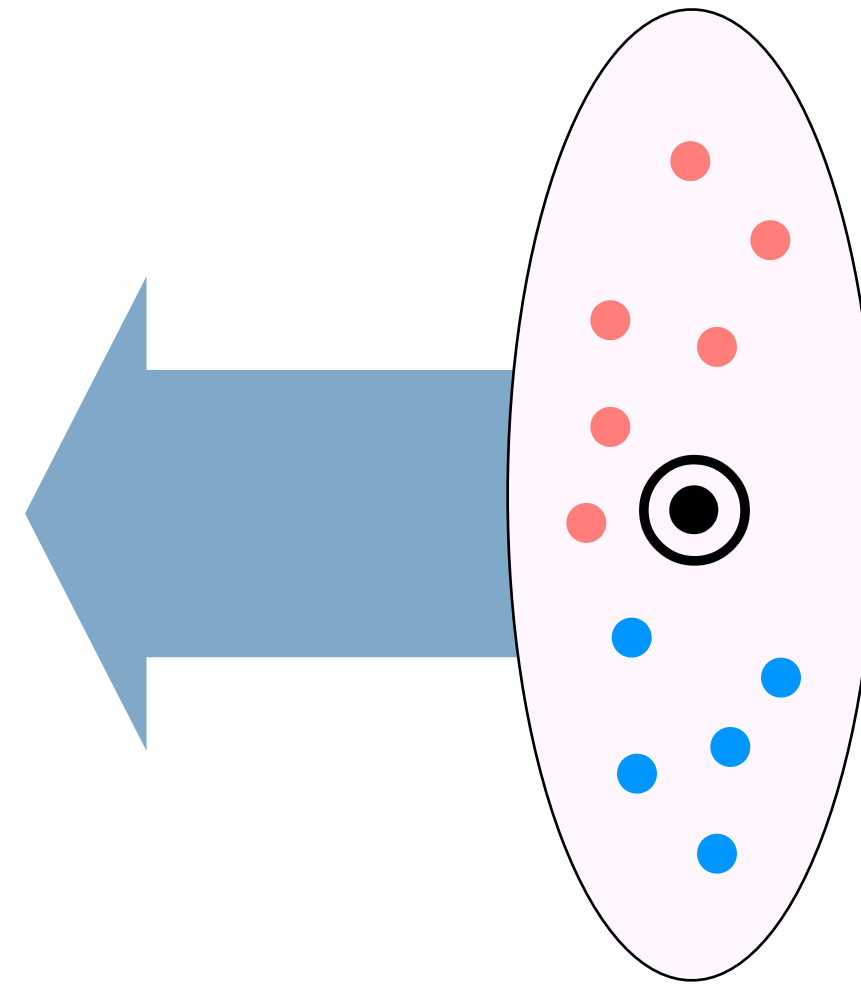
unpolarized quark in **unpolarized** nucleon



Burkardt, PRD66 (02)

Constraining quark OAM with Sivers function

unpolarized quark in **transversely** pol. nucleon

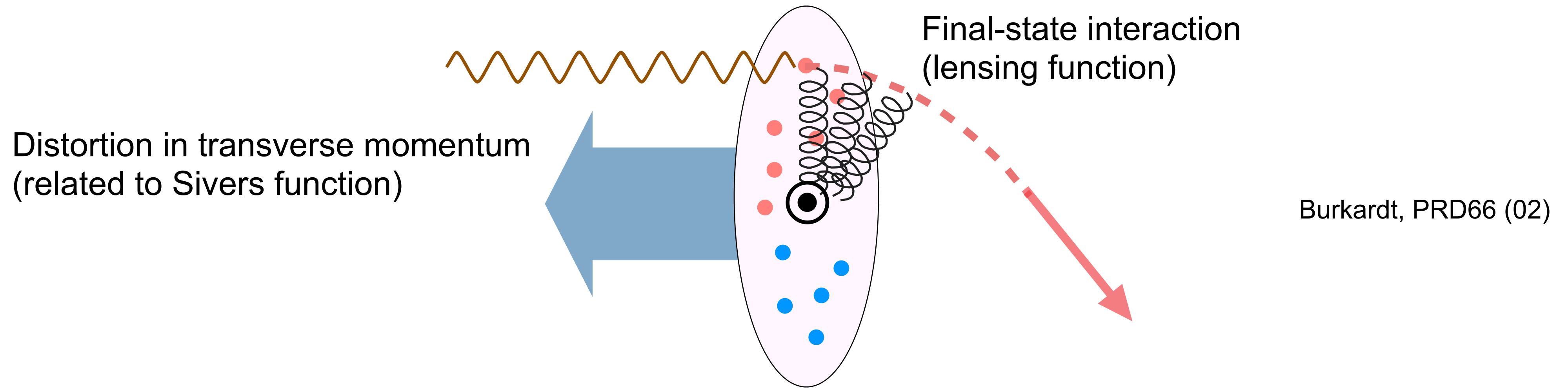


Distortion in impact parameter
(related to GPD E)

