DETERMINING THE HIGGS SPIN AND PARITY USING GLUON POLARIZATION

Wilco den Dunnen

arXiv:1304.2654

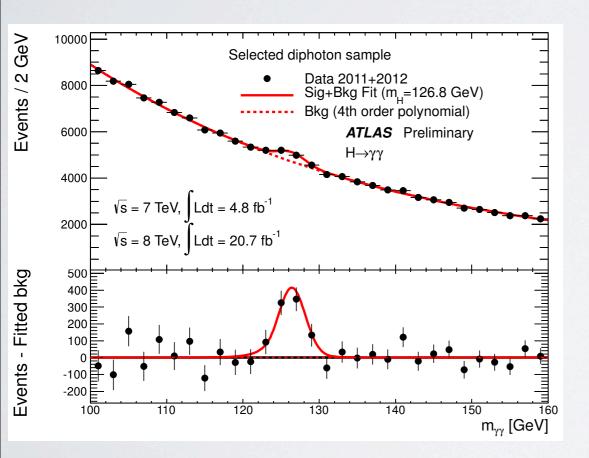
with Daniel Boer, Cristian Pisano and Marc Schlegel



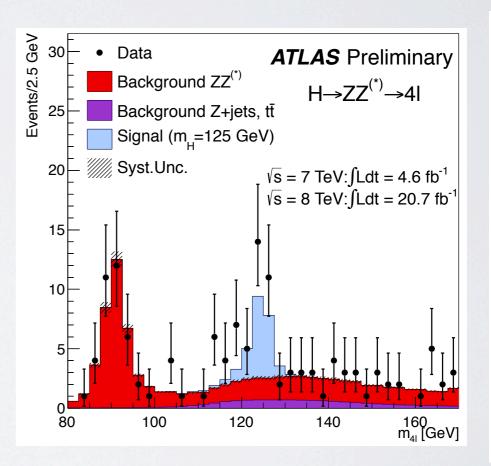


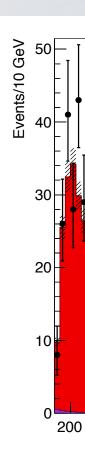
	postive parity	negative parity
spin-0	0+	0-
spin-I		
spin-2	2_m^+ 2_h^+ $2_{h'}^+$ $2_{h''}^+$	2_h^-

