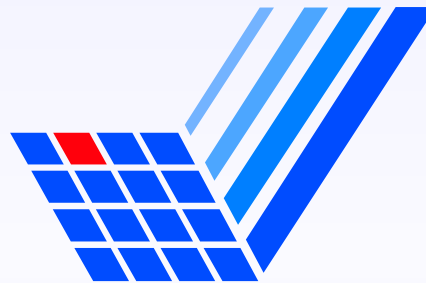


Radiatively Generated Isospin Violations in the Nucleon and the NuTeV Anomaly

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P. Jiménez Delgado

Theoretische Physik IV
Universität Dortmund

M. Gluck, P. J. D. and E. Reya, Phys. Rev. Lett. **95** (2005) 022002.

- **NuTeV** Weinberg angle $s_W^2 \equiv \sin^2 \theta_W \approx 3\sigma$ above world average.
- Among other sources, **isospin-symmetry violating pdfs** in the nucleon:

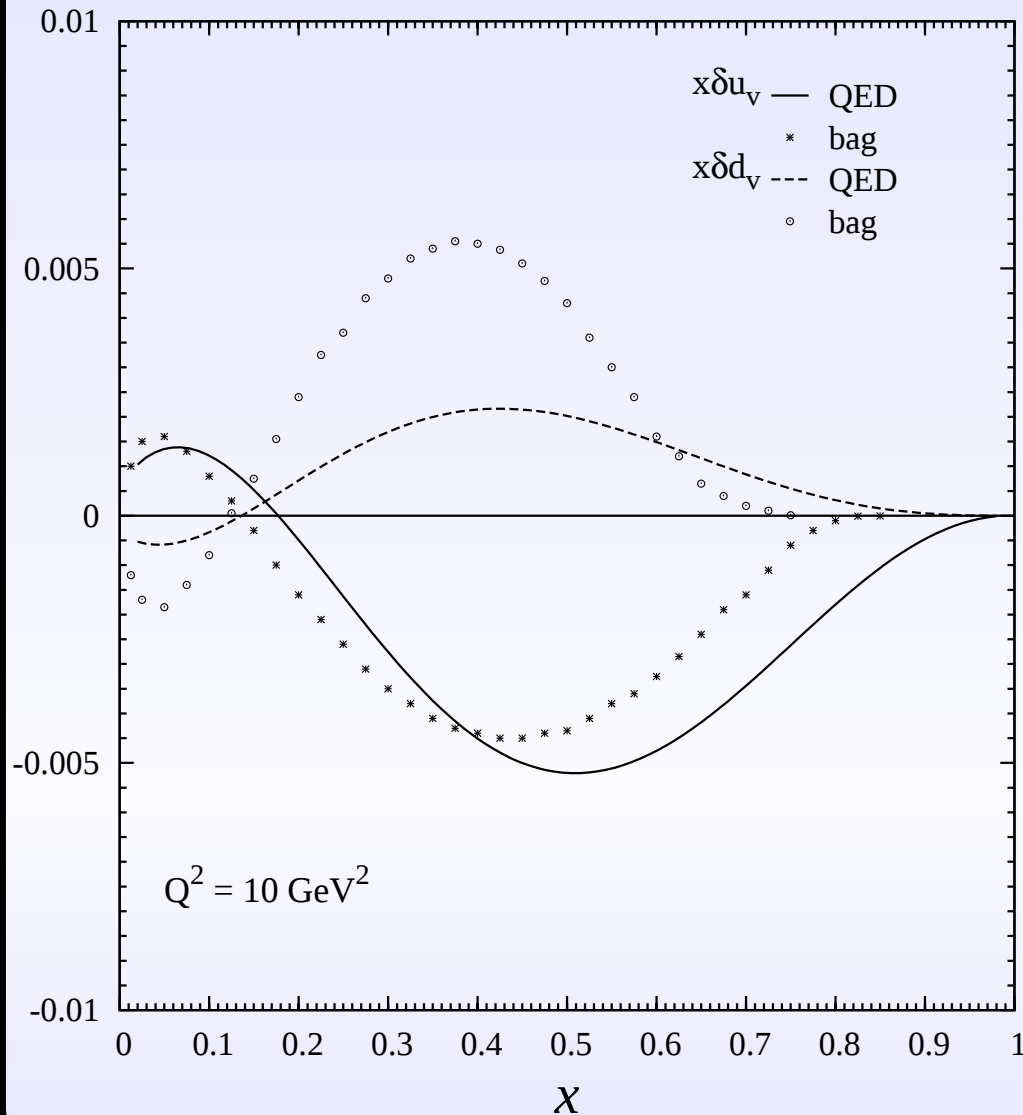
$$\begin{aligned}\delta^{(-)}_u(x, Q^2) &= {}^{(-)}u^p(x, Q^2) - {}^{(-)}d^n(x, Q^2) \\ \delta^{(-)}_d(x, Q^2) &= {}^{(-)}d^p(x, Q^2) - {}^{(-)}u^n(x, Q^2)\end{aligned}$$