Burkardt, PRD66 (02)

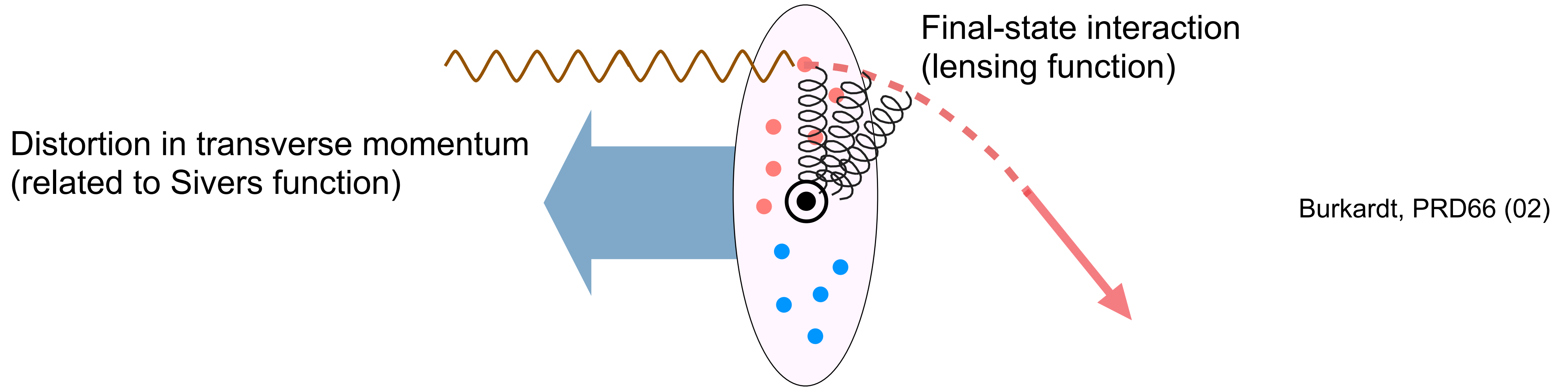
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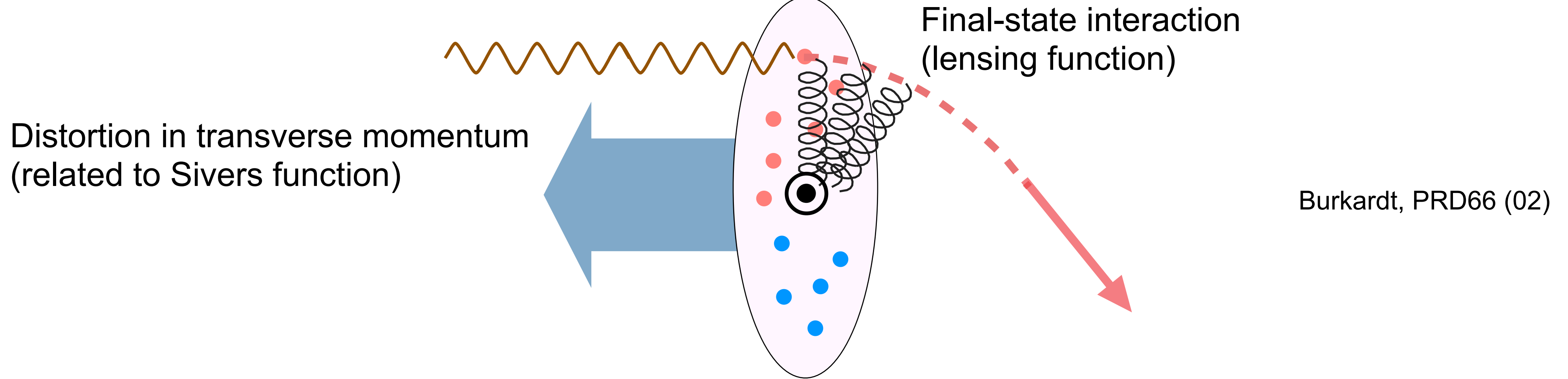
Burkardt, PRD66 (02)

$$- \int d^2 \vec{k}_T k_T^i \frac{\epsilon_T^{jk} k_T^j S_T^k}{M} f_{1T}^{\perp q}(x, \vec{k}_T^2) \simeq \int d^2 \vec{b}_T \mathcal{I}^{q,i}(x, \vec{b}_T) \frac{\epsilon_T^{jk} b_T^j S_T^k}{M} \left(\mathcal{E}^q(x, \vec{b}_T^2) \right)'$$

Sivers function
Lensing function
F.T. of E(x,0,t)

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Sivers function
Lensing function
F.T. of $E(x,0,t)$

inspired from model results

$$f_{1T}^{\perp(0)a}(x; Q_L^2) = -L(x) E^a(x, 0, 0; Q_L^2)$$

fitted to SIDIS data
(COMPASS, HERMES, JLab)

flavor independent
 $\frac{K}{(1-x)^\eta}$

first moment constrained
from anomalous magnetic moment

Bacchetta, Radici, PRL107(2011)