di-photon



$ZZ^* \rightarrow 4$



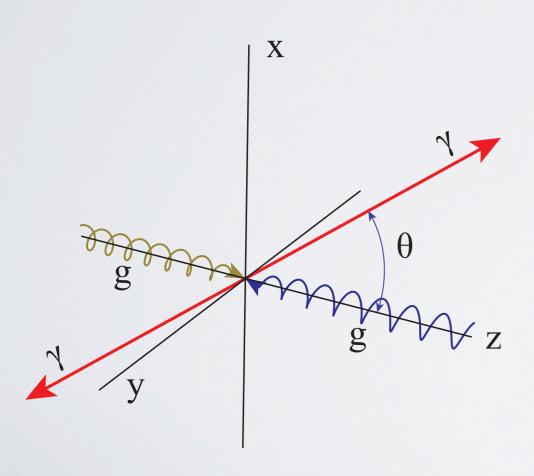


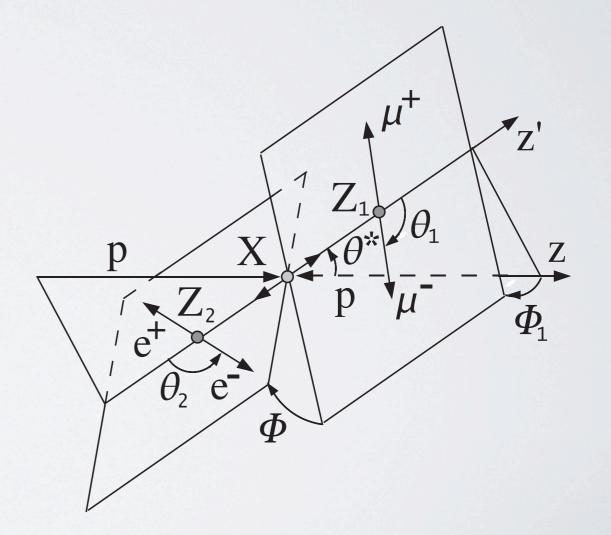
ATLAS-CONF-2013-012

ATLAS-CONF-2013-013

di-photon

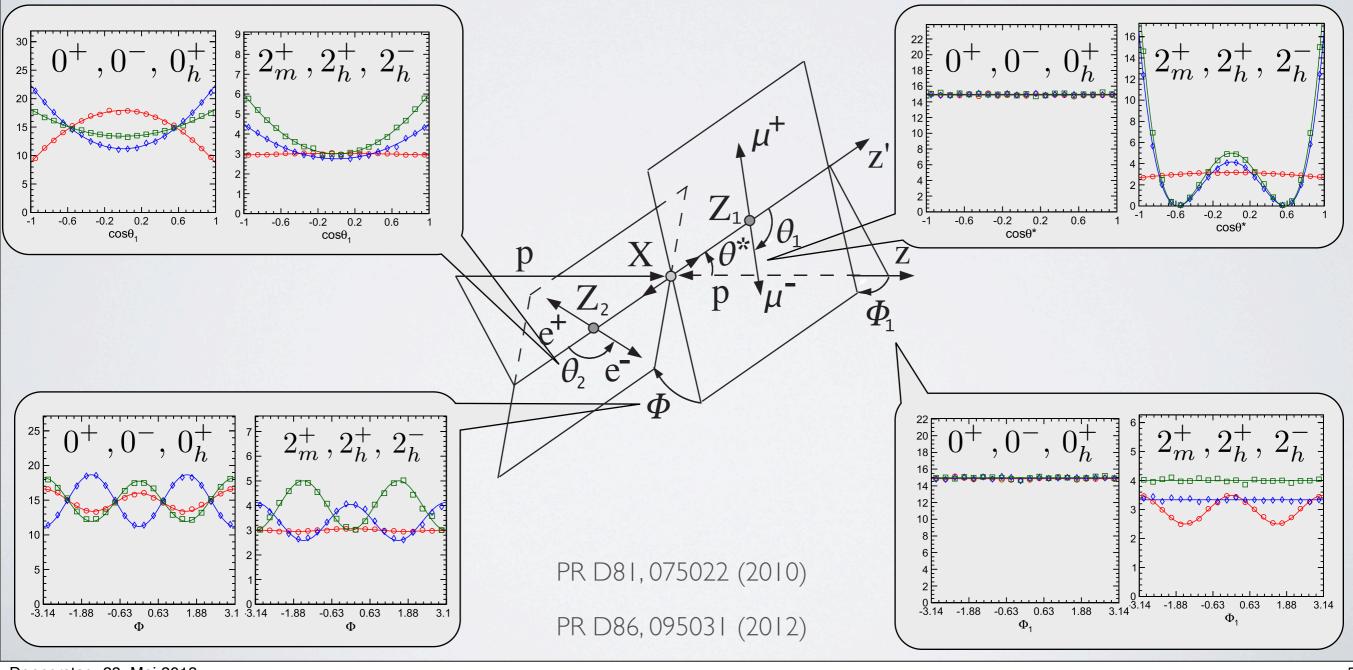




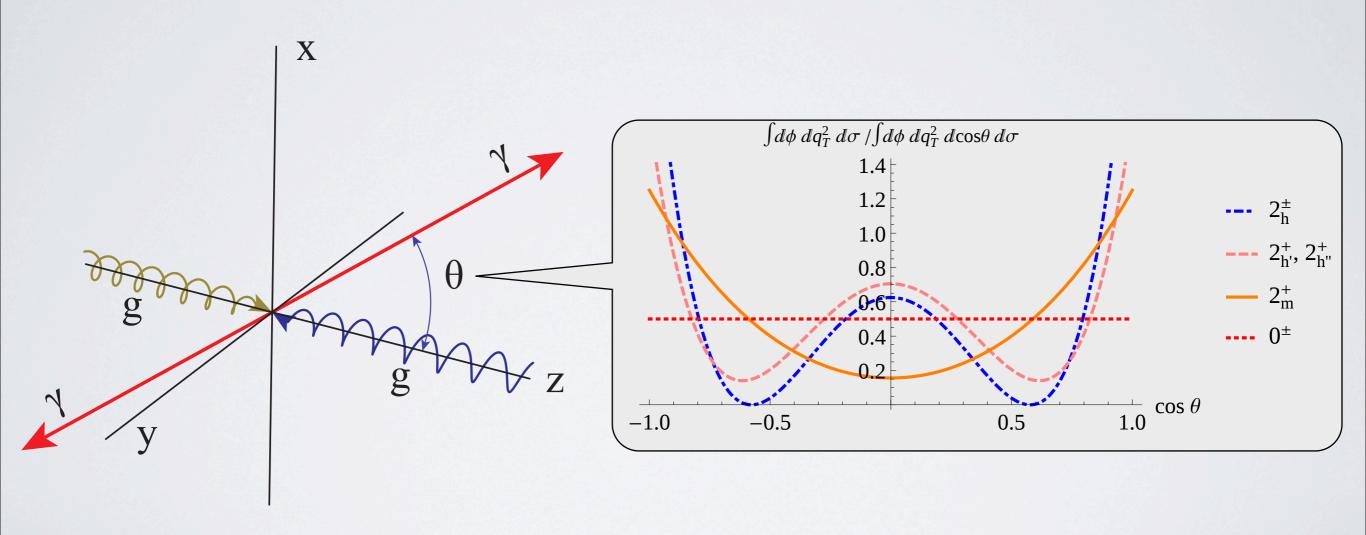


PR D81, 075022 (2010)

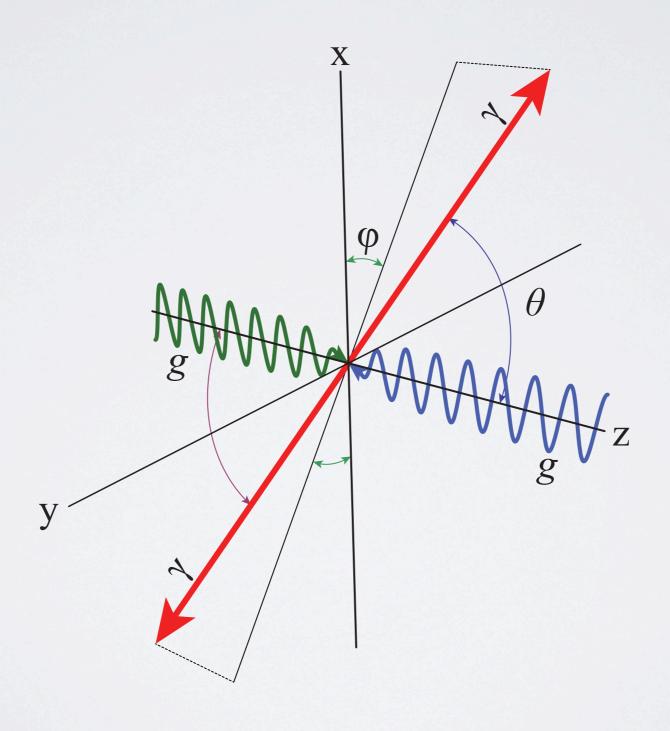
 $ZZ^* \rightarrow 4$



di-photon



WITH GLUON POLARIZATION



TMD FACTORIZATION

$$\frac{\mathrm{d}\sigma}{\mathrm{d}^4 q \mathrm{d}\Omega} \propto \int \!\! \mathrm{d}^2 \boldsymbol{p}_T \mathrm{d}^2 \boldsymbol{k}_T \delta^2 (\boldsymbol{p}_T + \boldsymbol{k}_T - \boldsymbol{q}_T) \mathcal{M}_{\mu\rho\kappa\lambda} \left(\mathcal{M}_{\nu\sigma}^{\ \kappa\lambda} \right)^* \Phi_g^{\mu\nu} (x_1, \boldsymbol{p}_T, \zeta_1, \mu) \Phi_g^{\rho\sigma} (x_2, \boldsymbol{k}_T, \zeta_2, \mu),$$