- Nonperturbative valence asymmetries due to mass differences $\delta m = m_d - m_u$ and $\delta M = M_n - M_p$ reduce $\approx 30\%$ (**Londergan *et al.*** within **bag model**).
- Additional asymmetries from **radiative QED effects**, at $\mathcal{O}(\alpha)$:

$$\begin{aligned}\frac{d}{d \ln Q^2} \delta^{(-)}_u(x, Q^2) &= \frac{\alpha}{2\pi} \int_x^1 \frac{dy}{y} P\left(\frac{x}{y}\right) {}^{(-)}u(y, Q^2) \\ \frac{d}{d \ln Q^2} \delta^{(-)}_d(x, Q^2) &= -\frac{\alpha}{2\pi} \int_x^1 \frac{dy}{y} P\left(\frac{x}{y}\right) {}^{(-)}d(y, Q^2)\end{aligned} \quad P(z) = (e_u^2 - e_d^2) P_{qq}^\gamma(z)$$

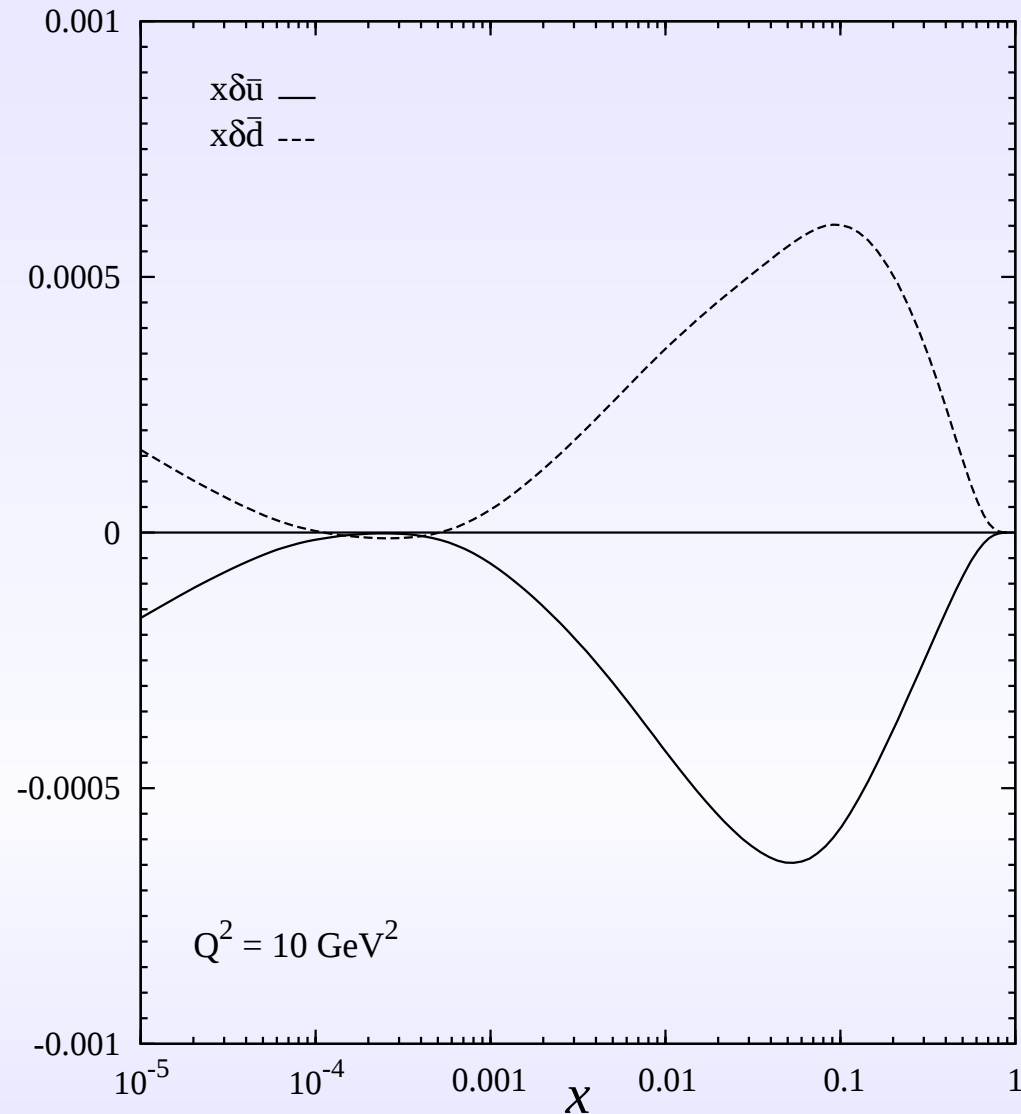
- Integrating with $m_q = 10$ MeV lower bound for a photon emitted by a quark:

$$\begin{aligned}\delta^{(-)}_u(x, Q^2) &= \frac{\alpha}{2\pi} \int_{m_q^2}^{Q^2} d \ln q^2 \int_x^1 \frac{dy}{y} P\left(\frac{x}{y}\right) {}^{(-)}u(y, q^2) \\ \delta^{(-)}_d(x, Q^2) &= -\frac{\alpha}{2\pi} \int_{m_q^2}^{Q^2} d \ln q^2 \int_x^1 \frac{dy}{y} P\left(\frac{x}{y}\right) {}^{(-)}d(y, q^2)\end{aligned} \quad \begin{array}{l} \text{with:} \\ {}^{(-)}q(y, q^2) = {}^{(-)}q(y, \mu^2) \\ \text{for } q^2 \leq \mu^2 \end{array}$$



- LO pdfs of the dynamical (radiative) parton model (GRV98) with $\mu^2 = 0.26 \text{ GeV}^2$.
- Comparable δu_v but different δd_v than bag model results from δm and δM .
- Similar result as MRST.

- Considerably smaller than valence asymmetries.
- Similar results obtained for CTEQ4 pdfs which have valence-like sea distributions at the input scale.
- May be tested in precision Drell–Yan and DIS experiments employing neutron as well as proton targets.



- **NuTeV** experiment: High precision (anti)neutrino-nucleon DIS. It is measured:

$$R^{\nu(\bar{\nu})}(x, Q^2) \equiv \frac{d^2 \sigma_{NC}^{\nu(\bar{\nu})N}(x, Q^2)}{d^2 \sigma_{CC}^{\nu(\bar{\nu})N}(x, Q^2)}, \quad N = \frac{1}{2}(p + n)$$

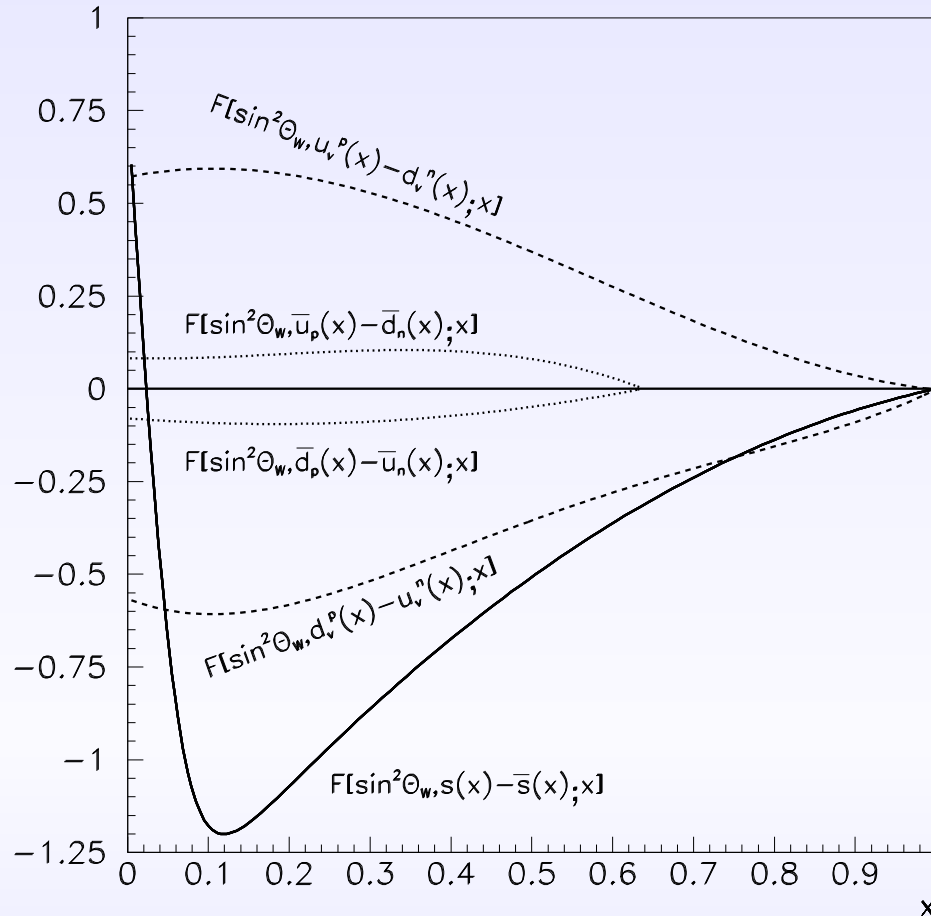
- Assuming $u = d$, $\bar{u} = \bar{d}$ and $s = \bar{s}$:

$$R^{\nu(\bar{\nu})} \equiv \frac{\sigma_{NC}^{\nu(\bar{\nu})N}}{\sigma_{CC}^{\nu(\bar{\nu})N}} = g_L^2 + r^{(-1)} g_R^2 \quad \text{with} \quad \begin{cases} g_L^2 = \frac{1}{2} - s_W^2 + \frac{5}{9} s_W^2 \\ g_R^2 = \frac{5}{9} s_W^2 \end{cases}$$

$$\Rightarrow \sin^2 \theta_W^{\text{NuTeV}} = 0.2277 \pm 0.0013 \pm 0.0009$$

- World average of electroweak measurements: $\sin^2 \theta_W^{\text{World}} = 0.2228(4)$.
- Analysis includes nonisoscalar target, higher-twists, charm production...
- It is explained considering **isospin violations** and **strange sea asymmetry** ($s > \bar{s}$ due to $M_\Lambda > M_K$).

- **Experimental** methods/functionals: $\Delta s_W^2 = \int_0^1 F[s_W^2, \delta q^{(-)}; x] x \delta q^{(-)}(x, Q^2) dx$



$$\sin^2 \theta_W^{\text{NuTeV}} = 0.2277(16) \longrightarrow \sin^2 \theta_W = 0.02234(16) \text{ agrees with } \sin^2 \theta_W^{\text{World}} = 0.2228(4)$$

- $\Delta s_W^2|_{\text{QED}} = -0.0011$ evaluated using our asymmetries $\delta q^{(-)}$.
- $\Delta s_W^2|_{\text{bag}} = -0.0015$
Londergan, Rodionov, Thomas, Sather, etc.
- $\Delta s_W^2|_{\text{strange}} = -0.0017$
J. Alwall and G. Ingelman.
- **Total shift** $\Delta s_W^2|_{\text{total}} =$
 $= \Delta s_W^2|_{\text{QED}} + \Delta s_W^2|_{\text{bag}} + \Delta s_W^2|_{\text{strange}}$
 $= -0.0043$

- For **simplicity**: Corrections to R_{PW}^- instead of $R^{\nu(\bar{\nu})}$:

$$R^- \equiv \frac{\sigma_{NC}^{\nu N} - \sigma_{NC}^{\bar{\nu} N}}{\sigma_{CC}^{\nu N} - \sigma_{CC}^{\bar{\nu} N}} = R_{PW}^- + \delta R_I^- + \delta R_s^- \text{ with } \begin{cases} R_{PW}^- = \frac{1}{2} - s_W^2 \\ \delta R_I^- \simeq \left(\frac{1}{2} - \frac{7}{6}s_W^2\right) \frac{\delta U_v - \delta D_v}{U_v + D_v} \\ \delta R_s^- \simeq -\left(1 - \frac{7}{3}s_W^2\right) \frac{S^-}{U_v + D_v} \end{cases}$$

where $\delta Q_v(Q^2) = \int_0^1 x \delta q_v(x, Q^2) dx$, $S^- = S - \bar{S}$, etc.

- This **overestimates** the corrections by $\approx 20\%$ – 40% :

$$\delta R_I^-|_{\text{QED}} = -0.0021 \text{ instead of } \Delta s_W^2|_{\text{QED}} = -0.0011$$

$$\delta R_I^-|_{\text{bag}} = -0.0020 \text{ instead of } \Delta s_W^2|_{\text{QED}} = -0.0015$$

$$\delta R_s^- = -0.0021 \text{ instead of } \Delta s_W^2|_{\text{strange}} = -0.0017$$

- $\delta R_{I,s}^-$ should be avoided and Δs_W^2 evaluated according to **NuTeV functionals** (sensitivity of their observables to cuts, etc.)

Summary

- We evaluate modifications $\delta \bar{q}^{(-)}(x, Q^2)$ to standard pdfs due to **QED $\mathcal{O}(\alpha)$ photon bremsstrahlung** within the dynamical (radiative) parton model.
- The **NuTeV Anomaly is removed** considering the impact of:
 - + these isospin asymmetries (QED).
 - + nonperturbative isospin asymmetries: quark and target mass differences.
 - + strangeness asymmetry ($s \neq \bar{s}$).
- Naive corrections to the PW relation **overestimate Δs_W^2** .

Research Status

- Working on NLO $\overline{\text{MS}}$ global analysis (radiatively generated, valence-like).
- Working on factorization scheme dependence of pdfs: DIS scheme analysis.

Δs_W^2	δu_v	δd_v	$\delta \bar{u}$	$\delta \bar{d}$	total
QED	-0.00071	-0.00033	-0.000019	-0.000023	-0.0011
bag	-0.00065	-0.00081	—	—	-0.0015