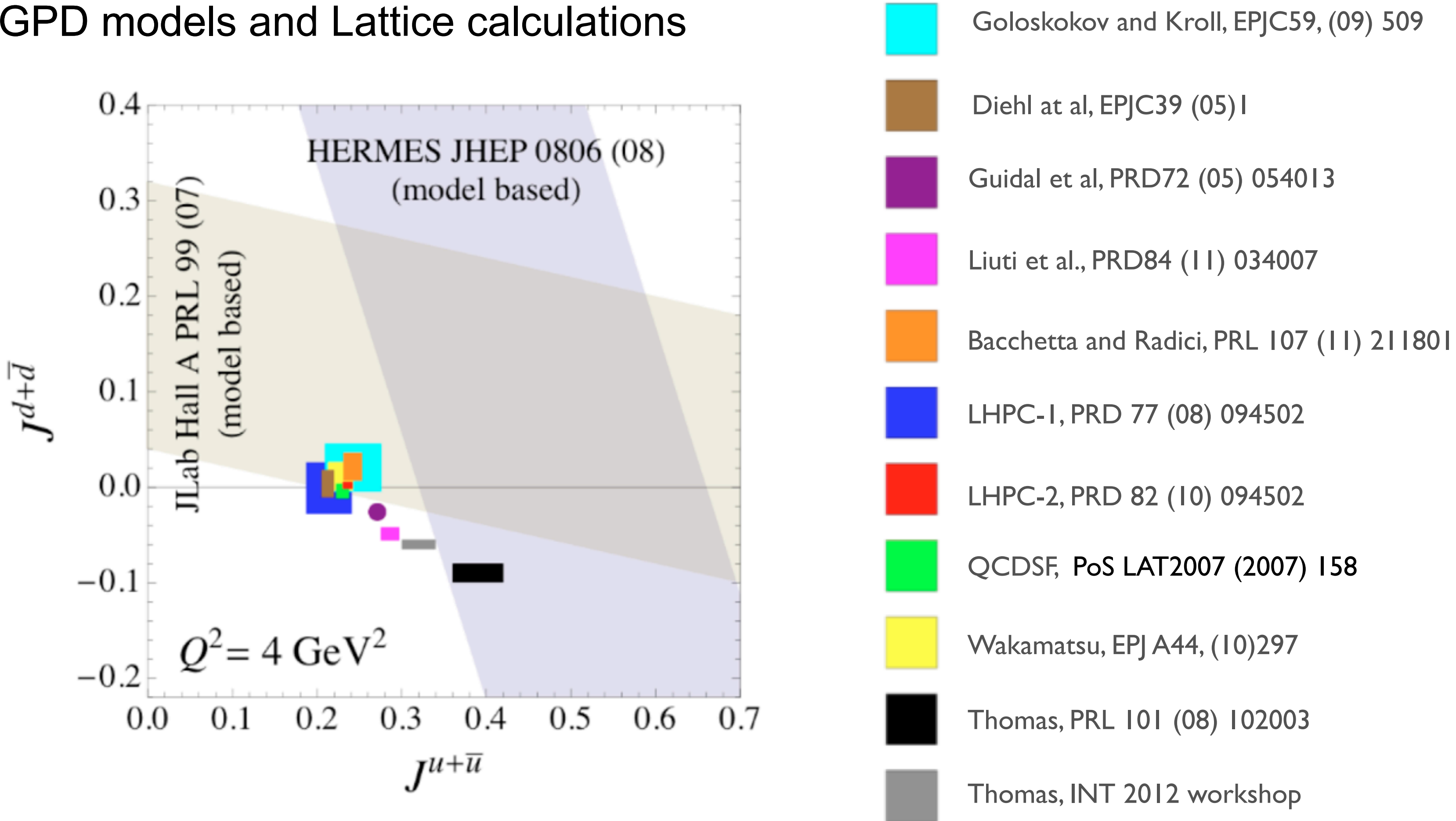
Results from **Sivers** ← **lensing** → **GPD**

Bacchetta, Radici, PRL107(2011)

$$\begin{aligned}
 J^u &= 0.229 \pm 0.002^{+0.008}_{-0.012}, & J^{\bar{u}} &= 0.015 \pm 0.003^{+0.001}_{-0.000}, \\
 J^d &= -0.007 \pm 0.003^{+0.020}_{-0.005}, & J^{\bar{d}} &= 0.022 \pm 0.005^{+0.001}_{-0.000}, \\
 J^s &= 0.006^{+0.002}_{-0.006}, & J^{\bar{s}} &= 0.006^{+0.000}_{-0.005}.
 \end{aligned}$$

($Q^2 = 4 \text{ GeV}^2$)

Comparing with GPD models and Lattice calculations



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❖ GTMDs Wigner Distributions

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❖ No direct connection between TMDs and OAM \Rightarrow need to use model-inspired connections

- use LCWF (eigenstate of quark OAM) to quantify amount of OAM in different observables
- model relation between pretzelosity and OAM
- OAM from model relation between Sivers function and GPD E