Hard scattering matrix element

Transverse Momentum Dependent (TMD) correlator

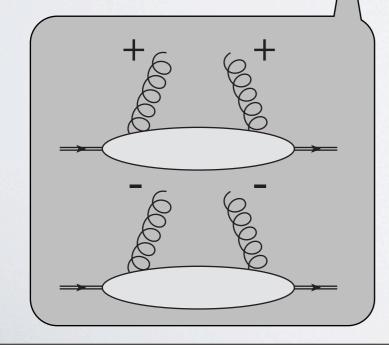
Ji, Ma, Yuan, JHEP07 (2005) 020

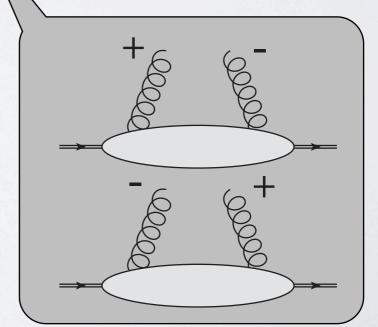
in principle also soft factor, but vanishes (up to NLO) for $\zeta_{1,2} \rightarrow 3/2\sqrt{S}$

TMD FACTORIZATION

$$\frac{\mathrm{d}\sigma}{\mathrm{d}^{4}q\mathrm{d}\Omega} \propto \int \!\!\mathrm{d}^{2}\boldsymbol{p}_{T}\mathrm{d}^{2}\boldsymbol{k}_{T}\delta^{2}(\boldsymbol{p}_{T}+\boldsymbol{k}_{T}-\boldsymbol{q}_{T})\mathcal{M}_{\mu\rho\kappa\lambda}\left(\mathcal{M}_{\nu\sigma}^{\kappa\lambda}\right)^{*}\Phi_{g}^{\mu\nu}(x_{1},\boldsymbol{p}_{T},\zeta_{1},\mu)\Phi_{g}^{\rho\sigma}(x_{2},\boldsymbol{k}_{T},\zeta_{2},\mu),$$

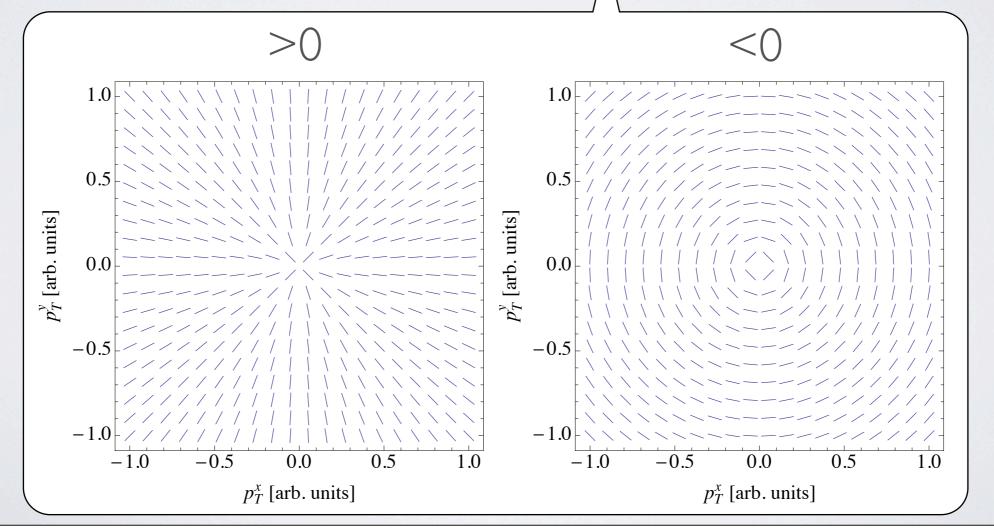
$$\begin{split} \Phi_g^{\mu\nu}(x, \boldsymbol{p}_T, \zeta, \mu) &\equiv 2 \int \frac{\mathrm{d}(\xi \cdot P) \, \mathrm{d}^2 \xi_T}{(xP \cdot n)^2 (2\pi)^3} \, e^{i(xP + p_T) \cdot \xi} \, \mathrm{Tr}_c \Big[\langle P | F^{n\nu}(0) \, \mathcal{U}_{[0,\xi]}^{n[-]} \, F^{n\mu}(\xi) \, \mathcal{U}_{[\xi,0]}^{n[-]} | P \rangle \Big]_{\xi \cdot P' = 0} \\ &= -\frac{1}{2x} \Big\{ g_T^{\mu\nu} f_1^g - \left(\frac{p_T^{\mu} p_T^{\nu}}{M_p^2} + g_T^{\mu\nu} \frac{\boldsymbol{p}_T^2}{2M_p^2} \right) h_1^{\perp g} \Big\} + \text{higher twist,} \end{split}$$



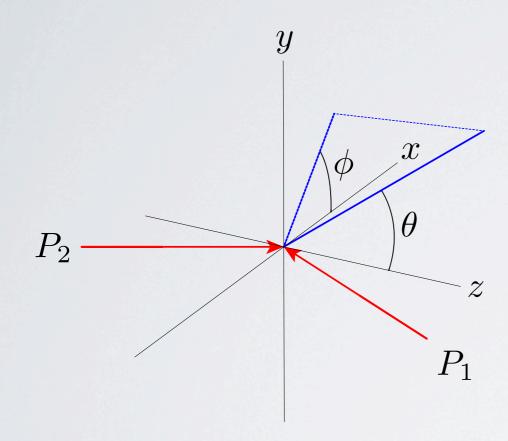


LINEARLY POLARIZED GLUON DISTRIBUTION

$$\begin{split} \Phi_g^{\mu\nu}(x, \boldsymbol{p}_T, \zeta, \mu) &\equiv 2 \int \frac{\mathrm{d}(\xi \cdot P) \, \mathrm{d}^2 \xi_T}{(xP \cdot n)^2 (2\pi)^3} \, e^{i(xP + p_T) \cdot \xi} \, \mathrm{Tr}_c \Big[\langle P | F^{n\nu}(0) \, \mathcal{U}_{[0,\xi]}^{n[-]} \, F^{n\mu}(\xi) \, \mathcal{U}_{[\xi,0]}^{n[-]} | P \rangle \Big]_{\xi \cdot P' = 0} \\ &= -\frac{1}{2x} \Big\{ g_T^{\mu\nu} f_1^g - \left(\frac{p_T^{\mu} p_T^{\nu}}{M_p^2} + g_T^{\mu\nu} \frac{\boldsymbol{p}_T^2}{2M_p^2} \right) h_1^{\perp g} \Big\} + \text{higher twist,} \end{split}$$



GENERAL STRUCTURE



Collins-Soper frame

$$\frac{\mathrm{d}\sigma}{\mathrm{d}^{4}q\mathrm{d}\Omega} \propto$$

$$F_{1}(Q,\theta) \,\mathcal{C}\left[f_{1}^{g}f_{1}^{g}\right] +$$

$$F_{2}(Q,\theta) \,\mathcal{C}\left[w_{2} \, h_{1}^{\perp g} h_{1}^{\perp g}\right] +$$

$$F_{3}(Q,\theta) \,\mathcal{C}\left[w_{3}(f_{1}^{g}h_{1}^{\perp g} + h_{1}^{\perp g}f_{1}^{g})\right] \cos(2\phi) +$$

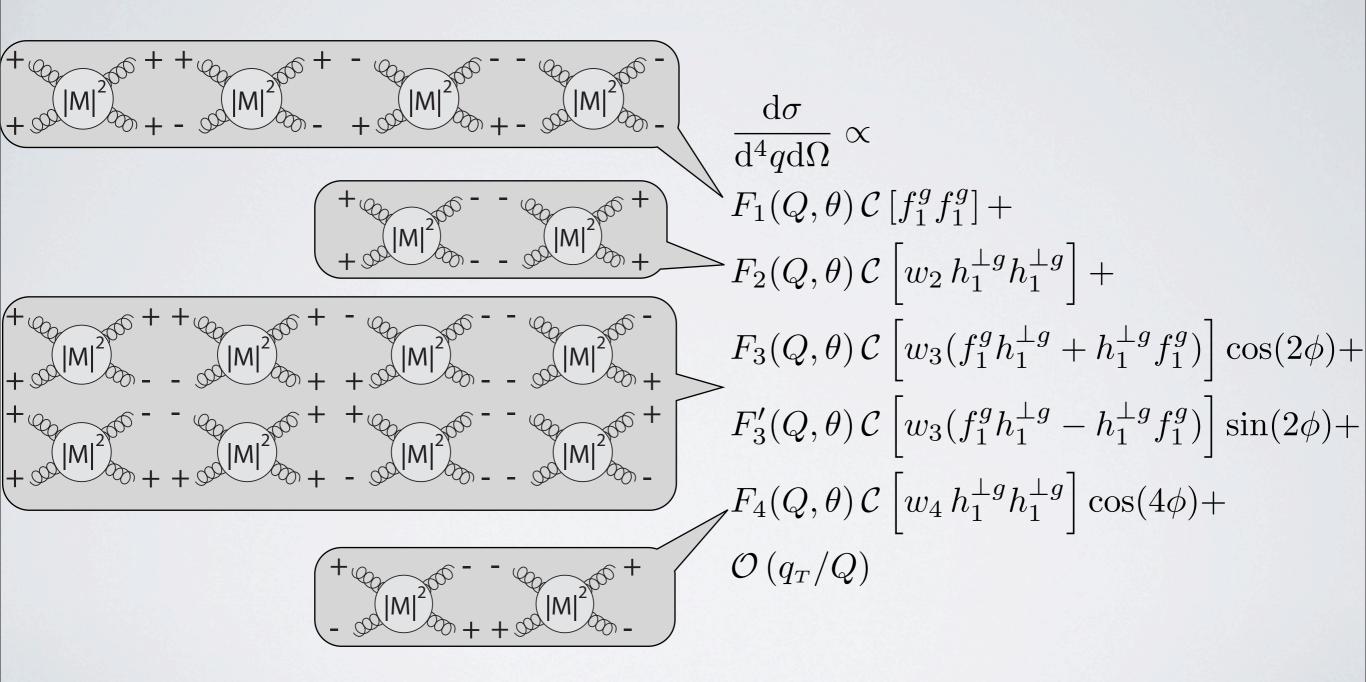
$$F_{3}'(Q,\theta) \,\mathcal{C}\left[w_{3}(f_{1}^{g}h_{1}^{\perp g} - h_{1}^{\perp g}f_{1}^{g})\right] \sin(2\phi) +$$

$$F_{4}(Q,\theta) \,\mathcal{C}\left[w_{4} \, h_{1}^{\perp g} h_{1}^{\perp g}\right] \cos(4\phi) +$$

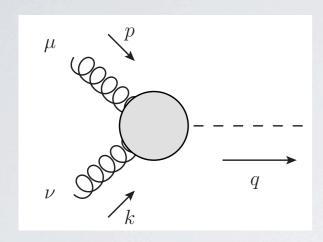
$$\mathcal{O}\left(q_{T}/Q\right)$$

ф azimuthal Collins-Soper angle

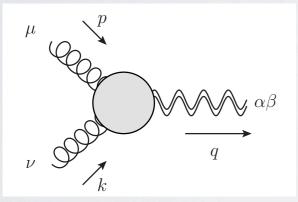
GENERAL STRUCTURE



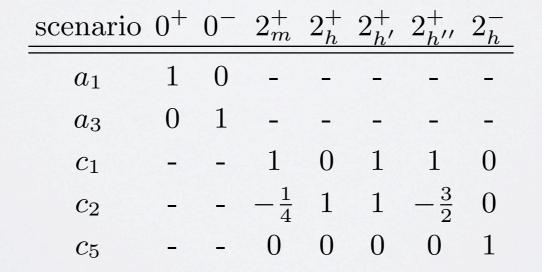
PARTONIC AMPLITUDES



$$= a_1 q^2 g^{\mu\nu} + a_3 \epsilon^{pk\mu\nu}$$



$$= \frac{1}{2}c_1 q^2 g^{\mu\alpha} g^{\nu\beta} + \left[c_2 q^2 g^{\mu\nu} + c_5 \epsilon^{pk\mu\nu}\right] \frac{(p-k)^{\alpha} (p-k)^{\beta}}{q^2}$$



0.0015

GENERAL STRUCTURE

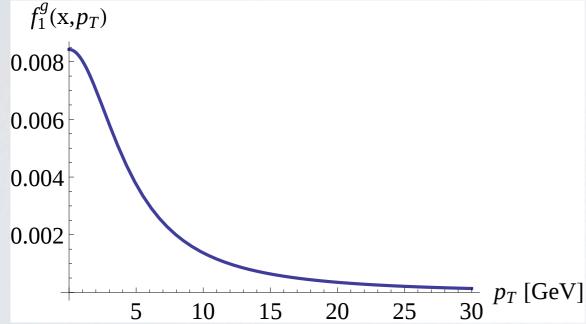
$$\frac{\mathrm{d}\sigma}{\mathrm{d}^{4}q\mathrm{d}\Omega} \propto \\ F_{1}(Q,\theta) \,\mathcal{C}\left[f_{1}^{g}f_{1}^{g}\right] + \\ F_{2}(Q,\theta) \,\mathcal{C}\left[w_{2} \, h_{1}^{\perp g} h_{1}^{\perp g}\right] + \\ F_{3}(Q,\theta) \,\mathcal{C}\left[w_{3}(f_{1}^{g}h_{1}^{\perp g} + h_{1}^{\perp g}f_{1}^{g})\right] \cos(2\phi) + \\ F_{3}'(Q,\theta) \,\mathcal{C}\left[w_{3}(f_{1}^{g}h_{1}^{\perp g} - h_{1}^{\perp g}f_{1}^{g})\right] \sin(2\phi) + \\ F_{4}(Q,\theta) \,\mathcal{C}\left[w_{4} \, h_{1}^{\perp g} h_{1}^{\perp g}\right] \cos(4\phi) + \\ \mathcal{O}\left(q_{T}/Q\right) \end{aligned}$$

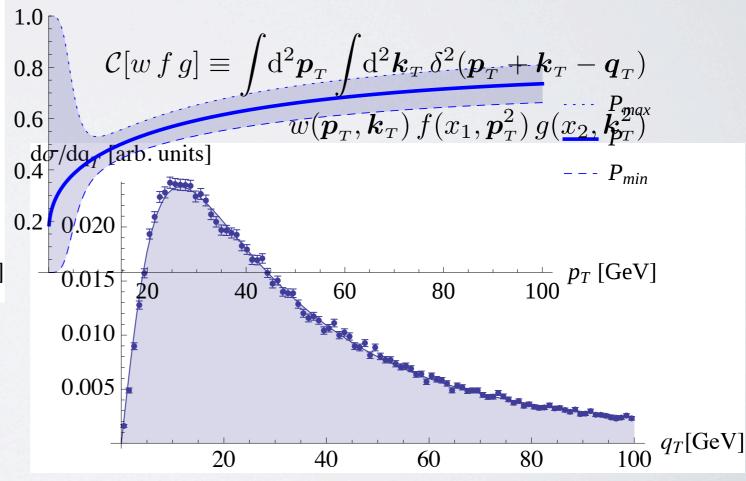
- Unpolarized distribution
- gluon TMD
- f_1^g

- Polarized distribution
- linearly polarized gluon distribution

 $-h_1^{\perp g}$

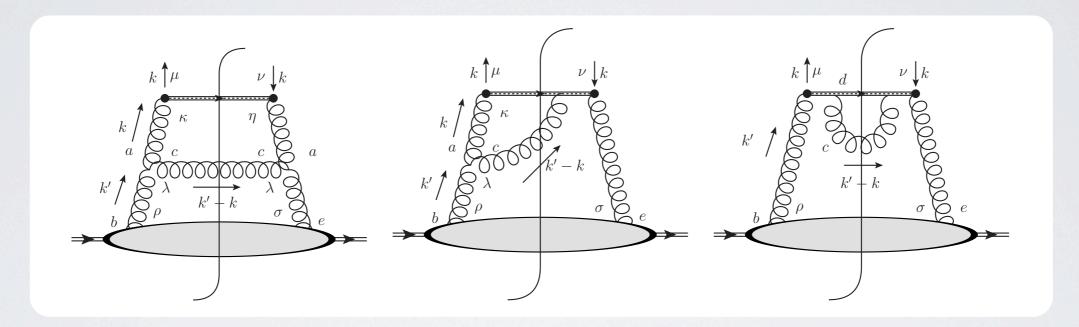
$$f_1^g(x, \mathbf{p}_T^2, \frac{3}{2}\sqrt{s}, M_h) = \frac{A_0 M_0^2}{M_0^2 + \mathbf{p}_T^2} \exp\left[-\frac{\mathbf{p}_T^2}{a\mathbf{p}_T^2 + 2\sigma^2}\right]$$

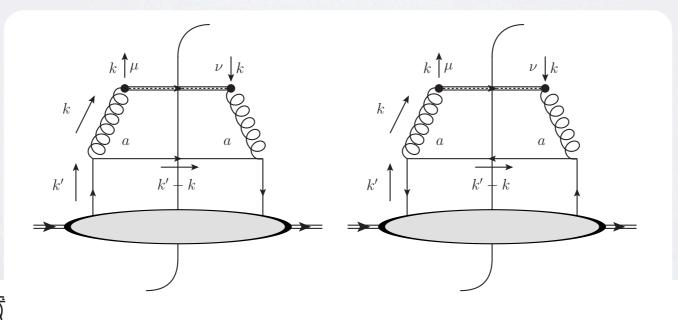




POWHEG+Pythia 8 Higgs qT distribution

$$\Phi_g^{\mu\nu}(x, \boldsymbol{p}_T, \zeta, \mu) \equiv 2 \int \frac{\mathrm{d}(\xi \cdot P) \,\mathrm{d}^2 \xi_T}{(xP \cdot n)^2 (2\pi)^3} \, e^{i(xP + p_T) \cdot \xi} \, \mathrm{Tr}_c \Big[\langle P | F^{n\nu}(0) \, \mathcal{U}_{[0,\xi]}^{n[-]} \, F^{n\mu}(\xi) \, \mathcal{U}_{[\xi,0]}^{n[-]} | P \rangle \Big]_{\xi \cdot P' = 0}$$

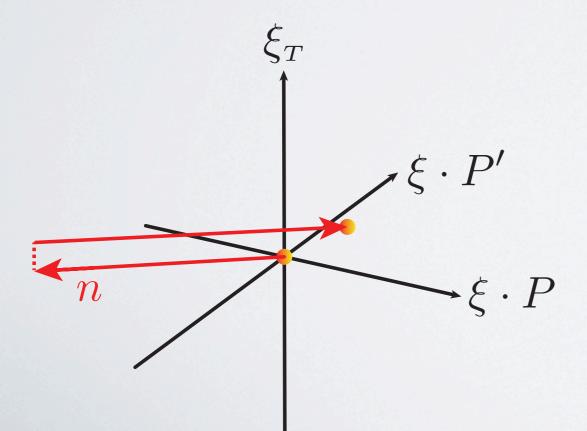




$$\Phi_{g}^{\mu\nu}(x, \boldsymbol{p}_{T}, \zeta, \mu) \equiv 2 \int \frac{\mathrm{d}(\xi \cdot P) \,\mathrm{d}^{2}\xi_{T}}{(xP \cdot n)^{2}(2\pi)^{3}} \, e^{i(xP + p_{T}) \cdot \xi} \, \mathrm{Tr}_{c} \Big[\langle P|F^{n\nu}(0) \, \mathcal{U}_{[0,\xi]}^{n[-]} \, F^{n\mu}(\xi) \, \mathcal{U}_{[\xi,0]}^{n[-]} |P\rangle \Big]_{\xi \cdot P' = 0}$$

$$\mathcal{U}_{[0,\xi]}^{n[-]} \equiv \mathcal{U}_{[0,-\infty]}^{n} \, \mathcal{U}_{[-\infty,\xi]}^{n}$$

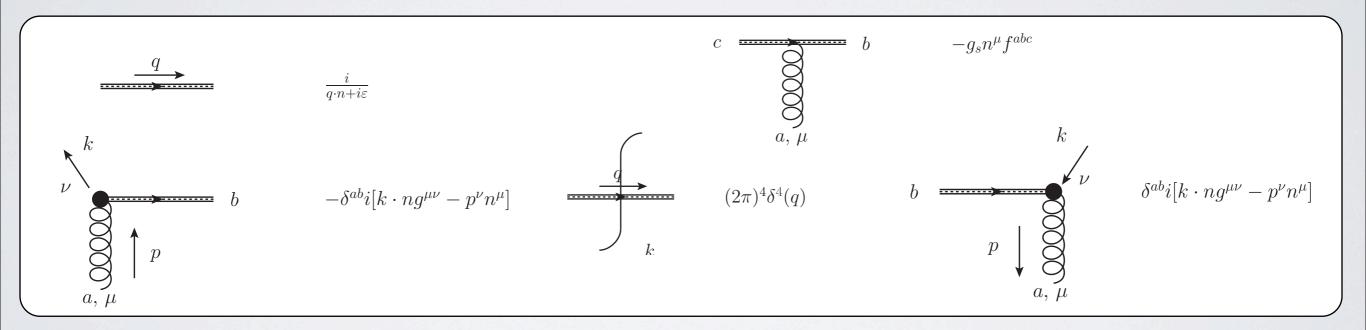
$$= \mathcal{P}e^{-ig \int_{0}^{-\infty} dy A^{n}(0+ny)} \mathcal{P}e^{-ig \int_{-\infty}^{0} dy A^{n}(\xi+ny)}$$



$$n \propto \frac{1}{\zeta} P + \frac{\zeta}{P \cdot P'} P'$$

$$\zeta \equiv \frac{2(n \cdot P)}{\sqrt{n^2}}$$

$$\Phi_g^{\mu\nu}(x, \boldsymbol{p}_T, \zeta, \mu) \equiv 2 \int \frac{\mathrm{d}(\xi \cdot P) \,\mathrm{d}^2 \xi_T}{(xP \cdot n)^2 (2\pi)^3} \, e^{i(xP + p_T) \cdot \xi} \, \mathrm{Tr}_c \Big[\langle P | F^{n\nu}(0) \, \mathcal{U}_{[0,\xi]}^{n[-]} \, F^{n\mu}(\xi) \, \mathcal{U}_{[\xi,0]}^{n[-]} | P \rangle \Big]_{\xi \cdot P' = 0}$$



in $\zeta \rightarrow \infty$ limit results compatible with the literature

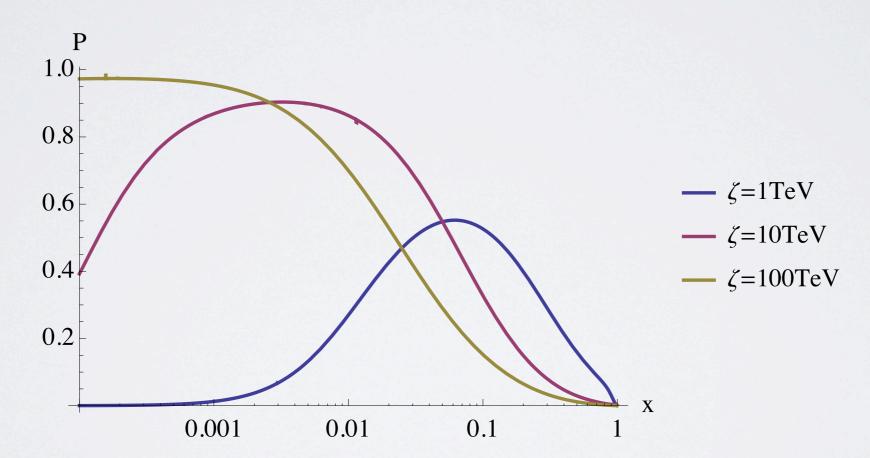
$$f_{1\text{tail}}^g(x, \mathbf{k}_T) = -\frac{4g_s^2}{(2\pi)^3 k_T^2} \int_x^1 \frac{\mathrm{d}z}{z} \left[C_A \left\{ \frac{1-z}{z} + z(1-z) + \frac{z}{(1-z)^+} \right. \right. \\ \left. - \frac{1}{2} \delta(z-1) \left[1 + \log \left(\frac{-k_T^2}{\zeta^2} \right) - \log x^2 (1-x)^2 \right] \right\} f_1^g \left(\frac{x}{z} \right) \\ \left. + C_F \frac{1 + (1-z)^2}{2z} \sum_{q,\bar{q}} f_1^q \left(\frac{x}{z} \right) \right]$$
 cf. Ji, Ma, Yuan JHEP07 (2005) 020

$$h_{1\text{tail}}^{\perp g}(x, \boldsymbol{k}_T) = \frac{4g_s^2}{(2\pi)^3 k_T^2} \frac{2M^2}{k_T^2} \int_x^1 \frac{\mathrm{d}z}{z} \frac{1-z}{z} \left[C_A f_1^g \left(\frac{x}{z} \right) + C_F \sum_{q,\bar{q}} f_1^q \left(\frac{x}{z} \right) \right]$$

cf. Sun, Xiao, Yuan PR D84, 094005 (2011)

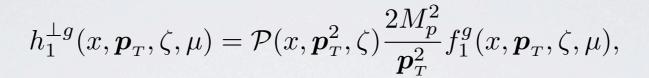
DEGREE OF POLARIZATION

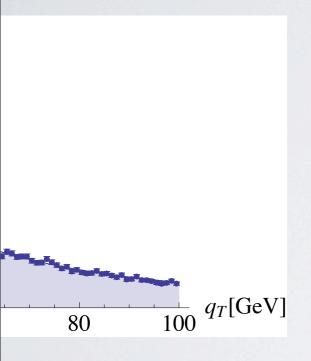
$$h_1^{\perp g}(x, \boldsymbol{p}_{\scriptscriptstyle T}, \zeta, \mu) = \mathcal{P}(x, \boldsymbol{p}_{\scriptscriptstyle T}^2, \zeta) \frac{2M_p^2}{\boldsymbol{p}_{\scriptscriptstyle T}^2} f_1^g(x, \boldsymbol{p}_{\scriptscriptstyle T}, \zeta, \mu),$$



CGC model predicts full polarization at small x: A. Metz and J. Zhou, Phys. Rev. D 84, 051503 (2011)

DEGREE OF POLARIZATION



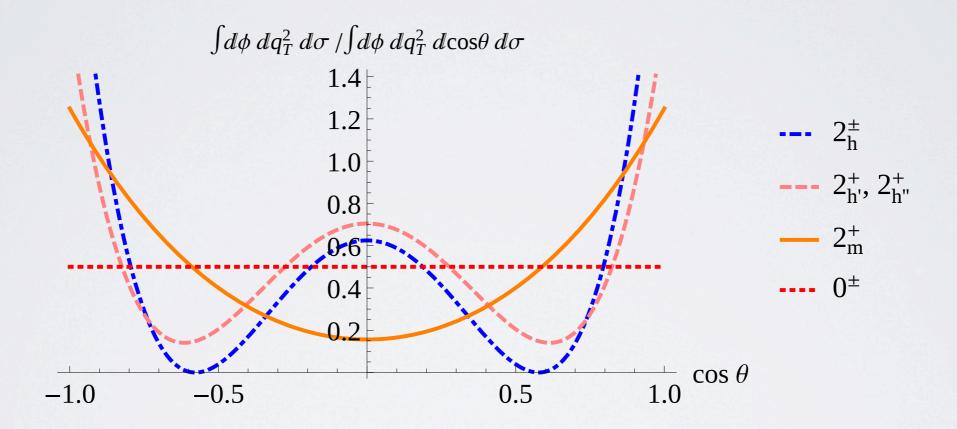


$$\mathcal{P}_{min} \equiv \frac{\boldsymbol{p}_{T}^{4}}{p_{0}^{4} + \boldsymbol{p}_{T}^{4}} 0.9 \, \mathcal{P}_{pQCD}(x, \boldsymbol{p}_{T}^{2}),$$

$$\mathcal{P} \equiv \mathcal{P}_{pQCD}(x, \boldsymbol{p}_{T}^{2}),$$

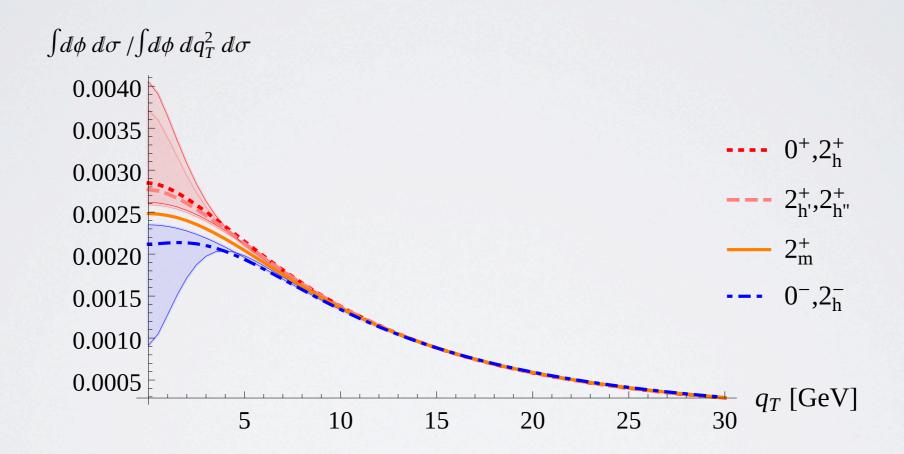
$$\mathcal{P}_{max} \equiv 1 - \frac{\boldsymbol{p}_{T}^{4}}{p_{0}^{4} + \boldsymbol{p}_{T}^{4}} \left[1 - 1.1 \, \mathcal{P}_{pQCD}(x, \boldsymbol{p}_{T}^{2})\right],$$

COSO DISTRIBUTION

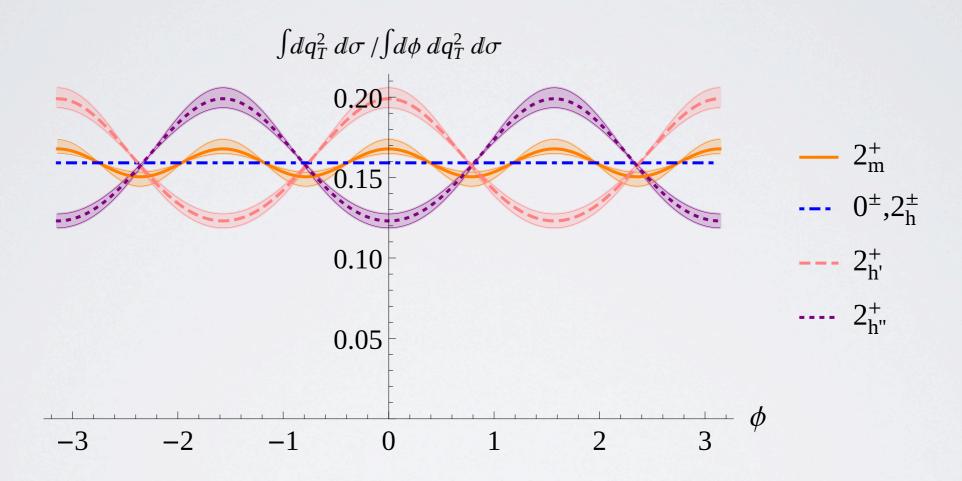


Donnerstag, 23. Mai 2013 22

TRANSVERSE MOMENTUM DISTRIBUTION



Ф DISTRIBUTION



REMARKS

• spin-2 with CP violating coupling ↔ sin2φ dependence

• gg \rightarrow box $\rightarrow \gamma\gamma$ background also ϕ dependent

• same can be done in the $H \rightarrow ZZ^*$ channel

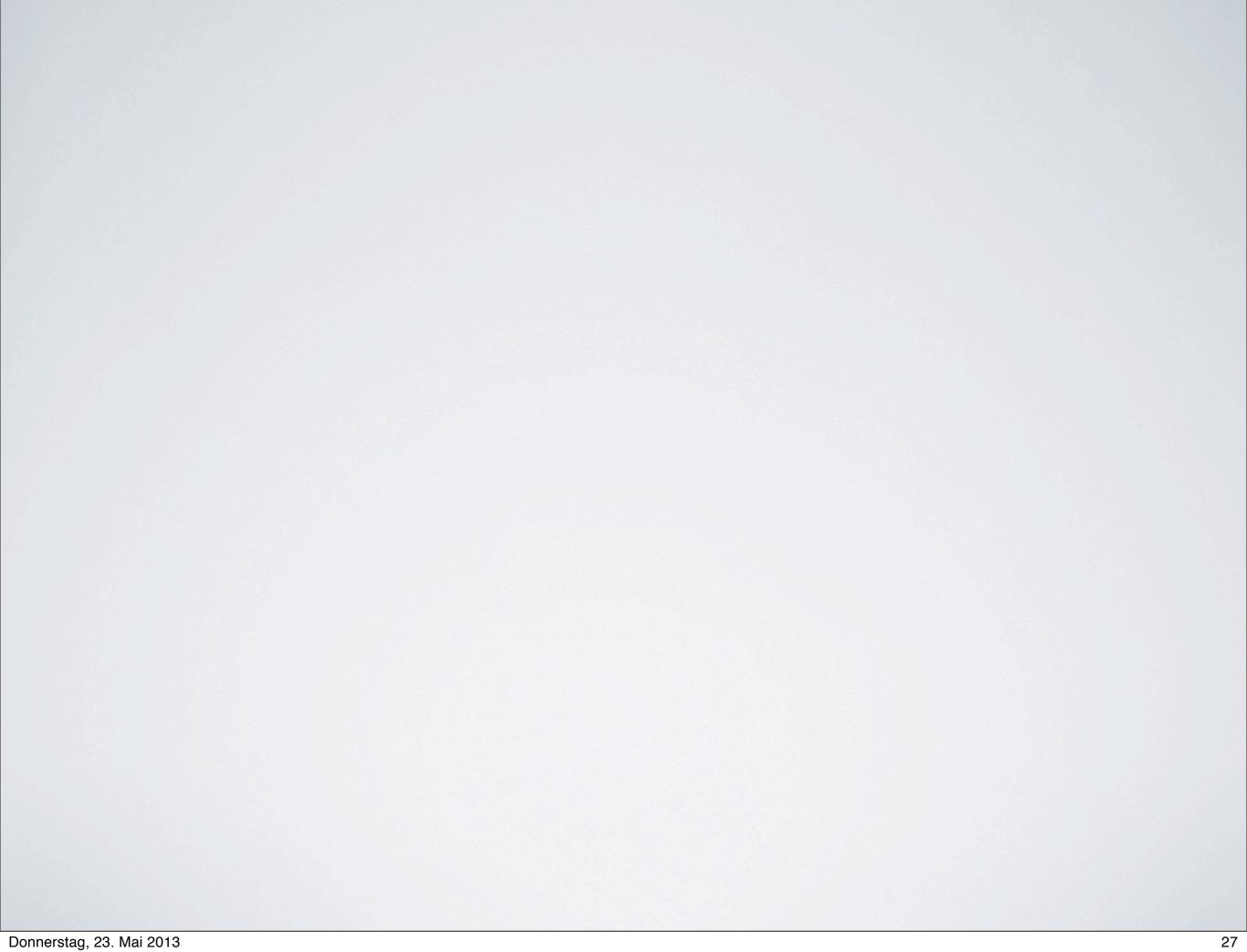
for gg → box → γγ see also
J. -W. Qiu, M. Schlegel and W. Vogelsang,
Phys. Rev. Lett. 107, 062001 (2011)

[arXiv:1103.3861 [hep-ph]]

CONCLUSIONS

- gluon polarization modifies both Higgs q_T and ϕ distribution
- q_T distribution modification different for positive/negative parity states
- φ distribution modification different for various spin-0 and spin-2 coupling scenarios

for more info see: arXiv:1304.2654



SPIN STATUS

May 15, 2013

0 ⁺ vs 2 _m ⁺	ATLAS	CMS
YY	99.3% CLs (2.7 σ)	_
/// *	94.7% CLs (1.9 σ)	86% CLs (1.5 σ)
ZZ*	83.4% CLs (1.4 σ)	98.5% CLs (2.4 σ)

PARITY STATUS

May 15, 2013

0 ⁺ vs 0 ⁻	ATLAS	CMS
YY	_	
W *	_	_
ZZ*	99.6% CLs (2.9 σ)	99.84% CLs (3.2 σ)